

ADA 071749

AFFDL-TR-79-3045

Volume II

ANGULAR VIBRATION OF AIRCRAFT
Volume II — Prediction Methods
for Angular Vibration

ANAMET LABORATORIES, INC.
APPLIED MECHANICS DIVISION
SAN CARLOS, CALIFORNIA 94070

April 1979

TECHNICAL REPORT AFFDL-TR-79-3045, Volume II
Final Report for Period August 1977 to April 1979

Approved for public release; distribution unlimited.

AIR FORCE FLIGHT DYNAMICS LABORATORY
AIR FORCE WRIGHT AERONAUTICAL LABORATORIES
AIR FORCE SYSTEMS COMMAND
WRIGHT-PATTERSON AIR FORCE BASE, OHIO 45433


20080818 020


NOTICE


When Government drawings, specifications, or other data are used for any purpose other than in connection with a definitely related Government procurement operation, the United States Government thereby incurs no responsibility nor any obligation whatsoever; and the fact that the government may have formulated, furnished, or in any way supplied the said drawings, specifications, or other data, is not to be regarded by implication or otherwise as in any manner licensing the holder or any other person or corporation, or conveying any rights or permission to manufacture, use, or sell any patented invention that may in any way be related thereto.

This report has been reviewed by the Information Office (OI) and is releasable to the National Technical Information Service (NTIS). At NTIS, it will be available to the general public, including foreign nations.

This technical report has been reviewed and is approved for publication.


LT MICHAEL W. OBAL
Project Engineer


RALPH N. BINGMAN
Technical Manager


RALPH L. KUSTER, JR., Colonel, USAF
Chief, Structures and Dynamics Division

Copies of this report should not be returned unless return is required by security considerations, contractual obligations, or notice on a specific document.

UNCLASSIFIED

SECURITY CLASSIFICATION OF THIS PAGE (When Data Entered)

REPORT DOCUMENTATION PAGE		READ INSTRUCTIONS BEFORE COMPLETING FORM
1. REPORT NUMBER AFFDL-TR-3045, Volume II	2. GOVT ACCESSION NO.	3. RECIPIENT'S CATALOG NUMBER
4. TITLE (and Subtitle) ANGULAR VIBRATION OF AIRCRAFT Volume II - Prediction Methods for Angular Vibration		5. TYPE OF REPORT & PERIOD COVERED Final Report 8/77 to 4/79
		6. PERFORMING ORG. REPORT NUMBER
7. AUTHOR(s) Conor D. Johnson Ernest B. Paxson Warren C. Gibson David Kienholz		8. CONTRACT OR GRANT NUMBER(s) F33615-77-C-3050
9. PERFORMING ORGANIZATION NAME AND ADDRESS Anamet Laboratories, Inc. 100 Industrial Way San Carlos, California 94070		10. PROGRAM ELEMENT, PROJECT, TASK AREA & WORK UNIT NUMBERS Project 2401 Task 240104 W.U. 24010408
11. CONTROLLING OFFICE NAME AND ADDRESS Air Force Flight Dynamics Laboratory AFSC/AFWAL, United States Air Force Wright-Patterson AFB, Ohio 45433		12. REPORT DATE April 1979
		13. NUMBER OF PAGES 433
14. MONITORING AGENCY NAME & ADDRESS (if different from Controlling Office)		15. SECURITY CLASS. (of this report) UNCLASSIFIED
		15a. DECLASSIFICATION/DOWNGRADING SCHEDULE
16. DISTRIBUTION STATEMENT (of this Report) Approved for public release; distribution unlimited.		
17. DISTRIBUTION STATEMENT (of the abstract entered in Block 20, if different from Report)		
18. SUPPLEMENTARY NOTES		
19. KEY WORDS (Continue on reverse side if necessary and identify by block number) Angular Vibration Angular Measurement Finite Element NASTRAN Statistical Energy Analysis Structural Test Structural Analysis Semi-Loof Structural Dynamics		
20. ABSTRACT (Continue on reverse side if necessary and identify by block number) This report describes development work in several distinct areas, all related to prediction of angular vibration of aircraft structures. Angular vibration in this context refers to dynamic rotations or changes in slope at specific points on a vibrating structure. It is of interest primarily in connection with high resolution optical and electro-optical systems. Efforts were directed at both low frequency vibration, where individual normal modes are known, as well as high frequency vibration where they are not. For low frequency predictions, improved accuracy per unit cost was sought by an evolutionary improvement to an		

UNCLASSIFIED

SECURITY CLASSIFICATION OF THIS PAGE (When Data Entered)

UNCLASSIFIED

SECURITY CLASSIFICATION OF THIS PAGE(When Data Entered)

existing finite element code. A state-of-the-art shell element, the Semi-Loof, was incorporated into NASTRAN by means of pre- and post-processors and DMAP instructions. Its accuracy is tested against both closed form solutions for simple cases and experiment for an actual fuselage. For high frequency predictions the method of Statistical Energy Analysis was pursued. A demonstration case involving two coupled plates is presented to show how SEA may be used to predict angular vibration in a situation where normal modes are too numerous to be predicted individually. Relationships between linear and angular vibration were developed for various structural forms. Theoretical error bounds are also derived for spectral measurements of angular vibration which are obtained by differencing of signals from translational sensors.

UNCLASSIFIED

SECURITY CLASSIFICATION OF THIS PAGE(When Data Entered)

FOREWORD

This report describes an investigation into the prediction methods of angular vibration of aircraft, performed by Anamet Laboratories, Inc., San Carlos, California, for the Air Force Flight Dynamics Laboratory, Air Force Systems Command, Wright-Patterson Air Force Base, Ohio, under Contract F33615-77-C-3050. This research was conducted under Project 2401, "Structural Mechanics," Task 240104, "Vibration Prediction and Control, Measurement and Analysis," Work Unit 24010408, "Angular Vibration of Aircraft." Lt. Michael W. Obal, (AFFDL/FBG) was the project engineer. This report is in two volumes: Volume I - Executive Summary; Volume II - Prediction Methods for Angular Vibration.

This program was conducted by the Applied Mechanics Division of Anamet Laboratories. Program Manager was Dr. Conor Johnson and Principal Investigators were Dr. Warren Gibson, Dr. David Kienholz, and Dr. Ernest Paxson. This research was performed between August 1977 and April 1979.

This report was submitted by the authors on 2 April 1979 for publication as an AFFDL Technical Report.

TABLE OF CONTENTS

<u>Section</u>	<u>Page</u>
I. INTRODUCTION	1
II. BACKGROUND	7
2.1 LITERATURE SEARCH	7
2.2 SELECTION OF LOW FREQUENCY METHODS	10
2.3 SELECTION OF APPROACH FOR HIGH FREQUENCY PREDICTION	15
III. LOW FREQUENCY METHODS	18
3.1 FINITE ELEMENT REVIEW	19
3.2 THE SEMI-LOOF ELEMENTS	24
3.3 COMPARISON AND EVALUATION OF SEMI-LOOF ELEMENTS	31
3.4 IMPLEMENTATION INTO NASTRAN	37
3.5 SUMMARY	39
IV. HIGH FREQUENCY METHODS	40
4.1 INTRODUCTION	40
4.1.1 Statistical Energy Analysis	42
4.2 FUNDAMENTALS OF SEA	43
4.2.1 SEA Parameters	48
4.3 ESTIMATION OF COUPLING LOSS FACTOR BY THE WAVE TRANSMISSION METHOD	50
4.3.1 Derivation in Terms of Waves	50
4.3.2 Derivation in Terms of Normal Modes	52
4.4 ESTIMATION OF ANGULAR RESPONSE	61
4.4.1 Relation Between Angular Response and Total Energy for a Uniform Plate	61
4.5 ESTIMATION OF INTERNAL LOSS FACTOR	67
4.5.1 Measurement and Use of Modal Damping Ratios	68

TABLE OF CONTENTS (Continued)

<u>Section</u>	<u>Page</u>
4.6 SEA WAVE TRANSMISSION EXPERIMENT	71
4.6.1 Objective	71
4.6.2 Experimental Hardware	73
4.6.3 Experimental Procedure	77
4.6.4 Experimental Results	80
4.7 CONCLUSION	94
4.7.1 Summary	94
4.7.2 Suggestions for Future Work	95
V. RELATIONSHIPS BETWEEN LINEAR AND ANGULAR VIBRATION IN AIRCRAFT STRUCTURAL COMPONENTS . . .	97
5.1 LINEAR-TO-ANGULAR RELATIONSHIP OF SPRING- SUPPORTED RIGID BAR SUBJECTED TO TEMPORALLY RANDOM CONCENTRATED LOAD	98
5.2 BEHAVIOR OF BEAMS SUBJECTED TO TEMPORALLY RANDOM LOADING CONDITIONS	106
5.3 NASTRAN ANALYSIS OF STIFFENED CURVED PANEL SUBJECTED TO TEMPORALLY RANDOM CONCENTRATED LOAD	113
5.4 VARIATION OF ANGULAR-TO-LINEAR DISPLACEMENT RATIOS IN STRUCTURAL COMPONENTS	128
VI. ANGULAR VIBRATION MEASUREMENT	130
6.1 THEORETICAL ERROR ANALYSIS	132
6.1.1 Narrow-Band Signal-to-Noise Ratio . .	132
6.1.2 Coherence Function for Angular Frequency Response	149
6.1.3 Typical Value of Spectral Discrete Difference	152
6.1.4 Flexural Error	154
6.2 EXAMPLES	159
6.2.1 Use of Error Formulas	159
6.2.2 Effect of Single Channel Noise . . .	163
6.2.3 Effect of Interchannel Gain and Phase Mismatch	164

TABLE OF CONTENTS (Continued)

<u>Section</u>	<u>Page</u>
6.3 CONCLUSION	171
VII. TEST AND EVALUATION OF PREDICTION METHOD . . .	172
7.1 DESCRIPTION OF THE TEST OBJECT	173
7.2 FINITE ELEMENT MODEL	179
7.3 TEST DESCRIPTION	182
7.3.1 Overview	182
7.3.2 Test Procedure	183
7.3.3 Test Equipment	191
7.4 RESULTS	194
VIII. SUMMARY AND CONCLUSIONS	215
REFERENCES	218
 APPENDIX A - SEMI-LOOF USER'S MANUAL	 221
A.1 INTRODUCTION	221
A.2 GENERAL DISCUSSION	222
A.3 PRELOOF INPUT	229
A.4 POSTLOOF INPUT	247
A.5 DECK SETUP	249
A.6 PROGRAMMER'S NOTES	256
 APPENDIX B - SEMI-LOOF PROGRAM LISTINGS	 261
B.1 CONTROL CARD PROCEDURES	262
B.2 DMAP ALTER LIBRARY (Partial Listing)	264
B.3 PRELOOF PROGRAM LISTING	267
B.4 POSTLOOF LISTING	373
 APPENDIX C - DETAILED DESCRIPTION OF SOFTWARE USED IN STATISTICAL ENERGY ANALYSIS EXPERIMENT . . .	 376
C.1 GENERAL	376
C.2 SUBROUTINE DESCRIPTIONS	377

TABLE OF CONTENTS (Concluded)

<u>Section</u>	<u>Page</u>
C.3 CALLING FORMATS	382
C.4 PROGRAM LISTING	389
APPENDIX D - FUSELAGE MODEL LISTING	397

LIST OF ILLUSTRATIONS

<u>Figure</u>		<u>Page</u>
1	Finite Element Mathematical Derivation	22
2	The Patch Test	23
3	Semi-Loof Quadrilateral Element Degrees of Freedom	26
4	Semi-Loof Triangular Element Degrees of Freedom	29
5	Semi-Loof Beam Element Degrees of Freedom	30
6	Cantilever Cylinder	32
7	Point-loaded Plate Evaluation	32
8	Plate Element Convergence Study	33
9	Curved Panel Configuration	36
10	Cumulative Power of Typical Translational Variable for In-flight Vibration	41
11	Cumulative Power of Typical Rotational Variable for In-flight Vibration	41
12	Two Bodies with Stiffness Coupling and Direct Excitation of One Body	45
13	Couples Plates Used for SEA Wave Transmission Experiment	74
14	Close-up of Connecting Link Instrumented to Measure Instantaneous Force and Acceleration	75
15	Normalized Transmitted Power $\langle \pi_{AB} \rangle_t / \langle E_A \rangle_t$	81
16	Normalized Energy of Indirectly Excited Body. Actual Damping Used.	82
17	Smoothed Power Spectral Density and R.M.S. Value of High-Frequency Angular Vibration at a Typical Point of a Uniform Rectangular Plate	84

LIST OF ILLUSTRATIONS (Continued)

<u>Figure</u>		<u>Page</u>
18	Internal Loss Factor vs. Frequency of Indirectly Excited Plate	85
19	Modulus of Acceleration Admittance at Center Point of Large Plate	88
20	Measured Coupling Point Mobility of Small Plate After Smoothing	89
21	Spectral Density of Transmitted Power, Smoothed and Unsmoothed Versions	90
22	Smoothed Spectral Density of Total Vibration Energy	91
23	Normalized Transmitted Power with Light Coupling Spring	93
24	Spring Supported Rigid Bar	98
25	Plot of θ_{rms}/y_{rms} for Simply-supported and Free-Free Beams Subjected to Temporally Random Loads	108
26	NASTRAN Model of Stiffened Curved Panel	114
27	Ratio of RMS Resultant Angular Displacement to RMS Linear Displacement vs. Grid Point Station for Curved Plate with Stiffeners, 20-2560 Hz.	116
28	Ratio of RMS Angular Displacement to Spatially Average RMS Linear Displacement vs. Grid Point Station for Curved Plate with Stiffeners, 20-2560 Hz.	117
29	Ratio of RMS Angular Displacement to RMS Linear Displacement vs. Grid Point Station for Curved Plate with Stiffeners, 20-2560 Hz.	120
30	Ratio of RMS Angular Displacement to RMS Linear Displacement vs. Grid Point Station for Curved Plate with Stiffeners, 20-40 Hz.	121

LIST OF ILLUSTRATIONS (Continued)

<u>Figure</u>		<u>Page</u>
31	Ratio of RMS Angular Displacement to RMS Linear Displacement vs. Grid Point Station for Curved Plate with Stiffeners, 40-80 Hz.	122
32	Ratio of RMS Angular Displacement to RMS Linear Displacement vs. Grid Point Station for Curved Plate with Stiffeners, 80-160 Hz.	123
33	Ratio of RMS Angular Displacement to RMS Linear Displacement vs. Grid Point Station for Curved Plate with Stiffeners, 160-320 Hz.	124
34	Ratio of RMS Angular Displacement to RMS Linear Displacement vs. Grid Point Station for Curved Plate with Stiffeners, 320-640 Hz.	125
35	Ratio of RMS Angular Displacement to RMS Linear Displacement vs. Grid Point Station for Curved Plate with Stiffeners, 640-1280 Hz.	126
36	Ratio of RMS Angular Displacement to RMS Linear Displacement vs. Grid Point Station for Curved Plate with Stiffeners, 1280-2560 Hz.	127
37	Measurement of Angular Acceleration by Differencing	133
38	Narrow-band Signal-to-Noise Ratio of Angular Acceleration Estimate Obtained by Differencing of Linear Accelerations $ \delta = 0.000$	142
39	Narrow-band Signal-to-Noise Ratio of Angular Acceleration Estimate Obtained by Differencing of Linear Accelerations $ \delta = 0.005$	143
40	Narrow-band Signal-to-Noise Ratio of Angular Acceleration Estimate Obtained by Differencing of Linear Accelerations $ \delta = 0.01$	144
41	Narrow-band Signal-to-Noise Ratio of Angular Acceleration Estimate Obtained by Differencing of Linear Accelerations $ \delta = 0.015$	145
42	Typical Narrow-band Signal-to-Noise for Piezoelectric Accelerometers	147

LIST OF ILLUSTRATIONS (Continued)

<u>Figure</u>		<u>Page</u>
43	Dimensionless Difference Signal $1-\alpha$ vs. Mode Shape ψ	155
44	Upper Frequency Limit for Differential Sensing of Angular Vibration	158
45	Noise PSD for Two Different Piezoelectric Accelerometers	160
46	Second Bending Mode in Slope Coordinates of Free-Free Beam Measured with Accelerometers of Different Noise Performance	165
47	Test Structure for Evaluating Angular Acceleration Measurement Methods	167
48	Measured PSD of Angular Acceleration of a Rigid Body with One Axis Fixed	170
49	Exterior View of Test Fuselage in As-Received Condition	174
50	Interior View of Test Fuselage in As-Received Condition	175
51	External View of Prepared Test Fuselage	176
52	Intenal View of Prepared Test Fuselage	177
53	Rear View of Prepared Test Fuselage Showing Construction Details	178
54	Semi-Loof Fuselage Model	180
55	Sketch of Test Object Showing Input Force Location and Direction (Arrow), Translational Response Sensor \bullet , and Rotational Response Sensor θ	184
56	Front View Sketch of Test Object Showing Input Force Location and Directional (Arrow), Translational Response Sensor \bullet , and Rotational Response Sensor θ	185
57	Partial Sketch of Test Object Showing Location of Rotational Sensor No. 5	186
58	Schematic of Test Setup	192

LIST OF ILLUSTRATIONS (Continued)

<u>Figure</u>		<u>Page</u>
59	Experimental Results, Response Point 1, Y Translation, Driving Point 1, Y Translation	197
60	Experimental Results, Response Point 21, Z Translation, Driving Point 1, Y Translation	198
61	Experimental Results, Response Point 22, Rotation about Y, Driving Point 1, Y Translation	199
62	Experimental Results, Response Point 23, Rotation about X, Driving Point 1, Y Translation	200
63	Experimental Results, Response Point 4, Rotation about X, Driving Point 2, Y Translation	201
64	Experimental Results, Response Point 5, Rotation about X, Driving Point 2, Y Translation	202
65	Experimental Results, Response Point 12, Rotation about Y, Driving Point 2, Y Translation	203
66	Experimental Results, Response Point 21, Z Translation, Driving Point 2, Y Translation	204
67	Experimental Results, Response Point 23, Rotation about X, Driving Point 2, Y Translation	205
68	Experimental Results, Response Point 32, Rotation about Z, Driving Point 2, Y Translation	206
69	Experimental Results, Response Point 22, Rotation about X, Driving Point 3, X Translation	207
70	Experimental Results, Response Point 23, Rotation about Y, Driving Point 3, X Translation	208
71	Analytical Results, Response Point 22, Rotation about Y, Driving Point 1, Y Translation	209
72	Analytical Results, Response Point 23, Rotation about X, Driving Point 1, Y Translation	210

LIST OF ILLUSTRATIONS (Concluded)

<u>Figure</u>		<u>Page</u>
73	Analytical Results, Response Point 32, Rotation about Z, Driving Point 2, Y Translation	211
74	Analytical Results, Response Point 22, Rotation about Y, Driving Point 3, X Trans- lation	212
75	Analytical Results, Response Point 23, Rotation about X, Driving Point 3, X Trans- lation	213
A-1	Stiffened Panel Properties	226
A-2	Local Coordinates for LOOF8 and LOOF6 Elements	227
A-3	Flow of Control from PRELOOF to NASTRAN to POSTLOOF	252

LIST OF TABLES

<u>Table</u>		<u>Page</u>
1	Key Words Used for High-Frequency Vibration Literature Search	9
2	Natural Frequencies and Modal Damping Ratios for Plate B	86
3	θ_{rms}/y_{rms} for Free-Free Beam Subject to Concentrated Load	111
4	Test Conditions for Short Beam Experiment Used to Investigate Mismatching Error	169
5	Description of Response Variables	187
6	Measured Frequency Responses	188
7	Test Equipment	193
8	Natural Frequencies of the Fuselage	195
9	RMS Values of PSD Responses of the Fuselage	196
A-1	Summary of PREL00F Input Deck	232

LIST OF SYMBOLS

a	Coordinate location ($x=a$) of concentrated random load
A	Cross-section area of beam (Section V)
A	Amplitude coefficient on the noise power spectrum (Section VI)
A_i	Stiffener area
A_{ij}, B_{ij}, D_{ij}	Force-strain parameters for stiffened shells
A_{ijkl}	Finite element master matrix
A_p	Plate area
b_i	Stiffener spacing
B	Amplitude coefficient on the signal power spectrum
B_{ij}	Finite element master vector
c	Transducing scale factor, engineering units/volt (Section VI)
c, c_1, c_2	Spring damping constants (Section V)
$c(f)$	Nominal frequency response of transducing channel
c_b	Phase velocity of flexural wave
c_ℓ	Extentional wave speed for a material, $\sqrt{E/\rho}$
C_i	Stiffener eccentricity
$C_i(f)$	Complex functions of frequency defined by Eq. (6.17)
$C_1 = \frac{\Lambda}{k} K_1$	Modal damping factors
$C_2 = \frac{\Lambda}{k} K_2$	
D	Amplitude of flexural wave
$e_i(t)$	Voltage produced by transducer channel i
e_o	Difference signal

LIST OF SYMBOLS (continued)

\vec{e}_x, \vec{e}_y	Unit vectors in x and y directions
E	Modulus of elasticity
$E_i(f)$	Fourier transform of $e_i(t)$
$E_i(t)$	Energy of i^{th} mode (in Section 4.4.1)
$E^{(i)}(t)$	Energy of i^{th} mode (in Section 4.5.1)
$E_A(f)$	Spectral density of total energy of body A
$E_A(t)$	Total energy (kinetic plus potential) of body A
f	Frequency in cycles/unit time
f_b	Center frequency of band b
f_c	Center frequency of a specified band
$f_i, f^{(i)}$	i^{th} natural frequency
$f_{l.m.}$	Log mean frequency of a band
f_r	r^{th} resonant frequency
F	Energy density interior to a structure
F_1, F_2	Forces at ends of rigid bar
g	Gravitational constant (396.04 lb _m -in/lbf-sec. ²)
G	Energy density on the boundary of a structure (Section III)
G	Shear modulus (Appendix A)
h	Plate thickness
$h_i(\tau)$	Impulse response to transducer channel i
h_1, h_2	Response single degree-of-freedom system to unit impulse
H(f)	Fourier transform of $h(\tau)$
H_i	Weighting factors for Gaussian integration
$H_i(f), c_i(f)$	Frequency response of transducer channel i

LIST OF SYMBOLS (continued)

$H_{oi}(f)$	Displacement/force frequency response between points o and i
H_A	Mechanical admittance of body A at a coupling degree of freedom ($=Z_A^{-1}$)
$H_1(\omega), H_2(\omega)$	Complex frequency responses
$H_1^*(\omega), H_2^*(\omega)$	Complex conjugates of H_1 and H_2 , respectively
i	$\sqrt{-1}$, also used as the mode index
I	Area moment of inertia of cross-section of beam about neutral axis
I_c	Mass moment of inertia for rigid bar about CG
I_i	Stiffener moment of inertia
J_i	Stiffener torsion constant
k	Exponent of envelope of signal power spectrum (Section VI)
k, k_1, k_2	Spring stiffness constants (Section V)
k_c	Stiffness of coupling spring
k_{xi}, k_{yi}	Wavenumbers in the x and y direction corresponding to the i^{th} mode
K_1, K_2	Modal stiffness constants
$\ell(t)$	Force input
$\ell_a(t), \ell_b(t)$	Coupling forces acting at points a and b of bodies A and B
$\ell_c(t)$	Force in a massless coupling element
$\ell_o(t)$	Prescribed input force
L	Length of rigid bar, length of finite beam
$L(f)$	Fourier transform of $\ell(t)$
$L_a(f), L_b(f)$	Fourier transforms of loads $\ell_a(t)$ and $\ell_b(t)$

LIST OF SYMBOLS (continued)

L_1, L_2	Distances from opposite ends to CG of rigid bar
m	Mass per unit length of beam (Section V)
m	Exponent of envelope of noise power spectrum (Section VI)
m_i	i^{th} diagonal entry of mass matrix \tilde{M} (Section 4.2) or i^{th} modal mass (Section 4.4.1)
m_A	Total mass of body A
M	Mass of rigid bar
\tilde{M}	Mass matrix
M_o	Minimum allowable value of flexural modulation function
M_1, M_2	Elements of mass matrix for spring-supported rigid bar system (Section V)
M_1, M_2, M_{12}	Shell moment resultants (Appendix A)
$n_i(t)$	Noise in channel i
$N_i(f)$	Fourier transform of $n_i(t)$
$N_{ij}(\xi, \eta)$	Finite element shape functions
N_A	Number of normal modes of body A within a specified frequency band
N_1, N_2, N_{12}	In-plane shell force resultants
$P(t)$	Temporally random concentrated load
$q(x)$	Spatial distribution of load on beam
q_n, q_k	Fourier coefficients for load function on beam
q_o	Constant representing load intensity on beam
q_1, q_2	Normal coordinates for spring supported rigid bar system
$R_{v\ell}(t)$	Cross correlation between $v(t)$ and $\ell(t)$

LIST OF SYMBOLS (continued)

$R_x(\tau)$	Auto correlation of $x(t)$
$S_{\bar{e}}(f)$	Average of $S_{e_1}(f)$ and $S_{e_2}(f)$
$S_p(\omega)$	Spectral density of random concentrated force
S_{po}	Constant value of spectral density for White Noise
$S_u(f)$	Autopower spectral density of $u(t)$
$S_{uv}(f)$	Cross power spectral density of $u(t)$ and $v(t)$
$S_x(\omega), S_y(\omega)$	Spectral densities of $x(t)$ and $y(t)$, respectively
$S_{w_r}, S_{\theta_\theta}, S_{\theta_z}$	Spectral density of $w_r, \theta_\theta, \theta_z$, respectively
$S_{NR}(f)$	Narrow-band signal-to-noise ratio for a quantity obtained by differencing
t	Time (Section IV)
t	Shell thickness (Appendix A)
T	Averaging time (Section 4.3.2) or sampling interval (Section 4.6.2)
$T_A(t)$	Kinetic energy of structure A
$u_i(\xi, \eta)$	Displacements
$\hat{u}_{k\ell}$	Finite element degrees of freedom
$U_A(f)$	Spectral density of $U_A(t)$
$U_A(t)$	Potential energy of body A
$v_a(t), v_b(t)$	Velocities at coupling points a and b of bodies A and B
$v_b^{(\ell_c)}(t)$	Velocity of coupling point b of body B when B is acted on only by the coupling force ℓ_c
$V_a(f), V_b(f)$	Fourier transforms of $v_a(t)$ and $v_b(t)$
w	Transverse shell displacement
$w(x, y)$	Out-of-plane displacement of a plate

LIST OF SYMBOLS (continued)

w_r	Displacement of panel normal to surface
x	Independent coordinate along undeformed beam neutral axis
\tilde{x}	Vector of displacements
x, y	Rectangular coordinates (in Section 4.4.1)
$x(t)$	Vertical displacement of CG of rigid bar
$x_a(t), x_b(t)$	Displacements at coupling points a and b of bodies A and B
x_1, x_2	Vertical displacement of ends of rigid bar
$\overline{x^2}$	Temporal mean of $[x(t)]^2$
$X_a(f), X_b(f)$	Fourier transform of $x_a(t)$ and $x_b(t)$
\tilde{x}_g, \tilde{y}_g	Unit vectors at the center point of a Semi-Loof shell element
$y(t)=L\theta(t)$	Relative displacement of one end of rigid bar with respect to opposite end
$y(x, t)$	Lateral displacement of beam longitudinal axis
$y_i(t)$	Displacement of channel i transducer
$\overline{y^2}$	Temporal mean of $[y(t)]^2$
$Y_i(f)$	Fourier transform of $y_i(t)$
Z_a	Mechanical impedance of body A at a coupling degree of freedom a
$\alpha(f)$	Spectral discrete difference function defined by Eq. (6.26)
α_i	Gauss points
α_1, α_2	Ratios L_1/L and L_2/L such that $\alpha_1 + \alpha_2 = 1$
β	Ratio of c_2/c_1 for rigid bar system, damping factor for beam material

LIST OF SYMBOLS (continued)

$\tilde{\beta}_1, \tilde{\beta}_2$	Modal damping factors
γ	Ratio of k_2/k_1
$\tilde{\gamma}$	Transverse shear strain vector resultant
$\gamma_{uv}^2(f)$	Coherence between $u(t)$ and $v(t)$
$\tilde{\gamma}_{XZ}$	Transverse shear strain on the boundary of an element
$r^2(f)$	Real function of frequency defined by Eq. (6.49)
$\delta(f)$	Complex channel mismatch function defined by Eq. (6.23)
$\delta(\xi)$	Discontinuity along an element edge $-1 \leq \xi \leq 1$
Δf	Width of a frequency band used to define the $\langle \rangle_t$ operator
Δx	Separation distance for transducers
$\epsilon_1, \epsilon_2, \epsilon_{12}$	In-plane shell strains
ζ_r	Viscous damping ratio of r^{th} mode
η	Ratio of distance to $P(t)$ from CG of rigid bar/length L
$\eta_i(t), \eta^{(i)}(t)$	i^{th} generalized coordinate
η_A	Internal loss factor for body A
η_{AB}	Coupling loss factor for power transmission from body A to body B
$\theta(t)$	Angular displacement of rigid bar (Section V)
$\theta(t)$	Rotation angle to be measured (Section VI)
$\theta(x, t)$	Angular displacement of beam longitudinal axis (Section V)
$\theta(x, y, t)$	∇w , gradient of $w(x, y, t)$ (Section IV)
θ_z	Angular displacement of panel about z-axis

LIST OF SYMBOLS (continued)

$\hat{\theta}_A$	Estimate of $\ddot{\theta}$ obtained with an actual transducing system
$\hat{\theta}$	Estimate of $\ddot{\theta}$ obtained with an ideal transducing system
θ_θ	Angular displacement of panel about θ -axis
κ	Radius of gyration for the cross-section of a beam or plate
λ	Wavelength of flexural wave
λ_n	Roots of characteristic frequency equation for FREE-FREE beam
Λ	Proportionality constant between spring damping and spring stiffnesses
ν	Poisson's ratio
ξ, η	Curvilinear shell coordinates
$\xi^{(i)}$	Viscous damping ratio of the i^{th} mode
π	Energy functional (Section III)
π	3.1416 (Sections IV and VI)
$\pi^{(i)}(t)$	Power input to the i^{th} mode
$\pi_{A, \text{DIS}}$	Power dissipated internally by body A
$\pi_{A, \text{IN}}$	Power input to body A from prescribed sources
$\pi_{AB, \text{TRAN}}$	Power transmitted from body A to body B
$\Pi_{AB}(f)$	Spectral density of power transmitted from A to B
ρ	Density of beam material
$\sigma_{e_i}^2$	Mean square signal in channel i
σ_n^2	Mean square noise
σ_u^2	Mean square value of the zero-mean variable $u(t)$
τ	Lag time in the correlation function

LIST OF SYMBOLS (continued)

$\phi(\xi, \eta)$	Any function to be integrated over an element domain
Φ_n, Φ_k	Eigenfunctions for FREE-FREE beam
$\chi_1, \chi_2, \chi_{12}$	Shell curvatures
$\psi_i(x, y)$	i^{th} normal mode shape
$\psi_o^{(r)}, \psi_i^{(r)}$	Amplitude of r^{th} normal mode shape at points o and i
ω	Frequency in radians/unit time
$\omega_i, \omega^{(i)}$	i^{th} natural frequency (Section IV)
ω_n, ω_k	Natural frequencies of n and k mode shapes for beam (Section V)
ω_r	$2\pi f_r$
Ω	Weighting factor for the centerpoint of Gaussian integration
$\mathcal{L}(u_i)$	Differential operator governing the interior of a structure
$\mathcal{B}(u_i)$	Differential operator governing the boundary of a structure
$\langle \rangle_t$	Expectation with respect to time, usually of the portion of a positive definite quantity within a specified frequency band
$\mathcal{F}(\)$	Fourier transform
$\mathcal{F}^{-1}(\)$	Inverse Fourier transform
$\text{Re}(\)$	Real part of
$\text{Im}(\)$	Imaginary part of
$ \ $	Modulus of
$\overline{}$ (overbar)	Smoothing with respect to frequency

LIST OF SYMBOLS (Concluded)

$\langle \rangle$	Ensemble averaging (when applied to products of Fourier transforms)
∇	Gradient operator in two dimensions
$E[\]$	Expectation
$(\)_{\text{rms}}$	Square root of the temporal mean of $(\)^2$
$(\)^{\cdot}$	Time derivative of $(\)$
$(\)'$	Spatial derivative
\wedge	Denotes estimate
$*$	Convolution operator
$*$	Complex conjugate (used as superscript)
$M(\cdot)$	Modulation function $\sin(\cdot)/(\cdot)$

SECTION I

INTRODUCTION

With the design and development of inertial sensing systems and laser beam control systems, angular vibration measurements and predictions have become as important, and in some cases more important, than translational vibration measurements and predictions. Also, very small amplitudes of angular vibration, on the order of a few microradians, have become important, especially in the design of airborne laser beam control systems. Angular vibration is the major factor in beam jitter of laser systems. Beam jitter is dynamic misalignment of a laser beam due to dynamic motion of the components of the optical train through which it passes. Beam jitter is normally random in time and is specified by the root mean square value. The optical train consists of both stationary and servo-controlled mirrors and sensors from beam initiation in the laser device through the pointing and tracking system.

Beam jitter is dependent on a number of factors. Among these are mechanical resonances of individual optical components and their supporting structures, dynamic behavior of the entire structure system, spacing of the optical components, mounting system characteristics, dynamic characteristics of the aircraft, and frequency content of loads. Since the laser or electro-optical system is mounted to the aircraft, and the aircraft is subjected to many loads, such as gusts, turbulence, etc., the dynamics of the aircraft must be predicted accurately. The dynamic characteristics of the aircraft that are used for design and analysis consist of both translational and angular vibrations. The elastic modes of an aircraft may be as low as 1 or 2 Hz. The lowest elastic modes of individual components of a laser system may be in the 150 to 300 Hz range. Therefore,

to perform complete frequency response analysis may require modal information from 1 Hz to 1000 to 2000 Hz.

The objectives of this contract were (1) to develop techniques for predicting the low and high frequency angular environment of aircraft; (2) to develop accurate angular vibration measurement techniques; (3) to develop techniques for predicting the low and high frequency angular vibration of combined airframe and electro-optical systems; and (4) to demonstrate these techniques by applying them to an aircraft-like structure and comparing the results with measured data. For this report, low frequencies are defined as those frequencies for which individual normal mode shapes can be predicted accurately or determined by tests.

The work on this contract was subdivided into three phases. The objective and accomplishment of Phase I was to identify methods which could be used to predict angular vibration for both low and high frequencies. Literature searches were performed to obtain information on past experience in the area of angular vibration. Also searched were methods and improvements to methods which may be useful in the prediction of angular vibration at low and high frequencies. For low frequencies, it was decided the main thrust should be directed towards obtaining more accuracy per degree of freedom in finite element analysis. This approach satisfies all of the objectives for the low frequency method and also gives a method which is flexible and adaptable for application to complex structures. The Semi-Loof shell element and its companion beam element were selected as the most promising candidates to meet this objective. The technique selected for combining components or structures, such as an aircraft and an electro-optical system, into a system analysis for low frequencies was component mode synthesis. The component mode synthesis technique treats each component or structure in terms of its modal description (obtained either from tests or

analysis). For high frequency analysis, statistical energy analysis (SEA) was selected. SEA treats the structure in terms of some averaged, or statistical, description. One of the most important features of the SEA method is its ability for making use of whatever level of detail is available in the description of a structure.

The major emphasis of Phase II was the detailed development of the methods chosen during Phase I. For the low frequency method, the approach taken was to implement Semi-Loof into the finite element code, NASTRAN. As originally derived, the Semi-Loof element was for use on static problems only. Therefore, the mass matrices for the elements had to be incorporated. To be able to model complex structures such as airframes, other enhancements to the element had to be made. Among these were orthotropic material properties, offset beams, smeared stiffeners, variable thicknesses for the shell elements, non-prismatic beam geometry, and distributed loads. The input was developed using NASTRAN bulk data card formats. A stress recovery and a rotational recovery capability were added. The element is incorporated into NASTRAN by means of pre- and post-processors and NASTRAN DMAP instructions. The processors incorporate a matrix assembly routine, and an extensive error checking capability. A User's Manual for Semi-Loof was written and is included in this report as Appendix A. During Phase II a number of sample problems were executed and the results were compared to NASTRAN solutions and tests. The results showed good improvement in accuracy per degree of freedom.

As stated earlier, statistical energy analysis (SEA) had been chosen as the best approach to predict the angular vibration environment of aircraft when only an averaged or statistical description of the structures is known. SEA has been used as a tool for acoustical analysis in the past but no work had been done toward employing this method to predict angular vibration environments. The theory of SEA was studied and the essential features of SEA which make it attractive for the

present purpose were identified. A formula for estimating coupling loss factor by the wave transmission method was rederived without assuming response to be in the form of traveling waves. An experiment was conducted where the wave transmission method was used to predict transmitted power and the equilibrium energy ratio between two coupled plates. This experiment demonstrated that SEA could be used to predict the energy ratio and power transfer coefficient between coupled plates in a high frequency region where finite element modeling was impractical. A relation was derived between the r.m.s. angle and the vibrational energy in frequency bands. An experiment was conducted to test this relation and show how r.m.s. angular displacement can be obtained from SEA results without detailed knowledge of individual mode shapes or natural frequencies. Software was developed for the minicomputer-based equipment used to perform these experiments.

A study was also conducted on relationships between linear and angular vibration for various structural components. A detailed study of the simplest structural system possessing both linear and angular degrees of freedom was conducted. The ratio of the mean square angular to mean square linear displacement was investigated for simply-supported beams, a simply-supported flat plate, a free-free beam, and a curved stiffened panel. A simple relation between mean square angular to mean square linear displacement was derived using wave theory.

In order to verify the prediction methods developed, reliable experimental data had to be obtained, especially dynamic rotations at specific points. The measurement method used to obtain these rotations was differencing of translational acceleration signals. A better quantitative understanding of limitations and error sources was desirable. Therefore, theoretical derivations were developed for estimating errors introduced by noise in individual channels, frequency-dependent gain and phase mismatching between channels, and flexure of the mounting surface. Effects of the first two error sources were

demonstrated by experiment and it was demonstrated that mismatch error can be reduced by appropriate data processing. Also, an expression was derived for coherence of a measured angular frequency response when angular response was obtained by differencing.

The major objective of Phase III was to apply the methods developed to a complex structure. A fuselage section of a fighter aircraft was chosen as the test structure. This structure was chosen because it was a fairly complex structure to analyze. A finite element model of the fuselage was developed using the Semi-Loof elements. All of the features that were developed for this element were employed in the modeling. An eigenvalue analysis of the structure was performed to determine the normal modes. A frequency response analysis was performed with random noise input at three different points. The responses at a number of translational and rotational degrees of freedom were output. A test of the fuselage was also performed. The fuselage was supported on a low stiffness mounting system and driven by a small shaker system at three different locations. Both force to linear acceleration and force to angular acceleration transfer functions were measured for drive and response points corresponding to points in the Semi-Loof model. After processing the data on the minicomputer-based modal analysis equipment, the experimental data was compared to the results predicted by Semi-Loof.

Section II of this report described the literature searches performed to obtain information on past experience in the area of angular vibration. Also searched were methods and improvements to methods which might have been useful in the prediction of angular vibration at low and high frequencies. This section also discusses some of the methods looked at and discarded, and the reasons behind the selection of the methods chosen for further study.

Section III describes the work performed on the Semi-Loof elements, which was the main low frequency method. This section gives a review of the theory of finite elements in general and the Semi-Loof element in particular. The implementation of Semi-Loof into NASTRAN is described in detail. Comparisons and evaluations are made using small simple problems, results of analysis using other elements, and test results.

Section IV describes the high frequency method, which is statistical energy analysis (SEA). A brief summary of the theory of SEA is presented. Derivations of relations and experiments are described which predict transmitted power, the equilibrium energy ratio, and angular response between two coupled plates.

Section V is a study of the relationship between linear and angular vibration for various structural components.

Section VI describes the work performed on some quantitative methods of estimating the reliability of measured power spectral density functions or frequency response functions for the case where the response quantity is an angular motion obtained by differencing of signals from linear transducers.

Section VII describes the fuselage finite element model and the dynamic test performed on this fuselage. A comparison of the results of the test and analysis is given.

Section VIII gives a summary of the work, draws conclusions from the progress made, and briefly describes additional work which should be done in the future.

SECTION II

BACKGROUND

2.1 LITERATURE SEARCH

The first task started in Phase I was a literature search to determine what had been done in the past for predicting angular vibration as well as high frequency translational vibration. As usual in an extensive literature search, one source led to another, often via topics which were not directly related to the immediate problem. Some of the most useful information and sources were obtained through personal communication.

Initial searching revealed little on angular vibration as such which was not previously known. However, numerous possibilities were identified within the overall field of vibration analysis which might lead to improved methods for prediction of angular vibration.

The data bases searched included DDC Technical Reports, NTIS, NASA, and COMPENDEX. The DDC (Defense Documentation Center) Technical Reports data bank contains more than 1,200,000 records dating back to March, 1953. DDC receives these reports from defense facilities and their contractors who are required to submit to DDC copies of each report that records scientific and technical data from defense sponsored research, development test, and evaluation. NTIS (National Technical Information Service) contains the complete Government Reports Announcements (GRA) file from the National Technical Information Service. It contains 641,000 citations of government research from over 240 agencies. The file, which dates back to 1964, is updated every two weeks, and is growing at the rate of 78,000 abstracts per year. The NASA data base contains the results of worldwide research and development activities in aeronautics, space, and supporting disciplines. It now contains nearly a million

documents which are abstracted and indexed. COMPENDEX is a machine readable version of the Engineering Index data base which provides abstracts and an index to the world's significant engineering literature and conference proceedings covering the time span of 1970 to the present. This data base covers 3,500 journals, publications, and papers from the proceedings of conferences as well as selected government reports and books. It contains over 90,000 abstracts and 655,000 citations.

The major topics that were searched included (1) linear vibration of aircraft, (2) angular vibrations, (3) finite elements or finite differences in combination with angular vibration, high frequencies, and statistics, (4) high frequency methods, (5) acoustical methods, and (6) statistical energy methods. These major topics led to many minor topics which were investigated to the extent of determining the usefulness in predicting angular vibration.

Automated searching on the key phrases "finite element" and "angular vibration" was unproductive. "Angular vibration" often turns out to mean torsional vibration of shafts. There simply does not seem to have been any finite element or other numerical work that focused specifically on angular vibrations as that term is understood in the current effort. Consequently, the searching strategy for the analytical portion of this effort was shifted rather early toward developments that might contribute indirectly to angular vibration predictions. This broadened the possibilities immensely (for example, a DDC search listed some 1300 references under "finite element"). Narrowing the search to shell elements, searching continued mostly in a manual mode, with many papers leading to others through references. Most of these were concerned with isoparametric elements which have received the most attention from researchers. A number of papers in this vein were reviewed. However, the Semi-Loof element, which was adopted for study and implementa-

tion, was obtained through a personal communication. Subsequently, other papers on Semi-Loof were discovered in the literature.

As noted previously in this section, literature searching under key words such as "angular vibration" yielded almost no research or development work not previously known. The few references uncovered on angular vibration of any kind were concerned with specific pieces of hardware rather than prediction methods. For high frequency angular vibration, the search strategy was to investigate high frequency and statistical methods in general without reference to any particular type of response variable. A list of key words and phrases used in the computer-aided searches is given in Table 1.

TABLE 1
KEY WORDS USED FOR HIGH-FREQUENCY
VIBRATION LITERATURE SEARCH

angular vibration
high frequency vibration
high frequency response
blast response
(statistical) and (vibrations)
statistical energy methods
high frequency approximations
mode slopes
mode rotations
(regression analysis) and (vibrations)
spatial averaging
spatial covariance
mode [modal] averaging
mode [modal] covariance
(statistical) and (structures)

spatial [spatially] random processes
(covariance) and (vibration[s])
(autocovariance) and (vibration[s])
Sommerfeld-Watson transformation
Poisson summation

It quickly became apparent during the literature search that a well developed body of theory called Statistical Energy Analysis (SEA) already existed and was at least partially applicable to the current problem. It also became clear that a survey of this field was not necessary because an excellent survey and tutorial report had been prepared for AFFDL in 1974 by R. H. Lyon of M.I.T. [1]*. This report was extremely valuable in acquiring a basic understanding of the underlying theory. An extensive bibliography organized by subject within SEA is contained in Reference [1].

2.2 SELECTION OF LOW FREQUENCY METHODS

The finite element method has risen to prominence in aircraft structural analysis in parallel with advancing computer capabilities. Its primary appeal lies in its generality and relative ease of use. Literature searching revealed thousands of titles containing the term "finite element" of which a large fraction would no doubt have some bearing on the problems at hand. Thus, there was never any question that this method would be the analysis tool used in the low frequency end of the angular vibration spectrum, with perhaps some exceptions for special cases or crude approximations.

With finite elements so widespread, there has been a gradual shift of emphasis away from purely theoretical developments toward questions like cost effectiveness and adaptability to well-established codes. In the present work, the goal was

* Numbers in brackets designate References at end of report.

to bring some advanced finite element technology from a "laboratory" status to a production environment. In practice, this means testing, evaluating, and "idiot-proofing" an element, and above all, tying it to an existing software system so as to take advantage of the large investment in auxiliary functions that such systems have.

While a number of theoretically appealing elements were reviewed, many of them were rejected mainly because they could not be used with NASTRAN. The goal of NASTRAN compatibility was adopted as it became clear that any other approach to assembling finite element software would be prohibitive in terms of the effort required in coding, checking, training users, and gaining their acceptance.

Implicit in this approach is the proposition that angular problems are not fundamentally distinguished from those of translational vibrations. There is no way to divorce an angular deformation from the associated translational deformation. Mathematically, one is the derivative of the other. Hence, the pursuit of better angular vibration was embodied in a search for better finite element methods in general. In other words, an evolutionary approach was taken rather than revolutionary.

The authors began this work with considerable experience in application of finite element techniques to aircraft structural analysis, much of it using NASTRAN. Following is a summary of the procedures and rules of thumb evolved by the authors in the course of this work.

First, the choice of mesh depends on a number of factors. The areas where response is more of interest need a finer mesh. Also, angular vibrations generally require a finer mesh than translational vibrations, and if the anticipated rotation is primarily about one axis, refinement perpendicular to that axis may be in order.

Boundary conditions must be handled carefully, often using auxiliary coordinate systems, in order that the allowable

motions in the model reflect the actual situation as nearly as possible.

Eigensolution strategies must be chosen carefully. When relatively high frequencies are desired, a sweeping procedure such as the Givens method is in order. In this case, a condensation step is usually called for. This may take the form of Guyan reduction, or generalized "dynamic reduction," both available in NASTRAN.

The debugging stage is just as important as the modeling stage. Following are some of the checks that Anamet personnel routinely use in verifying a finite element model:

- (1) Aspect ratios and degrees of skew are checked for plane and solid elements using a preprocessor developed by Anamet.
- (2) Plots of the undeformed structure are generated. Anamet's preprocessors NASSET and SAPLOT are sometimes used to generate partial views.
- (3) The weight calculated from the finite element model is compared with the actual weight of the structure.
- (4) Diagonals of the master stiffness and mass matrix are checked if a singularity or near-singularity is suspected.
- (5) Simple static cases are run, usually with dead weight acting, before dynamics is attempted.
- (6) For every NASTRAN matrix decomposition, the maximum ratio of matrix diagonal to factor diagonal is noted, as a check on the conditioning of the matrix.
- (7) The quantity ϵ_o is noted. For static runs where an equation of the form $Ku = \rho$ is solved, a residual vector $\delta P = \rho - K^{-1}u$ is calculated. This vector, which would be zero if no arithmetic round-off were present, has the form of a load vector. Its magnitude is assessed by comparing the work done by δP to that done by the actual load ρ , i.e. $\epsilon_o = \delta P^T u / \rho u$. For dynamic runs the rigid body matrix X should be singular. Its norm is calculated and compared to the stiffness matrix for the rigid-body coordinates, i.e. $\epsilon_o = ||X|| / ||K_{rr}||$. In both cases, for a successful

solution, ϵ_o should be near the order of computer precision, i.e. 10^{-10} to 10^{-14} .

- (8) When eigenvalues are extracted, the off-diagonal modal mass term is noted. This quantity indicates how nearly orthogonal the eigenvectors are.
- (9) When rigid-body modes are anticipated, the corresponding eigenvectors are checked to be sure that they are actually pure translations and/or rotations.
- (10) The first few elastic modes are checked for reasonableness. Frequencies can often be checked against estimates made by hand. Mode shapes are checked with plots.
- (11) Reaction forces are printed to insure that a constraint introduced for the purpose of eliminating a singularity has not inadvertently caused a spurious reaction force.
- (12) Strain energies are sometimes checked. A high concentration of strain energy in one or a few elements sometimes indicates a modeling error.

In angular vibration applications of finite elements it is important to assess the quality of the mode shapes that have been computed. Ideally one would like to be able to divide mode shapes into three groups. First, "good" modes; second, modes that individually are questionable, but as an aggregate possess reliable statistical characteristics; and third, unreliable modes. One rough rule that can be applied is to assume that only a certain percentage of the analysis set represents good modes. A clue to the unreliable modes (at high frequency) is a sudden fall off in modal density. Also, one should bear in mind the importance of mode shapes with respect to the response of interest. This may be assessed by requesting mode normalization such that each mode has a value of one at a specified degree of freedom. The modal mass would then be an indication of the importance a particular mode plays with respect to a particular degree of freedom. This assessment could be made for a number of degrees of freedom, including loaded d.o.f. as well as response d.o.f.

Very large finite element models, or models which combine finite element models with test data, may be executed with the technique of component modal synthesis. In component modal synthesis, each component is characterized in terms of a number of natural frequencies and associated mode shapes, modal masses, modal stiffnesses, and modal damping. This can be done in three different ways:

- (1) Perform a test using either a shaker or impulse loading. Record the modal information listed above.
- (2) Perform a computer analysis using a finite element or other appropriate program.
- (3) Analyze the component by hand.

There are two major advantages to describing a component in modal formulation. First, the equations governing the dynamic behavior of the component become uncoupled. Second, in the modal formulation it is possible to obtain adequate accuracy by using only a subset of the total modes of the component.

The components are then combined into a system analysis by writing equations of constraint between the interconnecting degrees of freedom of each component. These equations of constraint are the equations transforming the physical degrees of freedom to modal degrees of freedom. For each equation for the physical degree of freedom, the coefficients of the modal degrees of freedom are the components of the eigenvectors associated with that physical degree of freedom.

The accuracy of the component modal synthesis technique is dependent on several factors. First, the individual eigenvectors of the components must be accurately calculated, especially at the connection degrees of freedom (both translational and rotational). To model a system up to a given frequency, (i.e., N Hz.) each component's modal description must contain information greater than this frequency. A good rule of thumb is to pass to the system analysis modes for each

component up to one and a half times the highest frequency of interest. However, this is very dependent on the problem and the number of connection degrees of freedom. The greater the number of connection points, the greater the number of modes that must be passed. Not including the higher modes in the system analysis leads to a formulation which is too stiff (i.e., the predicted frequencies will be higher than the actual).

Using the procedures and techniques described, it is seen that large models may be solved, and the accuracy of the solution is realistic when compared with the elements used. Therefore, the kinds of advances that may bear on angular vibration problems include more accuracy per degree of freedom, judicious approximations that improve cost effectiveness, better numerical solution methods, and better ways of recovering dependent response quantities (in this case, angular deformations). It was decided that more accuracy per degree of freedom would be the most promising area, and that led to a selection of a new element described in Section III.

2.3 SELECTION OF APPROACH FOR HIGH FREQUENCY PREDICTION

The selection of an approach for prediction of so-called high frequency angular vibration was conditioned by the definition of low vs. high frequency as stated in the previous section. Modeling of a small stiff structure with a fundamental mode at 500 Hz. is not necessarily more difficult than modeling a larger structure with its first resonance at 5 Hz. The difficulty occurs when a model must be capable of predicting response to inputs over a frequency range which contains a very large number of modes. Angular responses are particularly difficult to predict because their modal series representations tend to converge more slowly than do those for translational responses and thus more modes must be known. Higher order modes will be sensitive to structural details which are too small to model economically and may not even be identical for structures built

from the same design. In effect, for high frequencies, one does not have a fixed description of the structure even in physical coordinates. The distinction between low and high frequencies for the purposes of this work is thus a functional one. Low frequency analysis implies that properties of individual vibration modes can be obtained, either by analysis or test. High frequency analysis presupposes that this level of detailed knowledge is unavailable.

This definition does not rule out the possibility of a purely empirical approach to high frequency prediction. In fact, this has been the basis for much high frequency translational vibration work in the past. One could collect the available data and attempt to correlate vibration levels with flight conditions, aircraft type, and some general structural description. It was decided in Phase I that this approach was not appropriate for the current contract. It would duplicate work being done already at AFFDL.

Based on the above considerations, it was decided during Phase I that the method of Statistical Energy Analysis (SEA) was the most promising candidate. While much of the specific SEA theory was unfamiliar to the investigators at this point, a number of attractive features were clear:

- (1) The method does not necessarily require information about individual normal modes of a structure in order to make response predictions. The lack of such information inevitably introduces some uncertainty into the predictions but this may be acceptable for many cases. The point is that averaged descriptor quantities such as approximate model density and total mass may be sufficient to make useful first estimates of response.
- (2) High model densities may actually be an advantage. Each mode contributing to response acts something like a statistical degree of freedom. As more modes contribute, their variability (i.e. the uncertainty as to the properties of any single mode) tends to become less important. For example, early SEA work was often associated with room acoustics where mode counts in the audio band may be in the hundreds of thousands.

- (3) It appeared that deterministic methods of analysis with which the investigators were intimately familiar could be used to good effect in SEA modeling. In particular, certain aspects of large scale finite element analysis and minicomputer-based experimental modal analysis appeared promising. It was suspected that a higher level of structural detail could be incorporated into an SEA model by utilizing these technologies which were not available when most of the basic SEA theory was worked out.
- (4) SEA modeling can be attempted using structural descriptions of varying detail. This was considered essential if a method were to be usable during both preliminary and prototype stages of design.

In hindsight, the investigators are convinced that the decision to focus on SEA for high frequency predictions was correct.

SECTION III

LOW FREQUENCY METHODS

Dynamic problems in general, and angular vibration in particular, require different approaches for different frequency ranges. We may classify these two approaches as deterministic and statistical. In a deterministic approach, the analyst models the entire structure mathematically in a way that strives to represent the behavior of any part of it under the most general loading. There are, of course, limits on the capability of any mathematical model in terms of the amount of detail that can be achieved in the predicted response. Ideally, the analyst knows that he can achieve greater accuracy as the mesh is refined, and can determine some error bounds for his model. In the frequency domain, these limitations can be expressed in terms of upper bounds on frequency. The lowest vibration modes almost always involve gross motions of the entire structure, and these modes are easy to predict analytically (for example, a simple beam model may be adequate to pick up a few modes of a complicated fuselage). With increasing frequency, mode shapes begin to have wavelengths at or below the mesh spacing, so that the model necessarily begins to break down. Aside from modeling errors, eigenvalue subroutines often have numerical trouble with the higher modes. Beyond the range where individual modes can be calculated accurately it may be possible to extract further information using aggregates of mode shapes.

However, the point is that a deterministic model must break down, and it is at this point (or hopefully before) that statistical approaches begin to make sense. The situation is analogous to the transition from microscopic to macroscopic thermodynamics where at some scale one must abandon any attempt to look at individual particles and start to look at averaged quantities.

Statistical methods will be taken up in Section IV. The rest of this section is devoted to a discussion of the finite element method in vibration problems, and to a new finite element that has been applied to angular vibration problems.

3.1 FINITE ELEMENT REVIEW

The finite element method is the dominant but not exclusive method in use today for numerical solutions to boundary value problems. Its popularity stems from its extreme flexibility with regard to the range of shapes and boundary conditions that may be handled. Before proceeding to a discussion of the Semi-Loof shell and beam elements, we review very briefly the mathematical foundations of the finite element method. Some understanding of these foundations is necessary in order to appreciate the development of the Semi-Loof elements. It may be fair to say that some users of finite elements need a better mathematical understanding of the method. Without this knowledge, users with a superficial view of finite elements as "building blocks" may misapply the method.

A finite element formulation may be derived beginning with an energy functional. For a two dimensional (ξ, η) problem, for example, one would have

$$\pi(u_i) = \iint F[u_i(\xi, \eta)] d\xi d\eta + \int G[u_i(\xi, \eta)] dS \quad (3.1)$$

where the u_i are independent variables (functions of ξ, η), F represents energy in the interior, and G the boundary of the region. At this point the problem could be one of thermodynamics, electrostatics, elasticity, or any other field where an energy principle governs. For structures problems, Eq. (3.1) most commonly is the potential energy, and the u_i are displacements.

In the traditional calculus of variations approach, a formal variation is carried out in terms of the variations of each independent function u_i . Setting the coefficients of

each δu_i to zero then yields the Euler equations, which are the differential equations and boundary conditions governing the problem.

$$\delta \pi = 0 = \sum_i \frac{\partial \pi}{\partial u_i} \delta u_i \quad (3.2)$$

$$\begin{aligned} \mathcal{L}_i(u_j) &= 0 \\ \mathcal{B}_i(u_j) &= 0 \end{aligned} \quad (3.3)$$

At this point the finite element approximation functions are introduced

$$u_i(\xi, \eta) \approx \sum_j N_{ij}(\xi, \eta) \hat{u}_{ij}, \quad i = 1, 2, \dots \quad (3.4)$$

wherein \hat{u}_{ij} are undetermined coefficients, or independent variables, and the N_{ij} are shape functions, usually polynomials in ξ and η . For a linear displacement formulation of a structures problem it is necessary that these functions satisfy displacement boundary conditions (prescribed boundary conditions) exactly while force (or "natural") boundary conditions are usually only approximated. We also note that the use of double subscripts is entirely arbitrary at this point.

Substitution of Eq. (3.4) into Eq. (3.3) then yields algebraic equations which may be written

$$\sum_{k, \ell} A_{ijkl} \hat{u}_{kl} + B_{ij} = 0, \quad i, j = 1, 2, \dots \quad (3.5)$$

which may then be solved numerically. Equation (3.5) is meant to represent a static, free vibration, transient dynamic, or frequency response problem, with time variation implicit in the A_{ijkl} and B_{ij} .

An alternative approach is possible. Instead of taking the variation first and then introducing approximation functions, one can reverse these steps. That is, introducing Eq. (3.4) into Eq. (3.1) renders it a function of the \hat{u}_{ij} ,

after the integration has been carried out. One may then set the partial derivatives of it with respect to each u_{ij} to zero to obtain the same algebraic Eq. (3.5). This alternative is attractive particularly when nonlinear terms are present, making it possible to apply minimization techniques directly to the potential energy. The two approaches are summarized in Figure 1.

The foregoing development is a description of the more general Rayleigh-Ritz method, and as yet nothing has been said about the essential feature of the finite element method. We now examine the nature of the approximation functions $N_i(\xi, \eta)$ and the independent variables (undetermined coefficients) that multiply them. Consider for illustration a plate bending problem. The plate is divided into finite elements by mesh lines. The shape functions are such that each function has a value only at one of the node points of the mesh, and in the interior of elements surrounding that node point (see Figure 2). The node point value is the independent variable \hat{u}_{ij} associated with that shape function. It is possible to consider the shape functions within only a single element, provided continuity requirements have been met. That is, two elements with a common edge must both predict the same displacement values at any point along that edge, given values for the corner-point degrees of freedom at either end of that edge. This being the case, it is possible to make models with elements of any size and shape, and to form matrices for the entire system by accumulating matrices generated for individual elements.

In recent years the "patch test" has come into favor as a means of evaluating a proposed finite element. It is based on the stipulation that any element must represent both rigid-body motion and constant-strain deformation in order to assure that a finer mesh yields a better solution. This is like saying that a series approximation needs to have at least the constant and first-order terms correct. In the patch test,

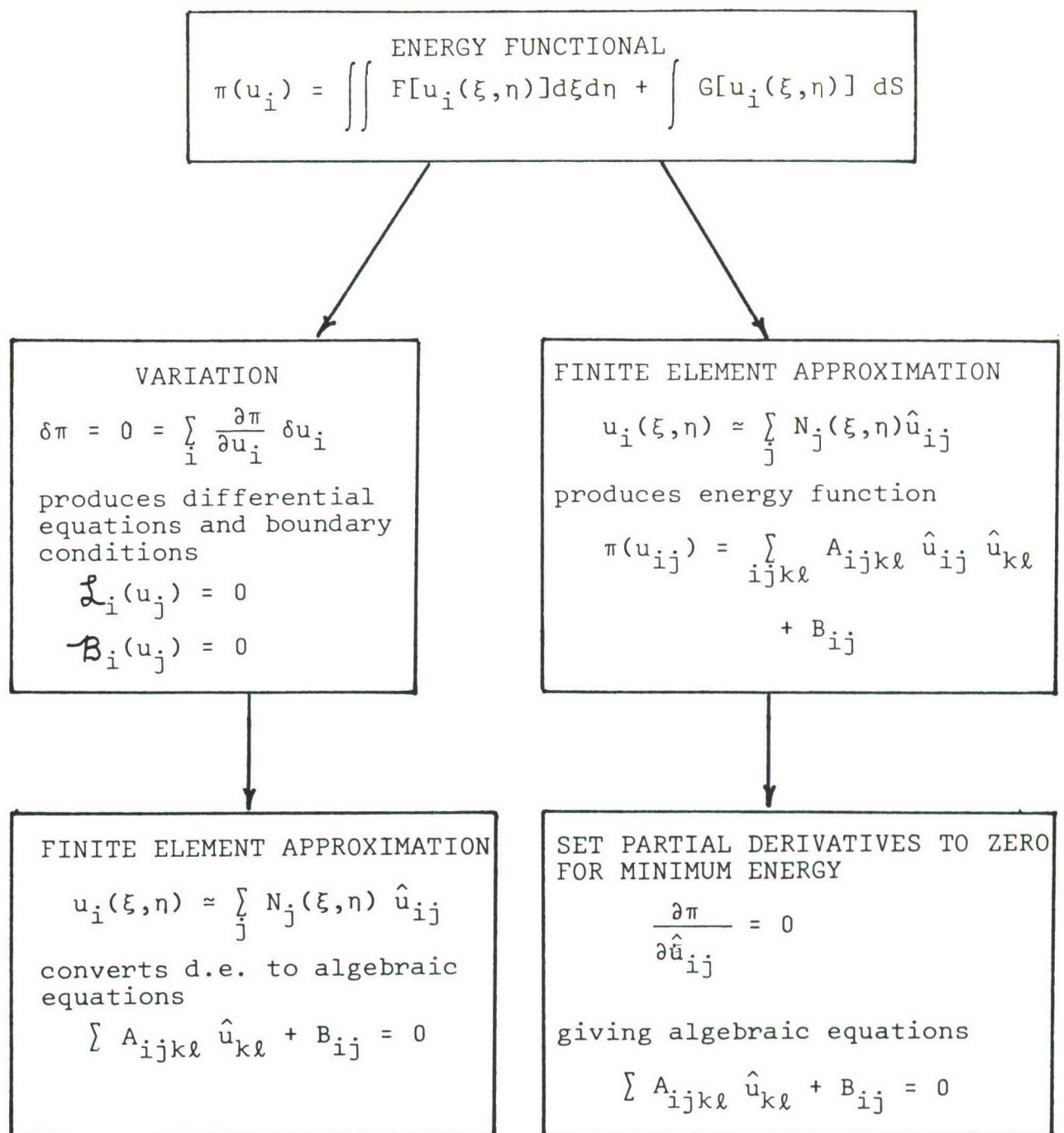


Figure 1 Finite Element Mathematical Derivation

these conditions are tested by taking a "patch" of elements and imposing rigid body or constant strain motions on the points on the boundary of the patch. The elements in the interior should then reproduce that state exactly.

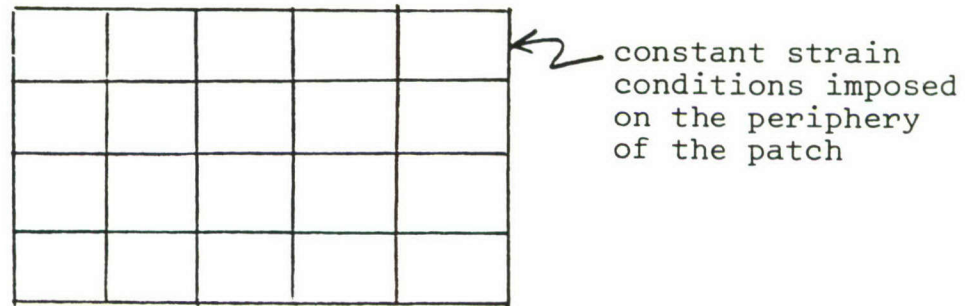


Figure 2 The Patch Test

3.2 THE SEMI-LOOF ELEMENTS

This section presents some of the background of the Semi-Loof elements that have been utilized in the present work. These elements were developed by Prof. Bruce Irons and his co-workers, and are best documented in Reference [2]. The developers are engineers and they are willing to venture into realms of approximations and adaptations where mathematicians might fear to tread, given the less than rigorous theoretical justifications. If past history is repeated, mathematicians will develop the missing theoretical basis for these and other advanced elements after engineers have been applying them successfully for some time.

First, the Semi-Loof shell elements are isoparametric elements, meaning that the functions used to describe the undeformed shape are the same as those used to describe displacements (i.e., the shape functions). This is a well known approach to shell elements, explained in texts such as Chapter 8 of Reference [3]. Also commonplace is the use of Gaussian quadrature for integration of element stiffness and mass. This is a form of numerical integration where the locations and weighting factors for sampling functions are chosen such as to minimize truncation errors when the integrands are polynomials. The locations are known as Gauss points. For two-point integration of a function on a normalized interval $(-1,1)$, for example, the Gauss points are located at $\pm 1/\sqrt{3}$ and the weighting factors are both unity.

The Semi-Loof elements begin to deviate from orthodox elements with the question of inter-element compatibility, or conformity as it is sometimes called. This concept was introduced in Section 3.1 and it asserts that displacements (and slopes, for bending elements) must be continuous across element boundaries as a condition of admissibility of the shape functions. The Semi-Loof elements are basically non-conforming

elements, but the non-conformity is of a special kind: such that any discontinuity and its first moment integrate to zero between elements. That is, if ξ is the normalized arc length along an element interface ($-1 \leq \xi \leq 1$), and $\delta(\xi)$ is a discontinuity, then

$$\int_{-1}^1 \delta(\xi) d\xi = 0$$

and

$$\int_{-1}^1 \xi \delta(\xi) d\xi = 0 \quad (3.6)$$

so that the elements can pass the patch test. This property is made possible by the use of Loof nodes which are located at the Gauss points along each element edge (i.e., $\xi = \pm 1/\sqrt{3}$). This idea stems from an early development by H. W. Loof [4] who proposed very high order elements providing exact solutions in element interiors, with shape functions collocated at a number of points along element interfaces.

A Semi-Loof isoparametric quadrilateral begins with 17 nodes and 43 degrees of freedom. The node points include four corner nodes, four midside nodes, and a center node (see Figure 3). The degrees of freedom are:

3 translations at each midside and corner node	3x8 = 24
Rotations about two axes at each Loof node and the corner node	2x(8+1) = 18
The "bubble function" $w = (1-\xi^2)(1-\eta^2)$ required by the patch test	= <u>1</u>
TOTAL	= 43

The 43 degrees are then reduced to 32 by imposition of the following constraints:

First, rotations about the normals at all eight Loof nodes are set to zero. These rotations introduce transverse shear strains, which are small for thin shell theory (8 constraints).

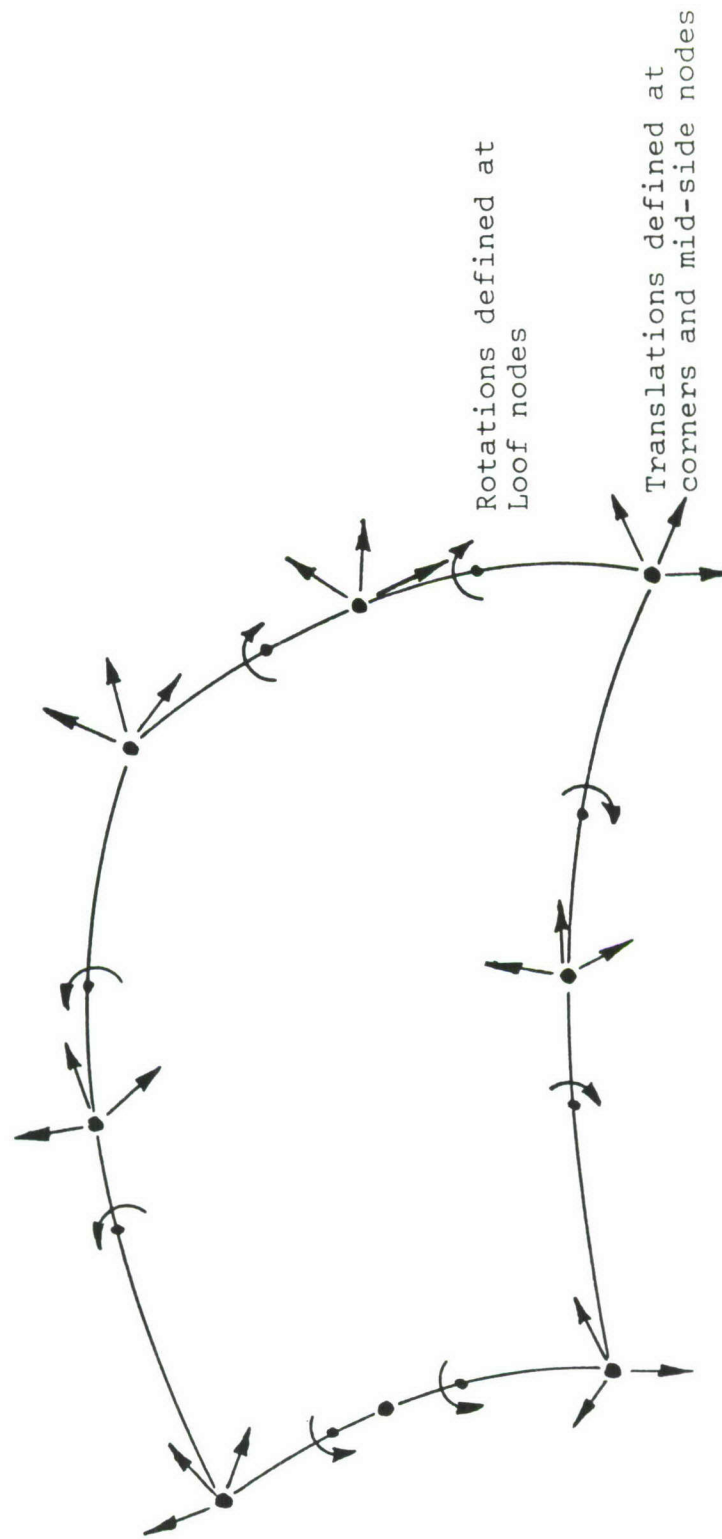


Figure 3 Semi-Loop Quadrilateral Element Degrees of Freedom

Second, two constraints are used to reduce out the center point rotations. These are:

$$\int \underline{X}_g \cdot \underline{\gamma} \, d(\text{area}) = \int \underline{Y}_g \cdot \underline{\gamma} \, d(\text{area}) = 0 \quad (3.7)$$

where \underline{X}_g and \underline{Y}_g are unit vectors at the center point in the $\xi = \text{constant}$ and $\eta = \text{constant}$ directions, and $\underline{\gamma}$ is the vector resultant of the two transverse shear strains. Irons found through experience that constraining this average shear strain produced better results than simply eliminating the two center point shear strains.

Finally, the constraint

$$\int \nabla \cdot \underline{\gamma} \, d(\text{area}) = 0 \quad (3.8)$$

is transformed into a boundary integral by Green's theorem

$$\int (\text{thickness}) \, \gamma_{XZ} \, d(\text{boundary}) = 0 \quad (3.9)$$

and used to reduce out the bubble function leaving $43-8-2-1=32$ degrees of freedom. These are shown in Figure 3.

Ideally, the stiffness matrix for the quadrilateral would have a rank of 26 (32 nodal variables minus 6 rigid body motions). But each integrating point can contribute, at most, six (the rank of the modulus matrix) for a total of 24 with 2x2 Gaussian quadrature. Thus, under these circumstances, the element contains two spurious mechanisms; that is, deformation states for which the predicted strain energy is zero. Experience has shown that for many cases this is not a problem, particularly when stresses are of primary interest. Nevertheless, this situation is unacceptable in a production environment. Two remedies are presently available, neither of which is ideal. The first is to introduce, somewhat arbitrarily, a fifth integration point at the center, with a small weighting factor, Ω ,

by

$$\int_{-1}^1 \int_{-1}^1 \phi(\xi, \eta) d\xi d\eta = \Omega \phi(0,0) + (1-\Omega/4) \sum_{l=1}^4 \phi\left(\pm \frac{1}{\sqrt{3}}, \pm \frac{1}{\sqrt{3}}\right) \quad (3.10)$$

The second alternative is to resort to 3 x 3 integration, for which the integration formula (taken from Table 8.1, Reference [3]) is

$$\int_{-1}^1 \int_{-1}^1 \phi(\xi, \eta) d\xi d\eta \approx \sum_{i,j=1}^3 H_i H_j \phi(\alpha_i) \phi(\alpha_j) \quad (3.11)$$

where the Gauss points are

$$\begin{aligned} \alpha_1 &= .7745966692 \\ \alpha_2 &= 0.0 \\ \alpha_3 &= -\alpha_1 \end{aligned} \quad (3.12)$$

and the weighting factors are

$$\begin{aligned} H_1 &= 5/9 \\ H_2 &= 8/9 \\ H_3 &= 5/9 \end{aligned} \quad (3.13)$$

In the NASTRAN implementation of Semi-Loof to be discussed subsequently, the choice of integration rules has been left to the user (2x2, 2x2+1, or 3x3).

Triangular elements have not yet been introduced yet. These elements are a fairly straightforward extension of the quadrilateral element. There are 24 degrees of freedom, including three displacements at each of three corner nodes and each of three midside nodes, plus six Loof rotations (Figure 4). There is no problem with spurious mechanisms, since three integration points provide a matrix rank of $3 \times 6 = 18$, plus six rigid body modes = 24, the size of the matrix.

The third Semi-Loof element is the curved beam, which is meant to join with the shell elements along their edges. This

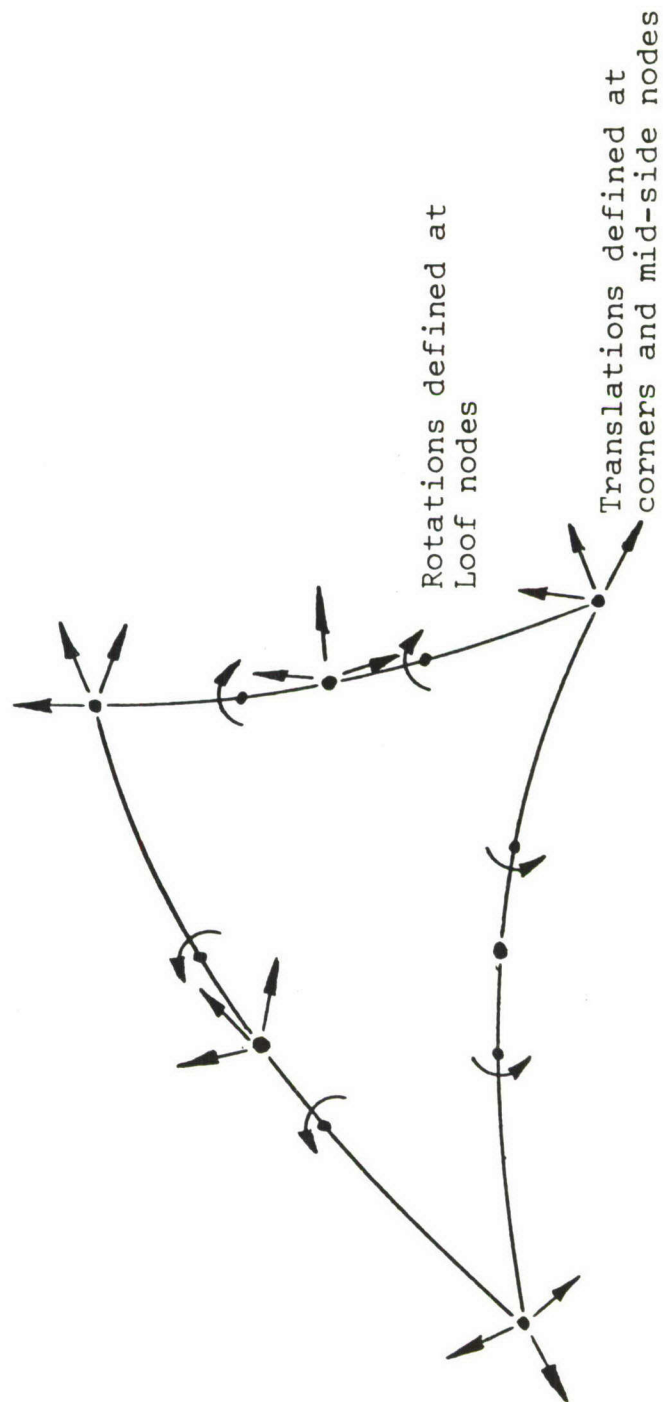


Figure 4 Semi-Loop Triangular Element Degrees of Freedom

element is documented in Reference [5]. Like the shell element, constraints are used to reduce out unwanted degrees of freedom. The element has two end nodes, a center node, and two Loof nodes, like the edge of a shell element. This is shown in Figure 5.

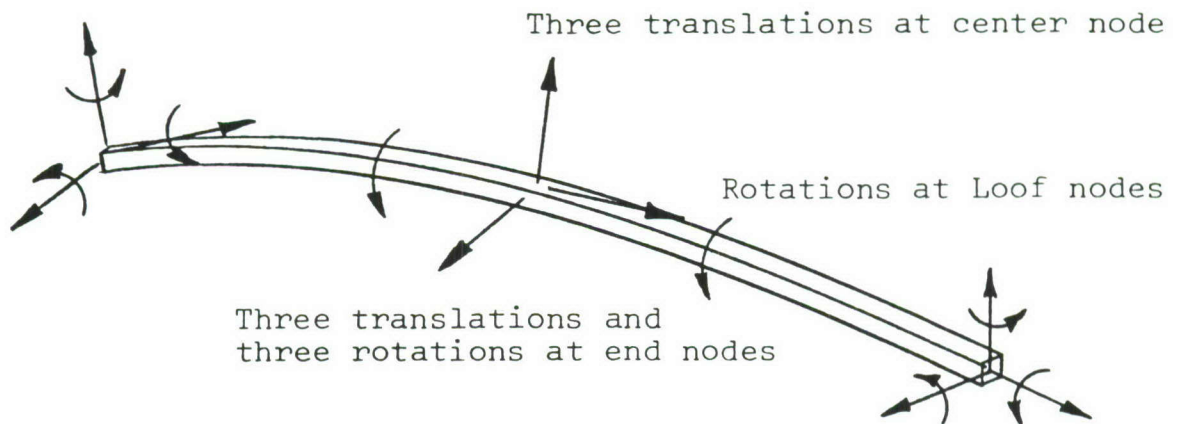


Figure 5 Semi-Loof Beam Element Degrees of Freedom

There are initially three displacements and three rotations at each end point, three displacements at the center point, and three rotations at each Loof node. Gauss point transverse shear strains are zeroed out, eliminating two unwanted rotation degrees of freedom at each Loof node. The final element has quadratic variation of both bending and torsion strains.

Prof. Irons has made available his shape function subroutines. When supplied with element geometry and isoparametric coordinates ξ, η , they return values of all the shape functions and their derivatives so that an engineer/programmer may change the integration method, or form different integrals without disturbing the basic shape function computations. These subroutines are the heart of the PRELOOF code, developed as part of this effort and described in Appendix A of this report.

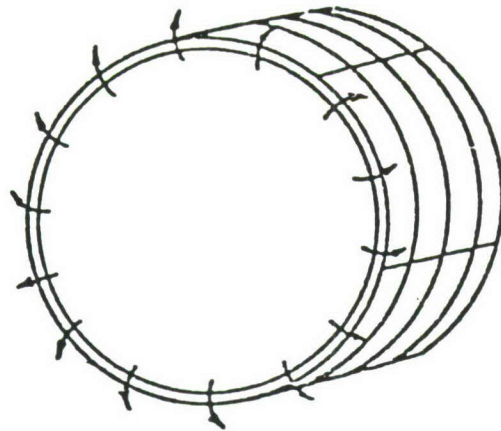
3.3 COMPARISON AND EVALUATION OF SEMI-LOOF ELEMENTS

A number of evaluations and comparisons of Semi-Loof elements have been published in the literature. At the time, Semi-Loof was in a "laboratory development" stage, and had not been applied much in production situations. Many of the elements that were compared to it were not in production status either. Additional evaluations were carried out as part of this effort with the specific aim of comparing Semi-Loof with the widely used plate elements in NASTRAN, since the final goal was to make Semi-Loof available to NASTRAN users.

We summarize very briefly some of the results reported in the literature. Irons [2] reports on a cantilever cylinder (Figure 6) with end moments. A very coarse mesh (60° elements) gave stresses accurate to within 1%. This is a rather spectacular result. Uniformly loaded square and circular plates also give excellent results. On the other hand, a point-loaded square plate is much less satisfying, though still acceptable (Figure 7). This may be a consequence of the fact that Semi-Loof corners can hinge. In practical situations, point loads are usually carried by beams, however.

Martins and Owen [6] compare Semi-Loof's performance with several other elements with respect to vibration of a square plate and of a cantilever cylindrical panel. Comparisons were presented as plots of relative error versus number of degrees of freedom employed, and each problem showed a substantial improvement for Semi-Loof over most other elements. The same authors [7] report favorable results for both elastoplastic and geometrically nonlinear problems. In this case, comparison with other elements is more difficult.

As an initial comparison of Semi-Loof with the standard QUAD2 (COSMIC NASTRAN) and QUAD4 (MSC/NASTRAN) plate elements, a simple plate bending problem was run. The results (Figure 8) were encouraging in that Semi-Loof outperformed both NASTRAN elements. A surprising result was that the QUAD4 did not converge any faster than QUAD2, although it did better for the



Radius = 10"
 Length from fixed end = 8"
 Thickness = 0.8"
 Poisson's Ratio = 0.3

Figure 6 Cantilever Cylinder

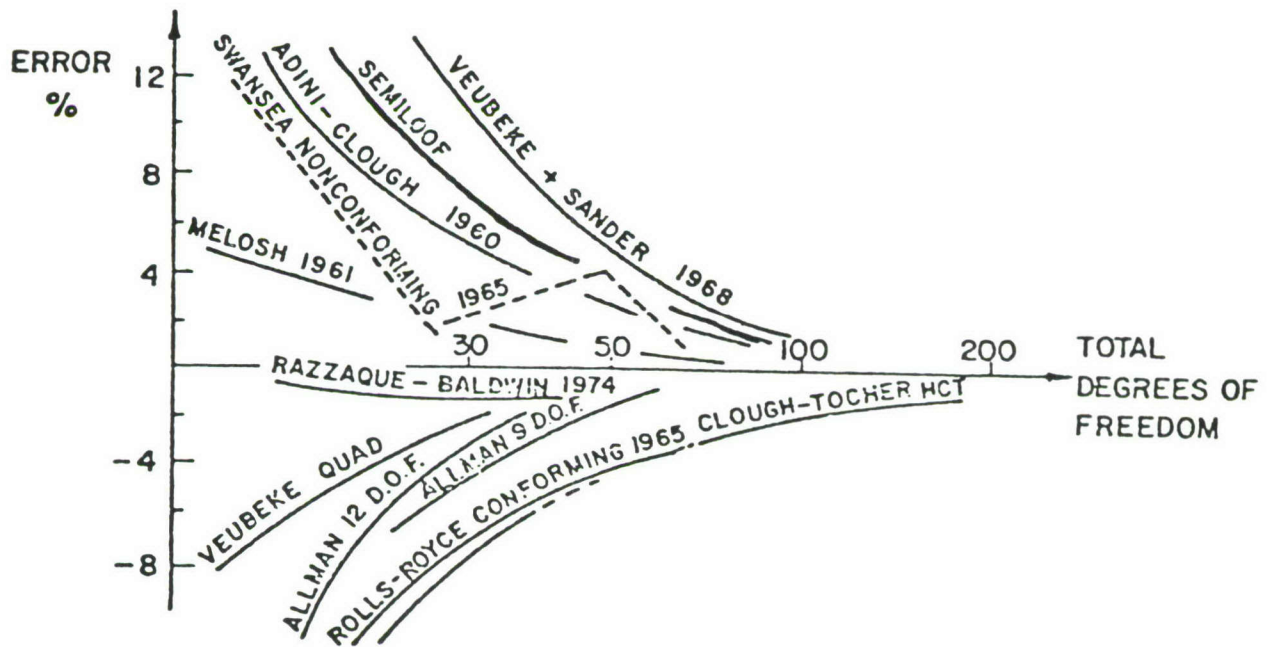


Figure 7 Point-loaded Plate Evaluation

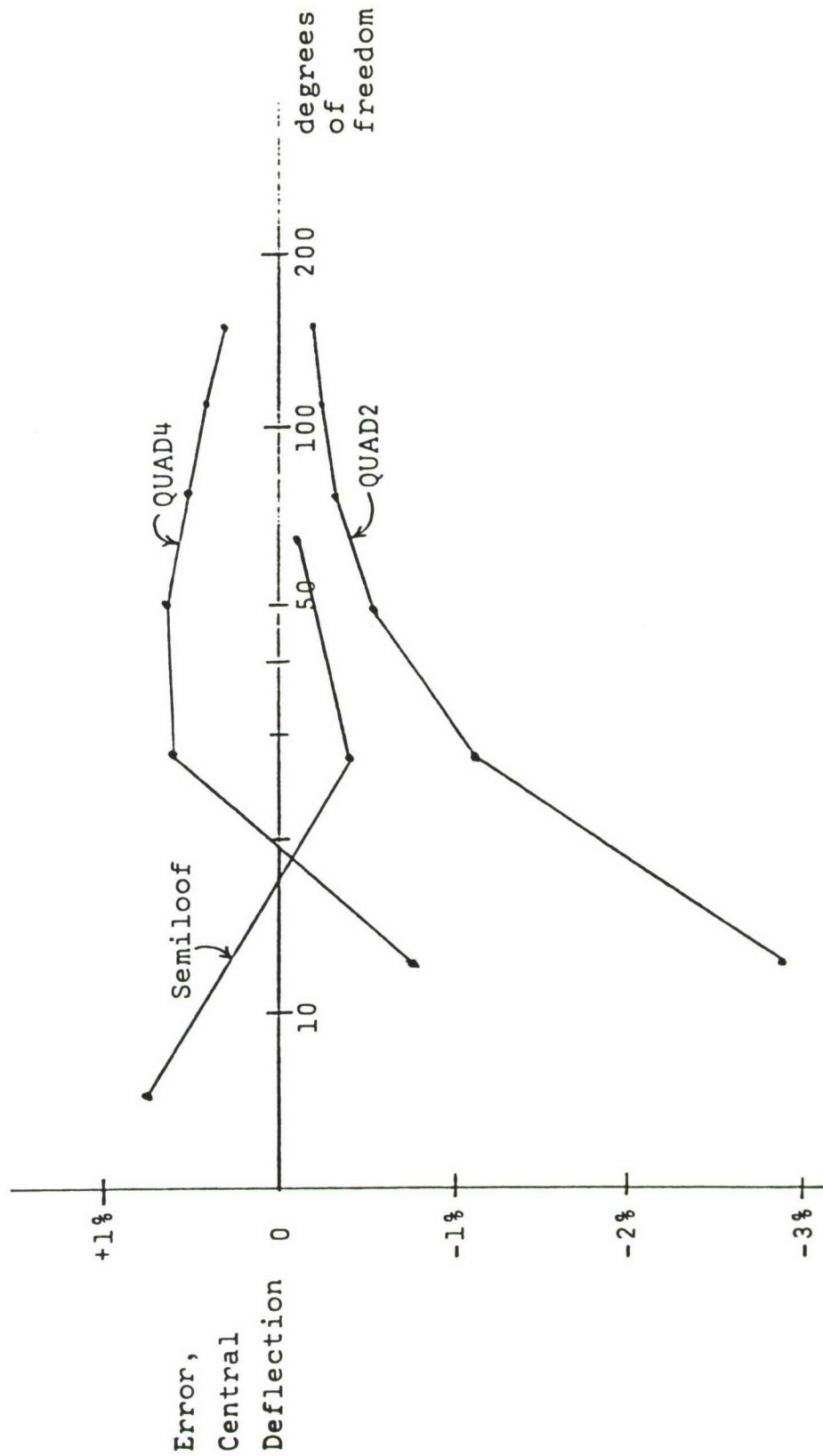


Figure 8 Plate Element Convergence Study

coarse mesh. Incidentally, the degree-of-freedom counts for these runs excluded in-plane displacements.

Next, Problem 3-1-1 of the NASTRAN Demonstration Problem Manual [8] was run. This is a vibration problem for a flat plate. The following frequencies were obtained:

Mode	Theoretical	QUAD1 10x20 Mesh 590 DOF	QUAD1 20x40 Mesh ~1200 DOF	Semi-Loof 4x7 Mesh 196 DOF
1	0.9069	0.9056	0.9066	0.9069
2	2.2672	2.2634		2.2632
3	4.5345	4.5329		4.5302

Although this exercise was mainly intended to check out the mass matrix coding and not necessarily to compare the two elements, there is some evidence of increased efficiency for Semi-Loof.

The third academic test problem that was selected was a cantilever cylinder with varying thickness. Although no analytical solution is known, it was possible to observe the rate of convergence (of the first fundamental frequency) as the models were refined (Semi-Loof versus QUAD4). The following table summarizes these runs:

	<u>Semi-Loof</u>	<u>QUAD4</u>
First Frequency	770.6	767.8
Degrees of Freedom (before omit)	504	900
Degrees of Freedom (after omit)	504	457
Number of Modes Computed	2	1
Total Run "Cost" Including Matrix Generation	<u>\$8.70</u>	<u>\$21.30</u>

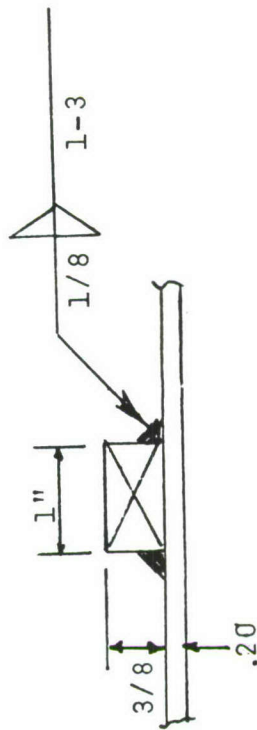
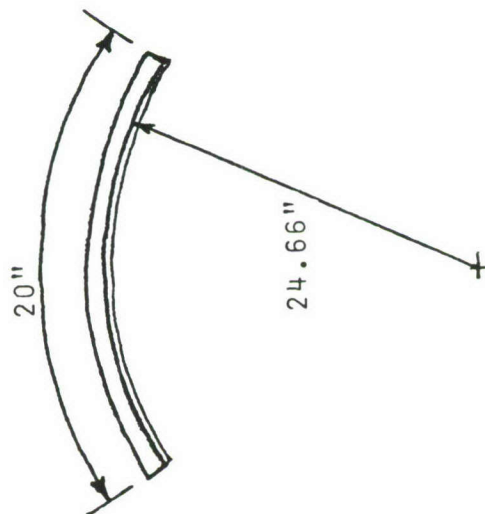
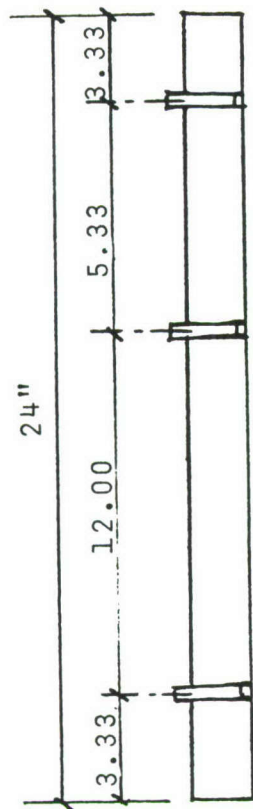
As a final test problem, a curved, stiffened aluminum panel was fabricated (see Figure 9). The panel was tested in the laboratory and was analyzed with both QUAD4 and Semi-Loof elements. The first few frequencies were:

<u>Mode</u>	<u>Description</u>	<u>Experimental</u>	<u>QUAD4 710 DOF</u>	<u>Semi-Loof 500 DOF</u>
1	First twisting	66.2	68.1	67.1
2	First bending	194.4	194.7	186.9
3	Second bending	228.1	239.0	232.5
4	Second twisting	241.9	247.3	241.8
5	Third bending	275.0	284.9	281.3

Note that the Semi-Loof model reproduces the twisting modes somewhat better than the bending modes. This is thought to be a reflection of an early deficiency in the Semi-Loof beam element: the lack of an offset specification to represent eccentric stiffeners. There was some difficulty in choosing accurate equivalent non-eccentric beam properties to model the stiffener. This deficiency was remedied later in the program and the appropriate code was added to PRELOOF.

During Phase I it was thought that rotary inertia, which is usually neglected in finite element mass matrices, might have more importance in angular vibration problems than is generally the case. Rotary inertia is the moment induced in a shell due to a unit angular acceleration about a line tangent to the shell's middle surface. If the two orthogonal rotations for a shell are $\partial w / \partial x$ and $\partial w / \partial y$, then, integrating over the shell surface, we would add the term

$$\iint \left[\left(\frac{\partial w}{\partial x} \right)^2 + \left(\frac{\partial w}{\partial y} \right)^2 \right] \frac{\rho t^3}{12} d\xi d\eta$$



Stiffener Section (Typical)

Material - Aluminum

Figure 9 Curved Panel Configuration

to the mass matrix, where the thickness, t , may vary with ξ and η . In a test case where rotary inertia was added, the natural frequencies changed in only the fourth decimal place. Rotation degrees-of-freedom typically changed in the sixth or seventh decimal place. For a thin shell it is clear that rotary inertia is entirely insignificant. For rotary inertia to be significant, either thicknesses or rotations would have to be large. If thicknesses were large, then the assumptions about zero transverse shear strain would be wrong. If rotations were large, then the linear strain-displacement relations would no longer hold. Although it costs virtually nothing to add rotary inertia to Semi-Loof, it would provide no real pay-off and would be inconsistent with the assumptions underlying the stiffness formulation.

3.4 IMPLEMENTATION INTO NASTRAN

The subroutines supplied by Irons do the basic work of defining shape functions for Semi-Loof elements and applying constraints to reduce out unwanted degrees of freedom. A primary goal of this effort was to make Semi-Loof elements available to NASTRAN users. Considerable effort was involved in writing code around the shape function routines to make this possible.

One of the first decisions faced was whether to modify NASTRAN at the source code level. This would have produced the "cleanest" result in that no interfacing would have been necessary. The alternative was to generate matrices in a pre-processing program and read them into NASTRAN. The former alternative was rejected for the following reasons:

- (1) There are two major versions of NASTRAN, MSC/NASTRAN and COSMIC NASTRAN. MSC/NASTRAN, used by most industrial and some government users, is not available at the source code level, and these users would presumably be shut out by a development limited to COSMIC NASTRAN, for which source code is available.

- (2) There would be no guarantee that future modifications to NASTRAN would be compatible with modifications made for Semi-Loof.
- (3) Link-editing NASTRAN is a time-consuming process, most likely an overnight computer run. This would hamper debugging and checkout considerably.

The NASTRAN DMAP language allows the user access to any matrix or table at any stage of execution. However, it is difficult to access element matrices and assembly tables. Consequently, it was decided that not only would element matrices be generated by PRELOOF, the preprocessing routine, but would be assembled into master matrices as well. This still allows for connection of Semi-Loof elements to other NASTRAN elements since the matrices read from PRELOOF are added to the matrices generated by NASTRAN for any additional elements that may be present. In addition to PRELOOF, a post-processing routine, POSTLOOF, was written for the purposes of recovering stresses and/or angular deformations. The latter are not independent variables at node points; hence, the option to recover them as secondary variables.

The file handling required to transfer information from PRELOOF to NASTRAN and from NASTRAN to POSTLOOF has been made painless by means of packaged control card procedures and DMAP Alters. Experience has shown that there is practically no penalty for doing preprocessing and post-processing outside of NASTRAN as a result.

PRELOOF and POSTLOOF input are described in Appendix A. Some programming notes for PRELOOF are given there, also. Briefly, PRELOOF does the following:

- (1) Reads and checks input data in NASTRAN Bulk Data format.
- (2) Sorts grid point numbers and establishes a sequencing table for degrees of freedom.

- (3) Generates PLOTTEL cards to enable plotting of shell elements, if desired.
- (4) Generates element stiffness matrices and, if requested, element mass matrices and/or load vectors. This involves formation of modulus matrices, generation of shape functions at integration points, and accumulation of stiffness terms.
- (5) Assembles element matrices into global matrices, which are written out to a file six columns at a time, in a form acceptable to NASTRAN.

3.5 SUMMARY

There has been considerable research activity in finite element analysis, but the results of these efforts have been disappointingly slow in reaching the production user. The demands made upon analysts by angular vibration problems in aircraft require that the best tools be available. Millions of dollars have been invested in finite element software, and the vast majority of these funds have been devoted to matters other than the mathematical formulation of elements; that is, such things as input checking, sorting, sequencing, matrix decomposition techniques, spill logic, internal file handling, and output formatting. It would be foolish to attempt to duplicate this massive investment just for the sake of introducing a new element; hence, the motivation for using NASTRAN with its full complement of solution strategies, output options, etc. It is believed that real progress has been made in delivering a new tool to the analyst for angular vibration and other structural problems. However, this new tool, in turn, intensifies the requirement that the analyst "know his elements." Semi-Loof by nature is more sophisticated than simple plate elements, especially with regard to integration procedures.

Part of a fighter fuselage was obtained and tested in the laboratory. The fuselage was modeled and analyzed with Semi-Loof elements, and tested in the laboratory. Section VII is a discussion and comparison of these results.

SECTION IV

HIGH FREQUENCY METHODS

4.1 INTRODUCTION

A basic assumption from the beginning of Anamet's work on angular vibration has been that deterministic methods of structural modeling would not, by themselves, be sufficient. Experience with airborne optical systems has shown that high frequency, low amplitude motion of optical components can seriously degrade system performance. This high frequency motion represents the aggregate contribution of numerous high order modes of vibration which are too sensitive to small details of construction to be reliably modeled by deterministic methods. While motions induced by high frequency disturbances may be small compared to low frequency contributions, the active servomechanisms used to control the optical beam have their own frequency limitations. They cannot effectively suppress and may actually amplify the effect of disturbances at frequencies beyond a few hundred Hz.

Typical cumulative power data for in-flight vibration at a particular fuselage station is shown in Figures 10 and 11 for a translational and rotational displacement, respectively. (Data courtesy of AFFDL.) It may be observed that the normalized cumulative power ($\int_0^f S(f)df/\sigma^2$) of the rotation converges to unity somewhat more slowly than does the corresponding function of the translation. In both cases only a few percent of total signal power is beyond 50 Hz. Nevertheless, the ability to accurately predict this portion of a response PSD is quite important.

Historically, prediction methods for high frequency vibration environments of airborne equipment have been based on empirical correlations of actual flight test data. While this approach has proven to be useful in many cases, it was considered outside the scope of the current contract. Efforts at

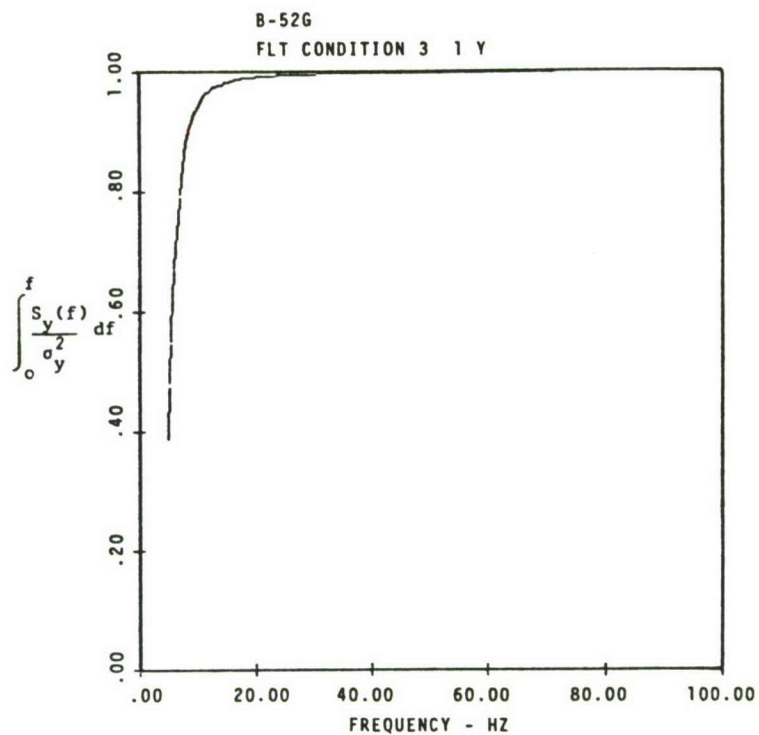


Figure 10 Cumulative Power of Typical Translational Variable for In-flight Vibration.

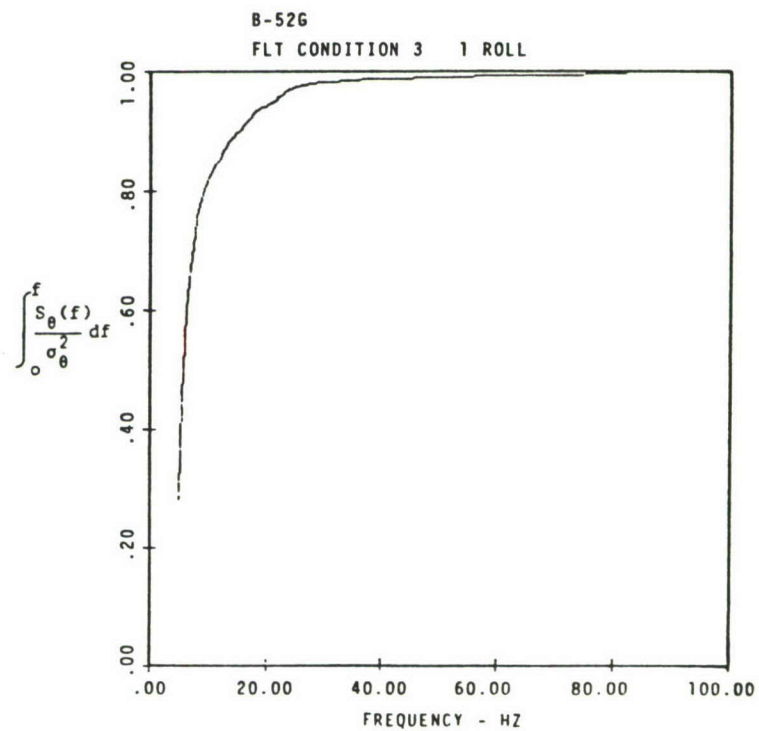


Figure 11 Cumulative Power of Typical Rotational Variable for In-flight Vibration.

Anamet have, therefore, concentrated on more general analytical methods derivable from first principles. The use of experimental data and, in particular, structural models assembled from test data, has by no means been ruled out, however.

The distinction between so-called high frequency and low frequency prediction methods which has been adopted in the current studies is a functional one and does not necessarily involve specific frequency ranges. High frequency prediction methods are defined simply as those which do not necessarily involve knowledge of individual normal modes of a structure. Low frequency methods are those which are based upon such knowledge whether it be obtained through analysis or test.

4.1.1 Statistical Energy Analysis

The definition of high frequency methods stated above leads inevitably to approaches which treat a structure in terms of some averaged, or statistical, description. In Phase I it was decided that the method of Statistical Energy Analysis offered the most promise for predicting high frequency angular vibration. A significant portion of the total effort in Phase II, as well as Phase I, was expended in gaining an understanding of the theoretical and practical basis of the method. This in itself has proven to be a formidable task, mainly because the theory is made up of contributions from several related but diverse fields: classical vibrations, random processes, wave mechanics, mechanical impedance methods, and acoustics. Reference [1] was found to be a most useful document in this regard due to its thoroughness and extensive bibliography.

The single most appealing feature of the SEA approach is its apparent facility for making use of whatever level of detail is available in the description of a structure. This was considered essential since a mode-by-mode description will, by assumption, not be known. The method also appears well suited to situations where a small, indirectly excited structure, such

as an optical element, is attached to a much larger structure, such as an airframe, by a small number of common degrees of freedom.

Section 4.2 of this chapter is devoted to a brief summary of the fundamentals of SEA. It is not intended to describe the entire basis of the method, but rather to point out the essential features which make it attractive for the present purpose. Attention will be focused on a case similar to that described in the preceding paragraph. A wave transmission method for obtaining coupling loss factors is reviewed in Section 4.3.1 and an alternate set of assumptions is presented in Section 4.3.2 which allows a similar formula for coupling loss factor to be obtained by considering a modal coordinate description. It is demonstrated how a convenient measurement procedure can be used to obtain the internal loss factor for a component. In Section 4.6 an experiment is described where the wave transmission method is used to predict transmitted power and equilibrium energy ratio between two coupled plates. Finally, an expression is derived and tested experimentally for predicting the r.m.s. angular displacement in frequency bands at an interior point of a uniform plate under broadband random excitation, given the total vibrational energy.

The purpose of the experiment was to gain some first-hand experience with SEA and to estimate its potential usefulness as a day-to-day engineering tool. In this regard, the use of mini-computer-based FFT methods is described as they pertain to the experiment.

4.2 FUNDAMENTALS OF SEA

Statistical Energy Analysis of vibration is set apart from classical deterministic methods by its use of time-averaged component energies and power flows between components as primary variables. If a component is a distinct mechanical substructure, its energy under stationary random vibration is always expressible as some form of inner product. For example, if component

A is a substructure described by an $N \times 1$ vector \tilde{x} of displacements and a lumped mass matrix \tilde{M} , the time averaged kinetic energy of A is

$$T_A = \frac{1}{2} \langle \tilde{\dot{x}}^T \tilde{M} \tilde{\dot{x}} \rangle_t = \frac{1}{2} \sum_{i=1}^N m_i \langle \dot{x}_i^2 \rangle_t \quad (4.1)$$

where m_i = i'th diagonal entry of \tilde{M}

$$\dot{x}_i = \frac{dx_i}{dt}$$

$\langle \rangle_t$ = expectation with respect to time

The time-averaging and summing operations (4.1) represent the transformation between the individual \dot{x}_i or $\langle \dot{x}_i^2 \rangle_t$ variables of classical deterministic analysis and the energy variables T_A of an SEA model. In general the sum as determined from an SEA model cannot be separated into its individual terms and hence, the identity of individual spatial points is lost. For situations where individual modes are unavailable, this mixing, or averaging process is basic to the method in that it allows equilibrium equations for component energies to be obtained under varying levels of assumptions regarding the nature of the substructure. The wave transmission method, for example, makes rather gross simplifying assumptions by assuming two uniform structures of infinite size. This is appropriate for very high frequency analysis since complications, such as shell curvature or lumped mass concentrations, tend to either have less effect than at low frequencies, or to affect response in only a local area with little change in total energy.

A second inherent characteristic of an SEA model is the decomposition of energy and power flow quantities by frequency. This is analogous to the decomposition of the mean square value

of a zero-mean stationary random quantity into its power spectral density.

$$\sigma_x^2 = \int_0^{\infty} S_x(f) df = \int_0^{f_1} S_x(f) df + \int_{f_1}^{f_2} S_x(f) df + \dots \quad (4.2)$$

It is possible to write equilibrium equations for component energies and inter-component power flow for linear structures under stationary excitation, in terms of frequency bands. The variables are energy or power between specified frequency limits. It should be noted that the terms power spectral density and energy spectral density now mean exactly what they say. This is in contrast to the more common usage of stochastic processes where $S_x(f)$ is called the PSD of x even though x^2 does not have the dimensions of power.

Equilibrium equations are written in terms of simple power balance relations for each component. Suppose two components A and B are connected by a light spring and component A is directly excited (Figure 12).

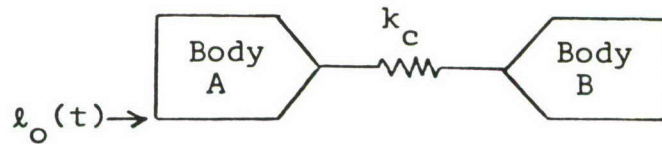


Figure 12 Two Bodies with Stiffness Coupling and Direct Excitation of One Body

$$\langle \pi_{A,IN} \rangle_t = \langle \pi_{A,DIS} \rangle_t + \langle \pi_{AB,TRAN} \rangle_t \quad (4.3)$$

$$\langle \pi_{B,IN} \rangle_t = \langle \pi_{B,DIS} \rangle_t + \langle \pi_{BA,TRAN} \rangle_t \quad (4.4)$$

where

$\langle \pi_{A,IN} \rangle_t$ = input power to body A from all ideal sources
(in this case $\ell_o(t)$)

$\langle \pi_{A,DIS} \rangle_t$ = power dissipated internally by body A

$\langle \pi_{AB,TRAN} \rangle_t$ = power transmitted from body A to body B

with similar definitions for $\langle \pi_{B,IN} \rangle_t$, $\langle \pi_{B,DIS} \rangle_t$, and

$\langle \pi_{BA,TRAN} \rangle_t$.

In equations (4.3) and (4.4) the $\langle \rangle_t$ operator implies time-averaged power in some frequency band. A separate set of equations may be written for each band of interest.

The power quantities on the right hand side of (4.3) and (4.4) are next expressed in terms of component energies by making use of two other basic SEA principles.

As usual in linear vibrations, the energy dissipated per cycle at a given frequency is assumed to be proportional to the amount present at that frequency. The proportionality constant for body A is called $\eta_A(f)$, the internal loss factor at frequency f . As usual, the equation is written in terms of energy and power within specific frequency bands.

$$\begin{aligned} \langle \pi_{A,DIS} \rangle_t &= 2 \pi f_c \eta_A(f_c) \langle E_A \rangle_t \\ \langle \pi_{B,DIS} \rangle_t &= 2 \pi f_c \eta_B(f_c) \langle E_B \rangle_t \end{aligned} \tag{4.5}$$

where

$\langle E_A \rangle_t$, $\langle E_B \rangle_t$ = time-averaged total energies of A and B
(kinetic + potential) within a frequency
band

f_c = center frequency of band

The relation between transmitted power and component energies represents the unique element of SEA analysis. It may be thought of as defining the so-called coupling loss factors $\eta_{AB}(f_c)$ and $\eta_{BA}(f_c)$

$$\langle \pi_{AB,TRANS} \rangle_t = 2\pi f_c [\eta_{AB} \langle E_A \rangle_t - \eta_{BA} \langle E_B \rangle_t] \quad (4.6)$$

For multi-mode structures η_{AB} represents the coupling of each mode of A to all modes of B and thus is a function of $N_B(f_c)$, the number of modes of B occurring within the band about f_c which was used to define the $\langle \rangle_t$ operator. All modes of a component have also been assumed to have equal energies which implies η_{AB} is directly proportional to N_B . Many methods have been used to estimate coupling loss factors. One of the simplest methods (which therefore makes relatively gross assumptions about the structure and excitations) is explained in Section 4.3.

4.2.1 SEA Parameters

For the simplest case where body B is driven only by its connection to body A, it is often assumed that $\langle E_B \rangle_t$ is small compared to $\langle E_A \rangle_t$ and that η_{BA}/η_{AB} is not large compared to unity. In this case (4.6) can be simplified to yield

$$\langle \pi_{AB,TRANS} \rangle_t = 2 \pi f_c \eta_{AB} \langle E_A \rangle_t \quad (4.7)$$

If B is not directly excited $\langle \pi_{B,IN} \rangle_t = 0$ and (4.4) becomes, after combining with (4.5) and rearranging

$$\langle E_B \rangle_t = \langle E_A \rangle_t \left[\frac{\eta_{AB}}{\eta_B + \eta_{BA}} \right] \quad (4.8)$$

It can be shown [9] that equipartition of energy between modes leads to a reciprocity relation between coupling loss factors

$$N_A \eta_{AB} = N_B \eta_{BA} \quad (4.9)$$

Combining (4.8) and (4.9)

$$\langle E_B \rangle_t = \langle E_A \rangle_t \left[\frac{\eta_{AB}}{\eta_B + (N_A/N_B) \eta_{AB}} \right] \quad (4.10)$$

Finally, if the previous assumption of $\eta_{BA} \langle E_B \rangle_t \ll \eta_{AB} \langle E_A \rangle_t$ is applied to (4.6) and the result combined with (4.4) and (4.5)

$$\langle E_A \rangle_t = \frac{\langle \pi_{A,IN} \rangle_t}{\omega(\eta_A + \eta_{AB})} \quad (4.11)$$

Equations (4.7), (4.10) and (4.11) can give some physical insight into the SEA parameters. For a given level of power input to A from the external source, the equilibrium energy level of A is set by the sum of η_A and η_{AB} (4.11). With $\langle E_A \rangle_t$ set, the power transmitted to B is proportional to η_{AB} (4.7). Also for a given level $\langle E_A \rangle_t$, the equilibrium energy level of the indirectly excited body $\langle E_B \rangle_t$ is set by a combination of

η_B , η_{AB} , and the mode counts (4.10). In the limit as the internal damping of B goes to zero, its energy level does not go to infinity as it would for a single body driven by an ideal source. Rather, the energy per mode $\langle E_B \rangle_t / N_B$ of the receptor body B becomes equal to $\langle E_A \rangle_t / N_A$, the energy per mode of the transmitter body. This last conclusion is a fundamental SEA result and is particularly significant in view of the difficulty of predicting internal loss factors in lightly damped structures. It implies that an error in η_B will produce a smaller percentage error in $\langle E_B \rangle_t$ (which is usually the quantity of most interest) and even an assumed value of zero internal loss in B will still produce an answer.

4.3 ESTIMATION OF COUPLING LOSS FACTOR BY THE WAVE TRANSMISSION METHOD

4.3.1 Derivation in Terms of Waves

One of the simplest methods for estimating the coupling loss factors for point-connected structures is the so-called wave transmission approach. It treats the motion of the coupled bodies in terms of traveling waves and represents the coupling loss factor in terms of mechanical impedance or admittance functions at the coupling points. It implies rather extreme simplifying assumptions when applied to real structures since simple traveling wave solutions exist only for structures with a high degree of uniformity. Examples would be infinite plates and beams of uniform section. Nevertheless, the method was pursued for several reasons:

- a. The simplicity of the model made the method a good first step in gaining familiarity with SEA.
- b. It appeared to be adaptable to one of the objectives of the SEA work; investigating the use of interactive minicomputer processing of test data to predict coupling loss factors.
- c. For sufficiently high frequency, the simplifying assumption of uniform structure may be quite reasonable for structures such as curved stiffened shells which are distinctly non-uniform at lower frequency.

The derivation of coupling loss factor in terms of coupling impedance using a wave model is given in Reference [10] and will not be repeated here except to slightly rearrange the final result. Equation 11 of Reference [10] is (refer to Figure 12):

$$2 \pi f_c \eta_{AB} = \frac{1}{m_A} \left| \frac{Z_A Z_B'}{Z_A + Z_B'} \right|^2 \operatorname{Re} \left(\frac{1}{Z_B} \right) \quad (4.12)$$

where

m_A = mass of body A

Z_A = impedance of body A at the degree of freedom to which one end of the coupling spring is attached

Z_B' = impedance looking into the series combination of body B and the coupling spring

We note that a typographical error occurred in Eq. 11 of Reference [10] and so it differs slightly from (4.12) [11]. If Z_B is the impedance of body B at the point where it attaches to the light spring, then

$$\frac{1}{Z_B} = \frac{1}{Z_B} + \frac{1}{Z_C} \quad (4.13)$$

where

$$Z_C = Z_C(f) = \frac{k_C}{i2\pi f}$$

k_C = stiffness of coupling spring

Combining (4.12) with (4.13) and using the definition of mobility $H(f)$ as the reciprocal of impedance $Z(f)$:

$$2\pi f_c \eta_{AB} = \frac{1}{m_A} \frac{\text{Re}(H_B)}{|H_A + H_B + H_C|^2} \quad (4.14)$$

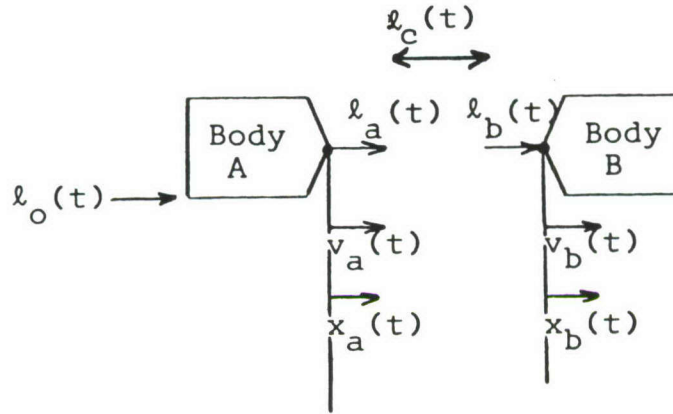
Equation 4.14 is nearly the formula for coupling loss factor upon which the experiment described in section 4.4 is based. However for application to real (i.e. finite) structures, an important modification is necessary. The explanation of this will follow more naturally after an alternate derivation has been given.

4.3.2 Derivation in Terms of Normal Modes

The formula for coupling loss factor, Eq. (4.14), may be derived without explicit reference to traveling waves. The procedure is similar in principal to Reference [10] except that the notions of normal modes and power spectral density functions are used rather than waves of complex exponential amplitude. It is felt that the former terminology would allow a clearer explanation of the experiment which follows.

Consider two structures A and B of Figure 12 connected by a light spring of stiffness k_c . Body A only is directly excited by load $\ell_o(t)$.

Replacing the coupling spring by its force $\ell_c(t)$



where

$\ell_a(t), \ell_b(t)$ = forces on bodies A and B at the connect points

$v_a(t), v_b(t)$ = velocities of the connect points

$x_a(t), x_b(t)$ = displacements at the connect points

$\ell_o(t)$ = external force on body A

$\ell_c(t)$ = coupling force in spring

Then if $\ell_c > 0$ implies tension in the spring

$$\ell_a = \ell_c \quad (4.15)$$

$$\ell_b = -\ell_c \quad (4.16)$$

The mechanical impedance of body B at the attachment point b is defined as

$$Z_b(f) \triangleq \frac{L_b(f)}{V_b(f)} \quad (4.17)$$

where

f = frequency

$$\begin{aligned} L_b(f) &= \text{Fourier transform of } l_b(t) \\ &= \mathcal{F}[l_b(t)] \end{aligned}$$

$$\begin{aligned} V_b(f) &= \text{Fourier transform of } v_b(t) \\ &= \mathcal{F}[v_b(t)] \end{aligned}$$

It should be noted that the definition of $Z_b(\omega)$ implies that $v_b(t)$ is the velocity of B at b when $l_b(t)$ is the only force acting on B. To emphasize this fact, we will write

$v_b^{(l_b)}(t)$ to denote the velocity of b under this condition.

The superscript in parenthesis will also carry over to Fourier transforms.

Transforming (4.16) and combining with (4.17)

$$Z_b(f) = \frac{-L_c(f)}{V_b^{(l_c)}(f)} \quad (4.18)$$

Similarly

$$Z_a(f) = \frac{L_c(f)}{V_a^{(l_c)}(f)} \quad (4.19)$$

The coupling force is

$$l_c(t) = k_c[x_b(t) - x_a(t)] \quad (4.20)$$

Transforming (4.20)

$$L_c(f) = k_c[X_b(f) - X_a(f)] \quad (4.21)$$

Using the derivative relation for Fourier transforms

$$V(f) = i2\pi f X(f) \quad (4.22)$$

gives

$$L_c = \frac{k_c}{i2\pi f} (V_b - V_a) \quad (4.23)$$

If both structures are assumed to behave linearly we may use superposition

$$\begin{aligned} v_a &= [v_a \text{ due to } \ell_o \text{ with } \ell_c = 0] \\ &\quad + [v_a \text{ due to } \ell_c \text{ with } \ell_o = 0] \end{aligned} \quad (4.24)$$

Putting (4.24) in terms of the notation previously defined

$$v_a = v_a^{(\ell_o)} + v_a^{(\ell_c)} \quad (4.25)$$

Transforming (4.25)

$$V_a(f) = V_a^{(\ell_o)}(f) + V_a^{(\ell_c)}(f) \quad (4.26)$$

From (4.26) and (4.19), dropping the argument

$$V_a = V_a^{(\ell_o)} + \frac{L_c}{Z_a} \quad (4.27)$$

Similarly

$$V_b = V_b^{(\ell_o)} - \frac{L_c}{Z_b} \quad (4.28)$$

But $V_b^{(\ell_o)} = 0$ since body B won't move if ℓ_c is removed. Then

$$V_b = -\frac{L_c}{Z_b} \quad (4.29)$$

Combining (4.23), (4.27) and (4.29)

$$L_c = \frac{-V_a^{(\ell_o)}}{\frac{1}{Z_a} + \frac{1}{Z_b} + \frac{1}{Z_c}} \quad (4.30)$$

where

$$\begin{aligned} Z_c(f) &\triangleq \text{impedance of spring element with one end fixed} \\ &= \frac{k_c}{i2\pi f} \end{aligned}$$

Next, we find an expression for $\langle \pi_{AB} \rangle_t$ in terms of $V_a^{(\ell_o)}$. $\langle \pi_{AB} \rangle_t$ is the time averaged power being transmitted from body A to body B through the coupling spring. If $\pi_{AB}(t)$ is the instantaneous power being received at b then

$$\pi_{AB}(t) = v_b(t) \ell_b(t) \quad (4.31)$$

$$= -v_b(t) \ell_c(t) \quad (4.32)$$

If the vibration is assumed to be stationary random and ergodic, we can time average (4.32) to get

$$\langle \pi_{AB} \rangle_t = \lim_{T \rightarrow \infty} \frac{-1}{T} \int_0^T v_b \ell_c dt = - \langle v_b \ell_c \rangle_t \quad (4.33)$$

Introducing the definition of cross correlation

$$R_{v_b \ell_c}(\tau) \triangleq \langle v_b(t) \ell_c(t+\tau) \rangle_t \quad (4.34)$$

From (4.33) and (4.34)

$$\langle \pi_{AB} \rangle_t = - R_{v_b \ell_c}(0) \quad (4.35)$$

The correlation $R(\tau)$ is related to power spectral density $S(f)$ by Wiener's theorem

$$R_{xy}(\tau) = \mathcal{F}^{-1}[S_{xy}(f)] \quad (4.36)$$

$$= \int_{-\infty}^{\infty} S_{xy}(f) e^{i2\pi f\tau} df \quad (4.37)$$

$$= \int_{-\infty}^{\infty} \langle X^*(f)Y(f) \rangle e^{i2\pi f\tau} df \quad (4.38)$$

where the last equation makes use of the stationarity assumption. Then letting $\tau = 0$ in (4.34) and (4.38) and combining with (4.33) gives

$$\langle \pi_{AB} \rangle_t = - \int_{-\infty}^{\infty} \langle L_c^* V_b \rangle df \quad (4.39)$$

where $*$ denotes complex conjugate and $\langle \rangle$ applied to products of Fourier transforms denotes ensemble averaging.

Now $v_b(t)$ and $l_c(t)$ are both real so $V_b(f)$ and $L_c(f)$ both have Hermetian symmetry; that is, $V_b(-f) = V_b^*(f)$ and $L_c(-f) = L_c^*(f)$. The integral in (4.39) is then

$$\int_{-\infty}^{\infty} \langle L_c^* V_b \rangle df = 2\text{Re} \left[\int_0^{\infty} \langle L_c^* V_b \rangle df \right] \quad (4.40)$$

Then

$$\langle \pi_{AB} \rangle_t = - 2\text{Re} \int_0^{\infty} \langle L_c^* V_b \rangle df \quad (4.41)$$

(4.41) suggests that we define $\Pi_{AB}(f)$, the spectral density of transmitted power, as

$$\Pi_{AB}(f) \triangleq -2\text{Re} \langle L_c^* V_b \rangle \quad (4.42)$$

The convention implied in (4.42), namely that Fourier transforms are double-sided and spectral density functions are single-sided, will be followed throughout this chapter.

From (4.41) and (4.42)

$$\langle \pi_{AB} \rangle_t = \int_0^\infty \Pi_{AB}(f) df \quad (4.43)$$

Since $\ell_c(t)$ is the only force acting on body B we can write, from (4.42)

$$\Pi_{AB} = -2\text{Re} \langle L_c^* V_b^{(\ell_c)} \rangle \quad (4.44)$$

Combining (4.18) and (4.44)

$$\Pi_{AB}(f) = 2\langle |L_c(f)|^2 \rangle \text{Re} \left(\frac{1}{Z_b(f)} \right) \quad (4.45)$$

Next define $E_A^{(\ell_o)}(f)$ as the spectral density of the time-averaged total energy (kinetic and potential) of body A with the coupling removed. That is, if $T_A^{(\ell_o)}(f)$ and $U_A^{(\ell_o)}(f)$ are respectively the spectral densities of kinetic and potential energies of the uncoupled body A

$$T_A^{(\ell_o)}(f) + U_A^{(\ell_o)}(f) = E_A^{(\ell_o)}(f) \quad (4.46)$$

It may be shown [12] that for the uncoupled body A driven only by $\ell_o(t)$ the time averaged kinetic and potential energies in any frequency band are equal.

$$T_A^{(\ell_o)} = U_A^{(\ell_o)} \quad (4.47)$$

Then from (4.46) and (4.47)

$$2T_A^{(\ell_o)}(f) = E_A^{(\ell_o)}(f) \quad (4.48)$$

Next, consider the quantity $\langle |V_a^{(l_o)}|^2 \rangle$. If the modal density of body A is very high in the frequency band of interest, and the mass distribution is fairly uniform, then the vibration field will be homogeneous. That is, the temporal mean square velocity in a frequency band will not vary much over the structure. Then, with the coupling spring detached, point a on body A is just a typical point whose velocity spectrum is a good estimate of the mass-averaged velocity spectrum of the whole structure. Thus, if m_A is the mass of body A

$$T_A^{(l_o)}(f) = m_A \langle |V_A^{(l_o)}(f)|^2 \rangle \quad (4.49)$$

It should be noted that the above assumptions are just those which are needed to admit a simple traveling wave description of the motion of body A. Thus, it is not surprising that the end result is similar to equation 11 of Reference [10], which was obtained using a wave approach.

From Eqs. (4.48) and (4.49)

$$E_A^{(l_o)}(f) = 2m_A \langle |V_A^{(l_o)}(f)|^2 \rangle \quad (4.50)$$

The fundamental theorem of SEA is

$$\int_{\Delta f} \Pi_{AB}(f) df = 2\pi f_c [\eta_{AB}(f_c) \int_{\Delta f} E_A(f) df - \eta_{BA}(f_c) \int_{\Delta f} E_B(f) df] \quad (4.51)$$

where $\eta_{AB}(f_c)$ and $\eta_{BA}(f_c)$ are the coupling loss factors evaluated for the frequency band Δf centered at f_c . It must be noted that equation (4.51) holds only in the mean. Each of the three terms represents an estimate of the mean value of a random variable and thus the terms are themselves random variables. In practical terms this implies that the product of modal density and Δf must be large enough to allow many modes to contribute if (4.51) is to hold for a particular trial [13].

When the coupling is light and only one body is directly excited, it is commonly assumed that the coupled and uncoupled energies of the directly excited body (body A) are essentially equal. Also, the energy of the indirectly excited body is assumed to be small compared to that of the directly excited. Under these assumptions (4.51) becomes

$$\int_{\Delta f} \Pi_{AB}(f) df = 2\pi f_c \eta_{AB} \int_{\Delta f} E_A^{(\ell_o)}(f) df \quad (4.52)$$

Combining (4.45) and (4.30)

$$\Pi_{AB}(f) = \frac{2 \langle |V_a^{(\ell_o)}|^2 \rangle}{\frac{1}{Z_a} + \frac{1}{Z_b} + \frac{1}{Z_c}} \operatorname{Re}\left(\frac{1}{Z_b}\right) \quad (4.53)$$

Substituting (4.53) and (4.50) into (4.52) and introducing the velocity admittance $H = Z^{-1}$ (also called the mobility) in place of impedance gives

$$\begin{aligned} & \int_{\Delta f} \frac{2 \langle |V_A^{(\ell_o)}|^2 \rangle \operatorname{Re}(H_b) df}{|H_a + H_b + H_c|^2} \\ &= 4\pi f_c m_A \eta_{AB} \int_{\Delta f} \langle |V_a^{(\ell_o)}|^2 \rangle df \end{aligned} \quad (4.54)$$

The mobility functions are defined at every frequency f both inside and outside the integration band Δf . For structures with the assumed high modal density, $H_A(f)$ and $H_B(f)$ can be expected to vary rapidly within the band and thus the factor $\operatorname{Re}(H_B)/|H_A + H_B + H_C|^2$ cannot be taken outside the integral on the left side of (4.54). However, if one returns to the basic assumptions of SEA the next step is clear. The energy equilibrium equations of the form (4.3)-(4.4) are written for frequency bands of finite width on the assumption that internal and

coupling loss coefficients can be treated as constant over each of these frequency bands. Thus, consistency requires that the quantity $\text{Re}(H_B)/|H_A+H_B+H_C|^2$ be replaced by a constant value which is, in some sense, best over the frequency band Δf . Intuitively, one suspects that the frequency average over Δf of the whole quantity would be appropriate. However, a slightly different approach is suggested by Manning [11] and is the one actually taken. The individual frequency responses $H_A(f)$, $H_B(f)$, and $H_C(f)$ are replaced by their frequency averages over the band Δf and centered at f . It can be shown [14] that for many types of structures, the mobility function of a finite structure, when averaged over a band which contains many modes, will equal the mobility of an infinite extension of the structure. Thus, replacing $H(f)$ by $\bar{H}(f_c)$ where

$$\bar{H}(f_c) = \frac{1}{\Delta f} \int_{f_c - \frac{\Delta f}{2}}^{f_c + \frac{\Delta f}{2}} H(f) df \quad (4.55)$$

amounts to replacing the actual structure by an equivalent one which is a best uniform approximation in a particular frequency range. Making the above substitutions for $H_A(f)$, $H_B(f)$ and $H_C(f)$ allows the mobilities to be taken outside the integral on the left. The remaining integrals then cancel and the result is

$$2\pi f_c \eta_{AB} = \frac{1}{m_A} \frac{\text{Re}[\bar{H}_B(f_c)]}{|\bar{H}_A(f_c) + \bar{H}_B(f_c) + \bar{H}_C(f_c)|^2} \quad (4.56)$$

Note that this result is identical to Eq. (4.14) except that the admittance functions H_A , H_B and H_C have been replaced by their smoothed (i.e., averaged with respect to frequency) counterparts. The final test of the approximations described above must, of course, be by experiment. One such experiment is described in Section 4.5.

4.4 ESTIMATION OF ANGULAR RESPONSE

In Section 4.3 of this report a method is described for calculating coupling loss factor in a particular situation. From the discussion of Section 4.2 it may be seen that this factor, along with internal (damping) loss factors, can be used to set up an SEA model for a simple two-component system coupled at a single degree of freedom. The solution of the SEA equilibrium equations will then be component energies and the inter-component power flow in frequency bands. However, these are not usually the quantities of direct interest. In the present case we would like to estimate mean square angular response and, if possible, its spectral distribution at specific points on a structure. In Section 4.5, the SEA modeling described in Sections 4.2 and 4.3 is carried out for a simple system consisting of two uniform plates coupled at one point. In this section we derive an expression for angular response in terms of component energy for this demonstration case. Its validity was tested by experiment and results are reported in Section 4.5.4.

4.4.1 Relation Between Angular Response and Total Energy For A Uniform Plate

Suppose a uniform rectangular plate is excited by broadband random forces. Assume that the smoothed spectral density of the total vibrational energy is fairly flat well into the frequency range where modal density (averaged over the smoothing bandwidth) is constant. Let the dynamic response of the plate be represented as

$$w(x,y,t) = \sum_i \eta_i(t) \psi_i(x,y) \quad (4.57)$$

where

x, y = coordinates parallel to plate edges

t = time

w = normal displacement of plate midplane

$\psi_i(x, y)$ = shape function of i^{th} normal mode of plate

$\eta_i(t)$ = amplitude of i^{th} normal mode.

It is assumed that if the support conditions allow rigid body modes, the forcing is such that their response is negligible compared to flexural modes in the high frequency range of interest. The desired response quantity is the amplitude of the total angle between the instantaneous normal to the plate at a point (x, y) and the normal to the undisturbed plate at that point.

$$|\theta(x, y, t)| = |\nabla w| \quad (4.58)$$

where

∇ = gradient operator

$$= \vec{e}_x \frac{\partial}{\partial x} + \vec{e}_y \frac{\partial}{\partial y}$$

\vec{e}_x, \vec{e}_y = unit vectors in x and y directions

combining Eqs. (4.57) and (4.58), squaring, and taking the expectation with respect to time.

$$\langle \theta^2(x, y, t) \rangle_t = \sum_i \sum_j \langle \eta_i \eta_j \rangle_t \frac{\partial \psi_i}{\partial x} \frac{\partial \psi_j}{\partial x} + \sum_i \sum_j \langle \eta_i \eta_j \rangle_t \frac{\partial \psi_i}{\partial y} \frac{\partial \psi_j}{\partial y} \quad (4.59)$$

For light damping, each of the mean square modal responses will have spectral densities which have large values in the neighborhood of their respective natural frequencies but are much smaller elsewhere. This implies

$$\langle \dot{\eta}_i^2 \rangle_t \approx \omega_i^2 \langle \eta_i^2 \rangle_t \quad (4.60)$$

and

$$\langle \eta_i \eta_j \rangle_t = 0 \text{ for } \omega_i \neq \omega_j \quad (4.61)$$

The time-averaged total energy of vibration in a lightly damped structure can be represented in modal coordinates as

$$\langle E \rangle_t = \sum_i \langle E_i \rangle_t = \sum_i m_i \langle \dot{\eta}_i^2 \rangle_t \quad (4.62)$$

where

$\langle E \rangle_t$ = time-averaged total energy, kinetic plus potential, of the plate

$\langle E_i \rangle_t$ = time-averaged energy of i^{th} mode

$$\dot{\eta}_i = \frac{d\eta_i}{dt}$$

m_i = modal mass of i^{th} mode.

As usual in SEA work, the $\langle \rangle_t$ operator applied to a positive definite quantity implies the division of that quantity by frequency bands; that is, if $G_u(f)$ is a single-sided spectral density function, then

$$\langle u^2 \rangle_t = \int_{f_c - \frac{\Delta f}{2}}^{f_c + \frac{\Delta f}{2}} G_u(f) df \quad (4.63)$$

Equation (4.63) defines the common SEA notation even though the center frequency f_c and bandwidth Δf are not expressly noted in the left-hand side. We assume that N modes are resonant within Δf and that only their responses need be considered in forming $\langle \theta^2 \rangle_t$.

Combining Equations (4.59) and (4.61) and simplifying

$$\langle \theta^2(x,y,t) \rangle_t = \sum_{i=1}^N \langle \eta_i^2 \rangle_t |\nabla \psi|^2 \quad (4.64)$$

Combining Equations (4.60), (4.62) and (4.64)

$$\langle \theta^2(x,y,t) \rangle_t = \sum_{i=1}^N \frac{\langle E_i \rangle_t}{m_i \omega_i^2} |\nabla \psi_i|^2 \quad (4.65)$$

Next we invoke the basic SEA assumption of equipartition of energy between the N modes of a single component in a band.

$$\langle E_i \rangle_t = \frac{\langle E \rangle_t}{N} \quad (4.66)$$

From Equations (4.65) and (4.66), writing out $|\nabla\psi|^2$

$$\langle \theta^2(x,y,t) \rangle_t = \sum_{i=1}^N \frac{\langle E \rangle_t}{m_i \omega_i^2 N} \left[\left(\frac{\partial \psi_i}{\partial x} \right)^2 + \left(\frac{\partial \psi_i}{\partial y} \right)^2 \right] \quad (4.67)$$

The frequency range of interest is high compared to the fundamental mode of the plate. Therefore, away from the boundaries the mode shapes will be approximately sinusoidal regardless of boundary conditions.

$$\psi_i(x,y) = \sin(k_{xi}x) \sin(k_{yi}y) \quad (4.68)$$

where k_{xi} and k_{yi} are allowable wavenumbers in the x and y directions. Their exact values will depend on boundary conditions, but in the high frequency regime, only their average density in wavenumber space is important. This density is independent of boundary conditions. In addition, by assuming a traveling wave solution, we may write the dispersion relation for a plate independent of the type of supports.

$$\omega_i^2 = (k_{xi}^2 + k_{yi}^2)^2 \frac{\kappa^2 c_\ell^2}{(1-\nu^2)} \quad (4.69)$$

where

ω_i = natural radian frequency of i^{th} mode

κ = radius of gyration of plate cross section

= $h/2\sqrt{3}$

h = plate thickness

c_ℓ = extensional wave speed

E = plate elastic modulus

ρ = plate mass density

ν = Poisson's ratio

Combining Equations (4.67) and (4.68)

$$\begin{aligned} \langle \theta^2(x,y,t) \rangle_t = \sum_{i=1}^N \frac{\langle E \rangle_t}{m_i \omega_i^2 N} [k_{xi}^2 \cos^2(k_{xi}x) \sin^2(k_{yi}y) \\ + k_{yi}^2 \sin^2(k_{xi}x) \cos^2(k_{yi}y)] \end{aligned} \quad (4.70)$$

In the wave transmission method for finding power transfer coefficients, it is assumed that many modes fall within the frequency band Δf and thus N is large compared to 1. We, therefore, assume that the "typical" point $(x'y')$ falls on the different mode shapes at points such that the quantities $k_{xi}x'$ and $k_{yi}y'$ are uniformly distributed over the range from 0 to 2π . Under this assumption we may replace the $|\nabla\psi_i|^2$ quantity on the right hand side of Equation (4.70) by its spatial average taken over many cycles in both x and y directions. We may also drop the x,y arguments on the left hand side, requiring only that $x'y'$ be well away from the plate boundaries. In that case

$$\langle \theta^2 \rangle_t = \sum_{i=1}^N \frac{\langle E \rangle_t}{m_i \omega_i^2 N} \frac{(k_{xi}^2 + k_{yi}^2)}{4} \quad (4.71)$$

For a uniform mass density, the modal mass corresponding to mode shapes given by Equation (4.68) is

$$m_i = \frac{M}{4} \quad (4.72)$$

where M is the plate physical mass. Combining Equations (4.69), (4.71) and (4.72)

$$\langle \theta \rangle_t = \frac{\langle E \rangle_t \sqrt{1-\nu^2}}{NMkc_\ell} \sum_{i=1}^N \frac{1}{2\pi f_i} \quad (4.73)$$

Once again, the assumption of high modal density can be used to simplify the result. If the natural frequencies f_i are assumed to be uniformly distributed over the frequency band Δf we have

$$\sum_{i=1}^N \frac{1}{f_i} \approx \frac{N}{\Delta f} \int_{f_c - \frac{\Delta f}{2}}^{f_c + \frac{\Delta f}{2}} \frac{df}{f} \quad (4.74)$$

Note that $N/\Delta f$ is the average modal density and $\Delta f/N$ is the average mode spacing. Combining Eqs. (4.73) and (4.74) and carrying out the integration.

$$\langle \theta^2 \rangle_t = \frac{\langle E \rangle_t \sqrt{1-\nu^2}}{2\pi M \kappa c_\ell} \frac{\ln \left[\frac{f_c + \frac{\Delta f}{2}}{f_c - \frac{\Delta f}{2}} \right]}{\Delta f} \quad (4.75)$$

or

$$\langle \theta^2 \rangle_t = \frac{\langle E \rangle_t \sqrt{1-\nu^2}}{2\pi M \kappa c_\ell f_{\ell.m.}} \quad (4.76)$$

where $f_{\ell.m.}$ is the log mean frequency of the band. For the bandwidths and center frequencies used in the coupled plate experiment, it is essentially equal to the arithmetic mean f_c . The final result is then

$$\langle \theta^2 \rangle_t = \frac{\langle E \rangle_t \sqrt{1-\nu^2}}{2\pi M \kappa c_\ell f_c} \quad (4.77)$$

The r.m.s. value of θ is obtained by summing $\langle \theta^2 \rangle$ over all bands and taking the square root.

$$\theta_{r.m.s.} = \frac{(1-\nu^2)^{1/4}}{(2\pi M \kappa c_\ell)^{1/2}} \left(\sum_b \frac{\langle E_b \rangle}{f_b} \right)^{1/2} \quad (4.78)$$

where

$\theta = |\nabla w|$, angle between the instantaneous normal to
 plate at a point x, y , and the normal to the undis-
 turbed plate at that point
 w = out-of-plane displacement
 ν = Poisson's ratio
 M = total plate mass
 $\kappa = h/2\sqrt{3}$, radius of gyration of plate cross section
 h = plate thickness
 $c_\ell = \sqrt{E/\rho}$, extensional wave speed
 E = elastic modulus
 ρ = mass density
 $\langle E_b \rangle$ = temporal mean energy in a frequency band b
 f_b = center frequency of band b

In a sense, the derivation leading to Eq. (4.78) is similar to
 the method of Lee and Whaley [15] for estimating angular vibra-
 tion from translational vibration. Like Lee and Whaley's method,
 it is based on a particular structural form (i.e., a beam or
 plate) and uses the idea of summing or averaging translational
 response over that form. However, it differs in that mass-
 weighted velocity PSD's are added to obtain the spectral den-
 sity of the vibrational energy. Also, this method does not
 require knowledge of boundary conditions or individual normal
 modes. Rather, it is assumed that many modes contribute and
 that low order modes are not dominant.

4.5 ESTIMATION OF INTERNAL LOSS FACTOR

In assembling an SEA model of a real structure, the most
 formidable task is usually the calculation of the coupling loss
 factors. However, the accuracy of the model in predicting
 equilibrium component energies is also quite dependent on the
 accuracy of internal loss factors which are input to the model.
 The damping mechanisms in built up structures are usually com-
 plex. Estimates of internal loss factor are almost always

empirical in nature. Measurements must be made on the structure in question if it is available or inferred from previous measurements on similar structures. In either case, a convenient method of obtaining damping data suitable for SEA modeling is highly desirable. Recent developments in mini-computer-based dynamic testing have provided a method which seems promising. It is described in the next section.

4.5.1 Measurement and Use of Modal Damping Ratios

While it was assumed at the outset that complete descriptions of individual normal modes of interacting structures would not, in general, be known, it is also true that this is a practical rather than a theoretical restriction. Modal densities of airframe structures may be high, but they are finite. It is quite easy to see the contributions of individual modes in the mobility function measured at a coupling degree of freedom if the measured function is sufficiently well resolved in frequency. It is also possible to estimate the equivalent viscous modal damping ratios of individual modes by trial-and-error fitting of a curve of specific mathematical form to the measured function [16, 17]. Again, the measurement may require a high degree of frequency resolution but this can be accommodated by commercially available hardware/software packages which implement the so-called "zoom" discrete Fourier transform [18]. In theory, each modal damping ratio is characteristic of the entire structure and is independent of the excitation used to measure it. In practice, one will obtain reliable modal damping estimates only for modes with significant amplitude at the coupling point. However, these are exactly the modes which will determine the properties of the coupled system, so obtaining their damping ratios should be adequate.

For the two component system which was tested, the desired damping quantity is

$$\omega_c \eta_B = \frac{\langle \pi_{AB} \rangle_t}{\langle E_B \rangle_t} \quad (4.79)$$

At first glance it might appear that the internal loss factor η_B could be obtained from the modal damping ratios $\xi^{(i)}$ simply by averaging $2\xi^{(i)}$ over all modes resonant in the band of interest. However, this method yields poor agreement with direct experimental measurement of η_B via Eq. (4.79) [19]. A better method, suggested by Manning [11], has been found to give excellent agreement and is explained below.

Since essentially all the vibrational energy of a particular mode is in a narrow band about its natural frequency, we may write, for the frequency band which contains $\omega^{(i)}$

$$\langle E_B^{(i)} \rangle_t = \frac{\langle \pi_{AB}^{(i)} \rangle_t}{\omega^{(i)} \eta^{(i)}} \quad (4.80)$$

where

$$\begin{aligned} \langle E_B^{(i)} \rangle_t &= \text{averaged energy in } i^{\text{th}} \text{ mode of B} \\ \langle \pi_{AB}^{(i)} \rangle_t &= \text{averaged power input to } i^{\text{th}} \text{ mode} \\ \omega^{(i)} &= \text{natural radian frequency of } i^{\text{th}} \text{ mode} \\ \eta_B^{(i)} &= \text{loss factor for } i^{\text{th}} \text{ mode of B} \\ &= 2\xi_B^{(i)} \\ \xi_B^{(i)} &= \text{viscous damping ratio of } i^{\text{th}} \text{ mode} \end{aligned}$$

If the total energy of the component is $\langle E_B \rangle_t$, then

$$\langle E_B \rangle_t = \sum_i^{N_B} \langle E_B^{(i)} \rangle_t \quad (4.81)$$

where the summation is over all N_B modes which are resonant within the band. From Eqs. (4.80) and (4.81)

$$\langle E_B \rangle_t = \sum_{i=1}^{N_B} \frac{\langle \pi_{AB}^{(i)} \rangle_t}{\omega^{(i)} \eta^{(i)}} \quad (4.82)$$

The input power to a specific mode will depend on its properties such as shape and modal mass, which are presumably not available. We, therefore, approximate $\langle \pi_{AB}^{(i)} \rangle_t$ for all i by an averaged value $\langle \pi_{AB} \rangle_t / N_B$ where $\langle \pi_{AB} \rangle_t$ is the total input power. Making this substitution in Eq. (4.82) as well as replacing $\eta^{(i)}$ by $2\xi^{(i)}$ and $\omega^{(i)}$ by $2\pi f^{(i)}$

$$\frac{\langle \pi_{AB} \rangle_t}{\langle E_B \rangle_t} = 4\pi \left[\frac{1}{N_B} \sum_{i=1}^{N_B} \frac{1}{f^{(i)} \xi^{(i)}} \right]^{-1} \quad (4.83)$$

Comparing Eqs. (4.79) and (4.83)

$$\omega_c \eta_B = 4\pi \left[\frac{1}{N_B} \sum_{i=1}^{N_B} \frac{1}{f^{(i)} \xi^{(i)}} \right]^{-1} \quad (4.84)$$

The left-hand side of Eq. (4.83) is the quantity desired for SEA modeling. The right-hand side is data available from the complex curve fitting routine.

It should be noted that a fitting algorithm of the type described in [16] will return more complex poles than the structure actually has modes. This is necessary in order to obtain a good fit in the presence of the inevitable noise contamination of the measurement. However, it is usually a simple matter to pick out the correct complex poles by examination of the peaks in the frequency response and the complex residues associated with each indicated pole. A test of Eq. (4.83) is described in a later section where $\langle \pi_{AB} \rangle_t / \langle E_B \rangle_t$ is measured directly and compared with an estimate made by curve fitting of a single admittance function.

4.6 SEA WAVE TRANSMISSION EXPERIMENT

4.6.1 Objective

During Phases II and III of this contract an experiment was carried out to test the theoretical developments (described in Sections 4.3 and 4.4). The specific objectives of the exercise were quite numerous. As a whole, it was intended to be a first step towards the development of SEA modeling techniques, based on a combination of theory and experiment, which could be used in the development of airborne optical systems. More specific objectives included:

- (1) Compare SEA predictions with measurements for transmitted power and energy of an indirectly excited body over various frequency bands. It was later decided to make the comparisons on the basis of $\langle \pi_{AB} \rangle_t / \langle E_A \rangle_t$ and $\langle E_B \rangle_t / \langle E_A \rangle_t$ to reduce the time and cost of the experiment. In this latter format no prediction for $\langle E_A \rangle_t$ itself is needed and some ambiguity is removed from the results.
- (2) Compare predictions of angular response made from Eq. (4.78) and measured energy with directly measured values.
- (3) Acquire some hands-on experience with applying an SEA model to a real structure, albeit a very simple one.
- (4) Test the limits of the assumptions built into the wave transmission formula Eq. (4.56). In particular, the test was run with rather heavy coupling (stiff coupling spring) so that $\langle E_B \rangle_t / \langle E_A \rangle_t$ was about 0.2 over most frequency bands. In developing the SEA model $\langle E_B \rangle_t \ll \langle E_A \rangle_t$ had been assumed.
- (5) Investigate the usefulness of interactive mini-computer processing of test data for SEA modeling. Since SEA models are always constructed individually for different frequency bands, it was felt that FFT methods could be put to good use both in developing the model and in processing of test data for comparison with model predictions.
- (6) Develop preliminary software. A number of sub-routines were written for tasks which have a high

probability of occurring in any future investigation of more sophisticated SEA methods.

- (7) Evaluate numerical procedures for damping measurement. The definition of internal loss factor η_B as utilized in the two-component SEA model of Figure 12 is

$$\eta_B(f_c) = \frac{\langle \pi_{AB} \rangle_t}{2\pi f_c \langle E_B \rangle_t} \quad (4.79)$$

Measurement of η_B directly from this definition is time-consuming even for a simple structure. It may be impractical for anything resembling an airframe since $\langle E_B \rangle_t$ must be assembled as a mass-weighted sum of smoothed velocity PSD's from many response points. A more practical method involving only measurement of coupling point mobilities was described in Section 4.5.

A word of caution is needed on two points regarding the experiment. While the software which was produced is supplied as an appendix to this report, it cannot be considered a finished product. Its purpose so far has been investigation of a method of analysis. A great deal more work can be anticipated prior to producing software suitable for SEA modeling of anything resembling an actual aircraft and optical system. The uniform plates chosen as test structures are not represented as being dynamically similar to the airframe-optical systems which motivate the overall effort. They are of interest simply because they are typical "high frequency" structures; that is, many modes contribute to response and properties of each individual mode are impractical to obtain, either by analysis or test.

Secondly, one should be aware of what constitutes a "correct" SEA prediction. For complex multi-modal structures, any SEA model is by necessity built up from a coarse and incomplete description of the actual. SEA theory is then used to estimate average or typical component energies and infer from them the mean square response (translation or rotation) at an

average or typical point. However, as Lyon [1] points out repeatedly, the analyst is not interested in averages. He desires to predict response for one specific realization; namely response at a transducer mount point on the prototype sitting in the lab. Thus, the variance of a particular trial about the mean must be considered. The derivation of confidence bands for estimates has not been made for the coupled-plate experiment. Rather, for this initial investigation, the test structure was simply chosen to have higher modal density and be much more uniform than a typical airframe. Both properties tend to reduce response variance. It was simply assumed that such variance as remained was held small enough to allow evaluation of the practical potential of SEA.

4.6.2 Experimental Hardware

The experimental set-up is shown in Figures 13 and 14. Two rectangular aluminum plates of slightly different sizes are suspended in parallel vertical planes by hanging them by long, compliant cords. Pendulum frequencies are low enough that in the frequency range of interest the plates behave as if they were in a free-free configuration. Plate dimensions are 65.12" x 48.12" x 0.090" and 79.12" x 48.12" x 0.090". The coupling link was constructed from 3/8" diameter aluminum rod, 5" in length, with a miniature preloaded universal joint at one end to reduce the moment transmitted into the indirectly excited plate (Figure 14). The connection to both plates is slightly off-center to avoid discriminating against either symmetric or antisymmetric bending modes. The connection to the indirectly excited plate is made by means of a short piece of 10-32 threaded rod which passes through a drilled hole in the plate and clamps it between an accelerometer on one side and a load cell on the other. The load cell then connects to one end of the universal joint. It was originally intended to use a light spring as the coupling element, but some difficulty was encountered in fabricating a spring which would behave as a pure

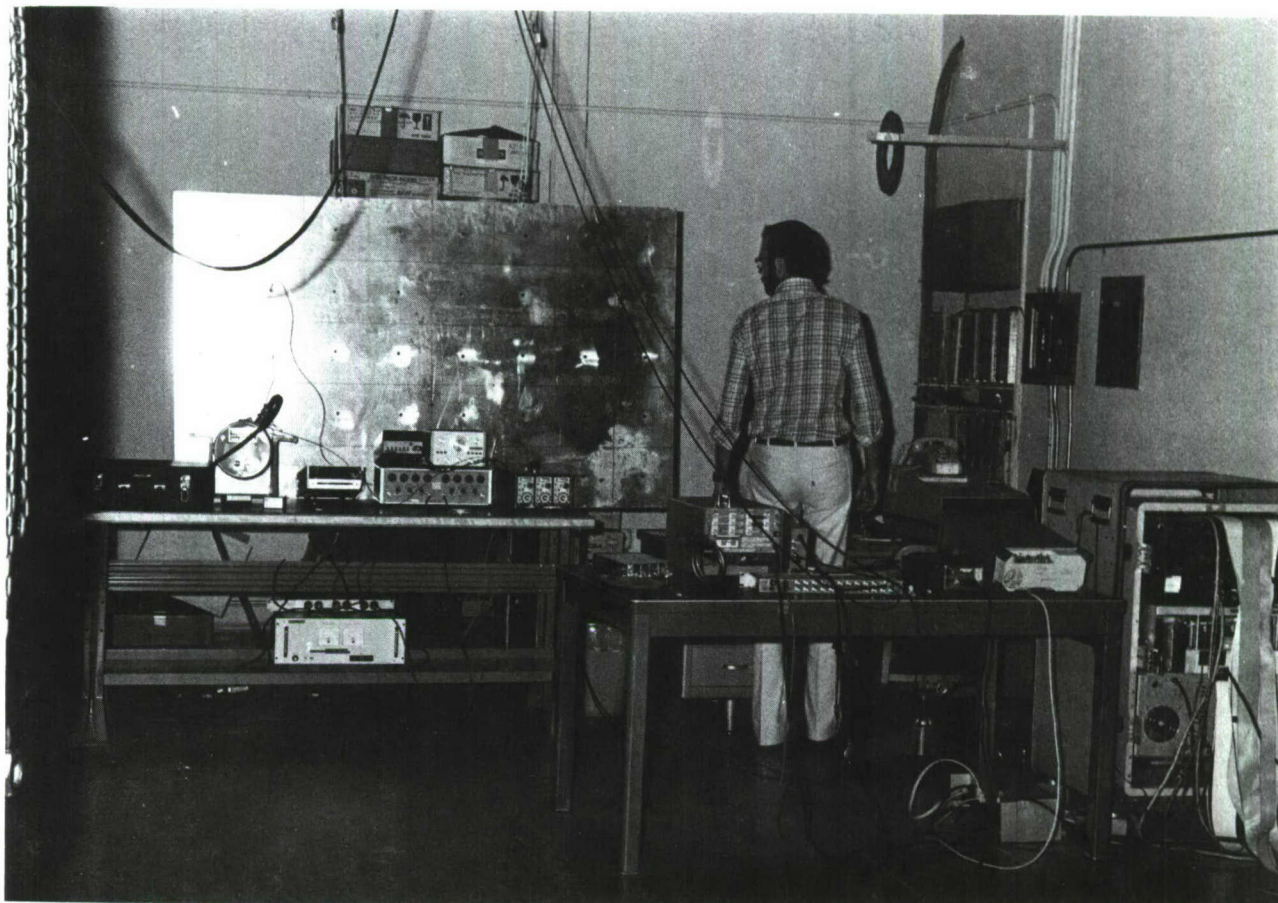


Figure 13 Couples Plates Used for SEA Wave Transmission Experiment

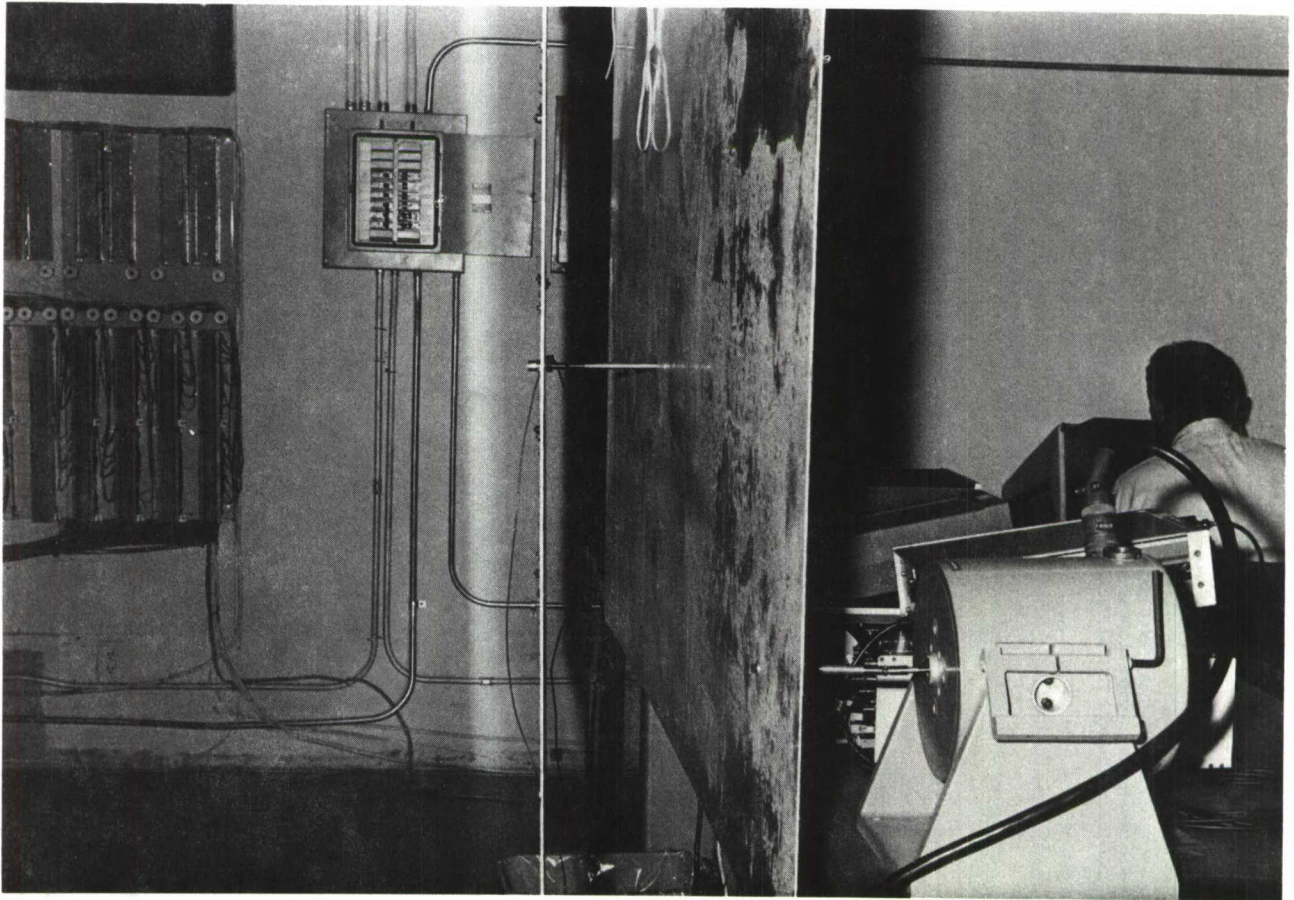


Figure 14 Close-Up of Connecting Link Instrumented to Measure Instantaneous Force and Acceleration

compliance over the desired 100-1000 Hz. band. An attempt was made to use a ring-shaped spring of 163 lbf/in. stiffness which would theoretically have given light coupling (i.e., $H_C \gg H_B$) above 350 Hz. The attempt failed due to resonances of the spring itself well below 1000 Hz. The exercise was, nonetheless, instructive and is described in Section 4.6.4.4.

Body A (the larger plate) was directly excited by a small electrodynamic shaker coupled through a universal joint. The drive signal was white Gaussian noise with band-limiting to 150-1000 Hz. Total power input to plate A was about 1 milliwatt and force level was about 0.3 lbf (rms).

The stiffness of the solid link was measured dynamically as 6163 lbf/in. While this implies fairly heavy coupling throughout the band, it was decided that it should be used, nonetheless. Since real structures are likely to be coupled simply by rigid or near-rigid connection of one or more degrees-of-freedom, it was felt that a deliberate violation of the $H_C \gg H_B$ assumption might be instructive. It may be noted from the derivation in Section 4.3.2 that $\langle E_B \rangle_t / \langle E_A \rangle_t \ll 1$ is the actual format of the light coupling assumption. Data presented later in this section show that this was not grossly violated.

The choice of structures for the wave transmission experiment was based on satisfying, as much as possible, the assumptions which are built into the derivation of Eq. 4.56. It was felt that this would allow attention to be focused on the more specialized question of assembling an SEA model from test data. A prerequisite for the test structures is that modal densities be high. For a uniform plate of finite size, the modal density $n(f)$ in modes per Hz. is [20]:

$$n(f) = \frac{\sqrt{3} A_P}{hc_\ell} \quad (4.85)$$

This implies modal densities of about 24 and 29 modes per 80 Hz. frequency band for the smaller (body B) and larger (body A) plates, respectively.

4.6.3 Experimental Procedure

The test procedure consisted basically of 12 steps.

- (1) Set up plates in the coupled configuration and establish a suitable drive level and drive spectrum shape.
- (2) Measure the smoothed spectral density of transmitted power from force and acceleration signals at the coupling point of plate B.
- (3) Measure the smoothed spectral density of total vibrational energy of each plate from accelerometer signals. A uniform grid of 40 locations on plate A and 35 locations on plate B was used.
- (4) Remove the coupling link and arrange the shaker, force cell, and accelerometer to measure the mobility function looking into the coupling point of plate B.
- (5) Repeat step (4) on plate A.
- (6) Compute frequency-averaged mobility functions for the desired frequency bands using the results of steps (4) and (5).
- (7) Compute the predicted coupling loss factor weighted by $2\pi f_c$ from the results of step (6) and Eq. (4.56).
- (8) Form the actual internal loss factor of plate B as $\langle \pi_{AB} \rangle_t / \langle E_B \rangle_t = 2\pi f_c \eta_B(f_c)$.
- (9) Compute the predicted energy ratio $\langle E_B \rangle_t / \langle E_A \rangle_t$ using the results of steps (7) and (8) and Eq. (4.10).
- (10) Use the results of steps (2) and (3) to compute actual transmitted power and component energy in frequency bands. Form the ratios $\langle \pi_{AB} \rangle_t / \langle E_A \rangle_t$ and $\langle E_B \rangle_t / \langle E_A \rangle_t$ for frequency bands.
- (11) Make comparison plots of actual and predicted values of $\langle E_B \rangle_t / \langle E_A \rangle_t$ and $\langle \pi_{AB} \rangle_t / \langle E_A \rangle_t$.
- (12) Rearrange the hardware to drive a single plate (the larger) and measure steady state energy and angular response at an interior point. Plot data as smoothed spectral densities to verify Eq. (4.78).

In principle, step (12) could have been carried out using an indirect drive to the plate to obtain a more nearly end-to-end comparison between theory and experiment. In fact, this step was carried out as something of an afterthought and the direct drive was used to reduce experimental time and cost. It was noted, however, that the smoothed PSD of force input to the plate was similar for either direct or indirect drive, so the comparison of predicted and measured $\langle \theta^2 \rangle_t$ and θ r.m.s. should be representative.

An error analysis of the differential acceleration method used to measure the spectral density of angular response is given as an example in Section 6.2.1.

In reference to step (11), a second set of predicted values for $\langle E_B \rangle_t / \langle E_A \rangle_t$ and $\langle \pi_{AB} \rangle_t / \langle E_A \rangle_t$ was computed using the theoretical mobility functions for infinite plates of the same material and thickness as the actual plates [21]. Thus, three sets of data are available for the comparison quantities noted in step (11):

- (a) Actual, as measured experimentally.
- (b) Predicted, using a theoretical infinite uniform plate approximation.
- (c) Predicted, using an infinite uniform structure approximation defined by frequency averaging of measured mobility functions.

Virtually all of the power and energy measurements were accomplished by first obtaining a discrete approximation to their spectral density functions by means of digital discrete Fourier transform (DFT) processing with ensemble averaging. This function could then be integrated over any desired frequency band to obtain the $\langle \rangle_t$ quantities of the SEA equilibrium equations. The only restrictions were that the desired frequency band must lie entirely below $1/2T$ and be an integer multiple of $1/NT$ where T is the sampling interval and N is the

time domain block size. Using digital methods, it was not necessary to specify either integration bandwidth or band center frequencies prior to data acquisition. In effect, many very narrow bands were defined and combined afterwards to obtain the desired bandwidth for the SEA model.

Source code listings and a description of a typical test run are provided in Appendix C. These are recommended as being the most detailed documentation of the procedure.

Spectral densities and mobility functions were measured with a 2.5 Hz. frequency resolution up to 1280 Hz. and then smoothed over 32 spectral lines (80 Hz.). A constant bandwidth was chosen in order to keep the number of interacting modes per band roughly constant. An interesting aspect of the digital procedure is that it is not necessary to restrict oneself to $1280/80 = 16$ frequency bands. It is actually more convenient to consider an 80 Hz. averaging band centered over each spectral line at 2.5 Hz. intervals. Thus the results, both predicted and actual, have the appearance of continuous curves, but actually represent a series of heavily overlapped discrete bands. A 40 Hz. band at either end of the base 0-1280 Hz. analysis range must, of course, be excluded from the smoothed results, but this is of no consequence for the intended purpose.

The DFT resolution band of 2.5 Hz. was chosen for a specific reason. It was known that this would be small compared to the smoothing band but still inadequate to obtain a good representation of the rapidly fluctuating mobility function for the lightly damped plates. However, it was desired to measure the effect of this shortcoming on the smoothed version of the function from which the coupling loss factors are actually estimated. If this moderate level of resolution could yield reasonable results, it would have important practical implications for testing of real structures. It should be noted that the actual power and energy comparison quantities are essentially invariant

with respect to DFT resolution. This is so because the smoothing with respect to frequency can be shown to be equivalent to starting with the same time data but dividing it into a larger number of shorter records. The shorter records would produce a coarser frequency resolution but the statistical confidence of the estimate of power within a given frequency band would remain the same.

4.6.4 Experimental Results

4.6.4.1 Principal Results

Figures 15 and 16 present the main results of the wave transmission experiment. Figure 15 compares the actual ratio of transmitted power to transmitting body energy with predicted values obtained using two different representations for coupling point mobility. Figure 16 shows a similar comparison for normalized receptor body energy, and is based on actual measured damping in plate B, as calculated by Eq. (4.79). Disagreement between measured mobility is about 1.5 to 6 dB (factor of 1.4 to 4) in the range beyond 300 Hz. The predicted value of $\langle \pi_{AB} \rangle_t / \langle E_A \rangle_t$ using an infinite plate approximation is somewhat better at 0 to 5 dB error (factor of 1 to 3.2). For $\langle E_B \rangle_t / \langle E_A \rangle_t$ the corresponding agreement figures are 0 to 3 dB for both prediction methods. In general, it appears at this point that the infinite plate approximation is the better of the two. However, some further data presented in the next section suggests that better DFT frequency resolution in the measurement of H_A and H_B would improve the accuracy of the frequency-averaged measured mobility approach. It also appears that a wider smoothing bandwidth would be appropriate and would improve the accuracy of both methods. At any rate, the results are quite encouraging. The receptor body energy is within a factor of two of the measured value in all frequency bands, which implies that a

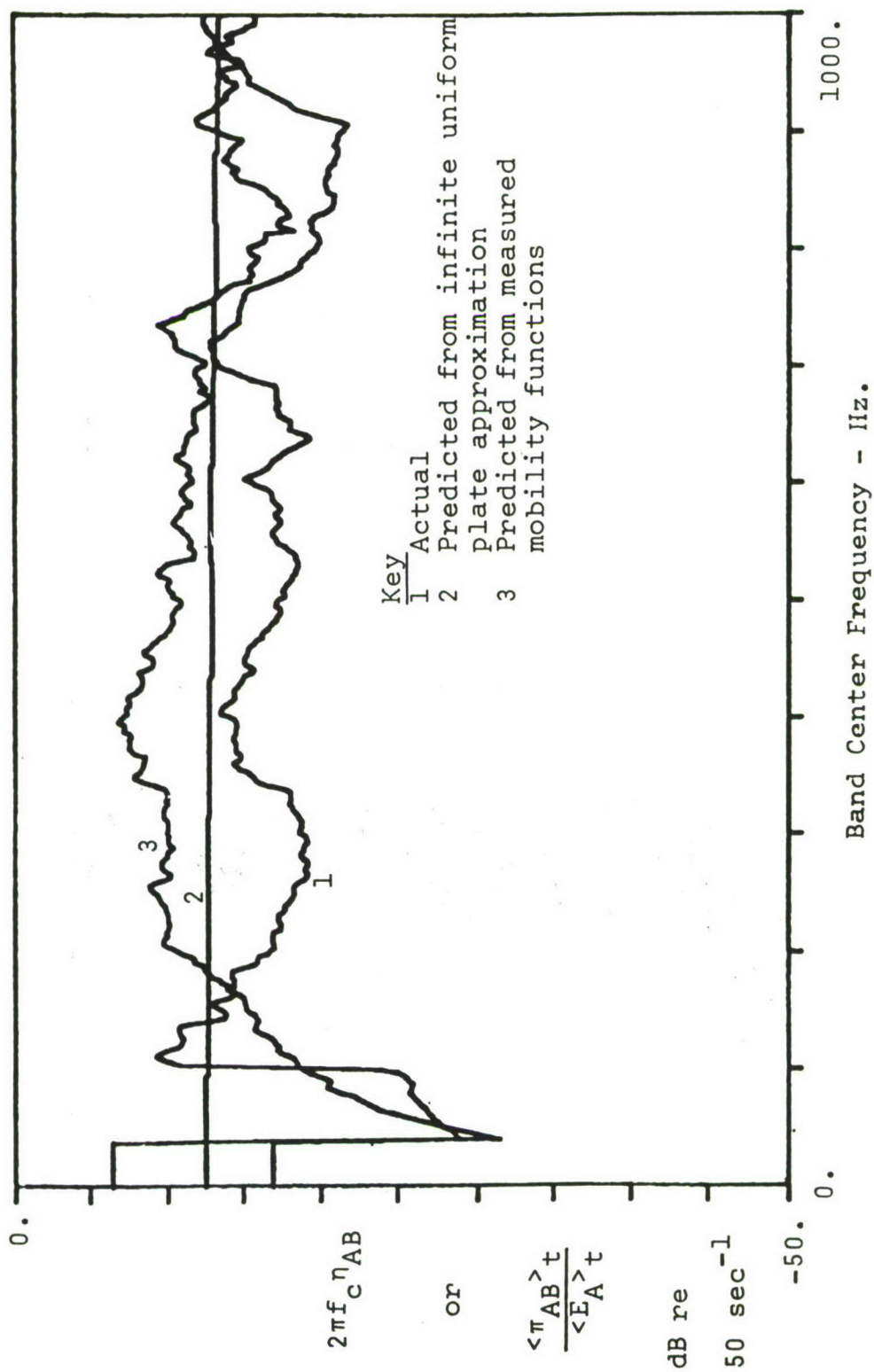


Figure 15 Normalized Transmitted Power $\langle \pi_{AB} \rangle_t / \langle E_A \rangle_t$

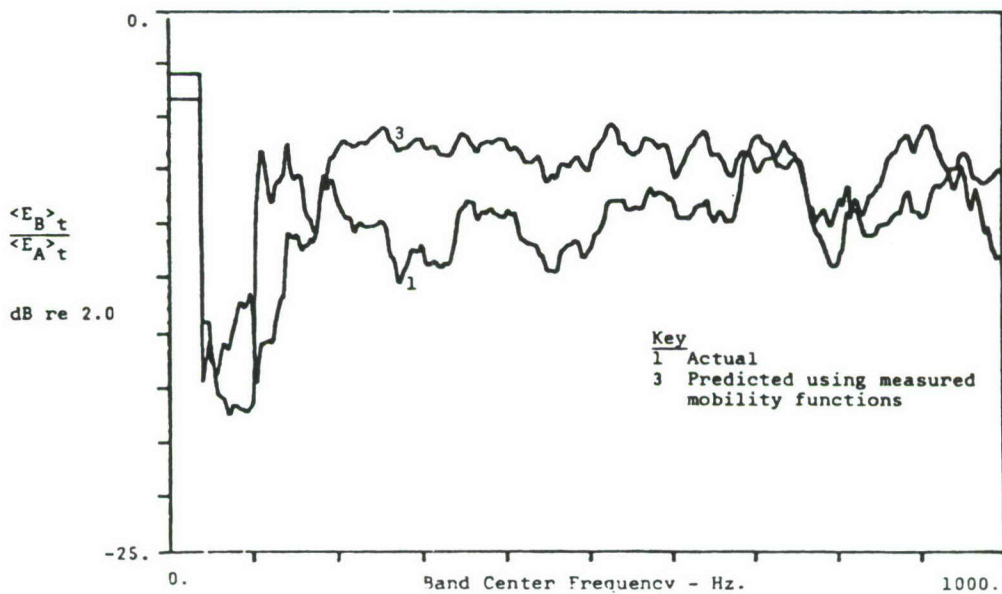
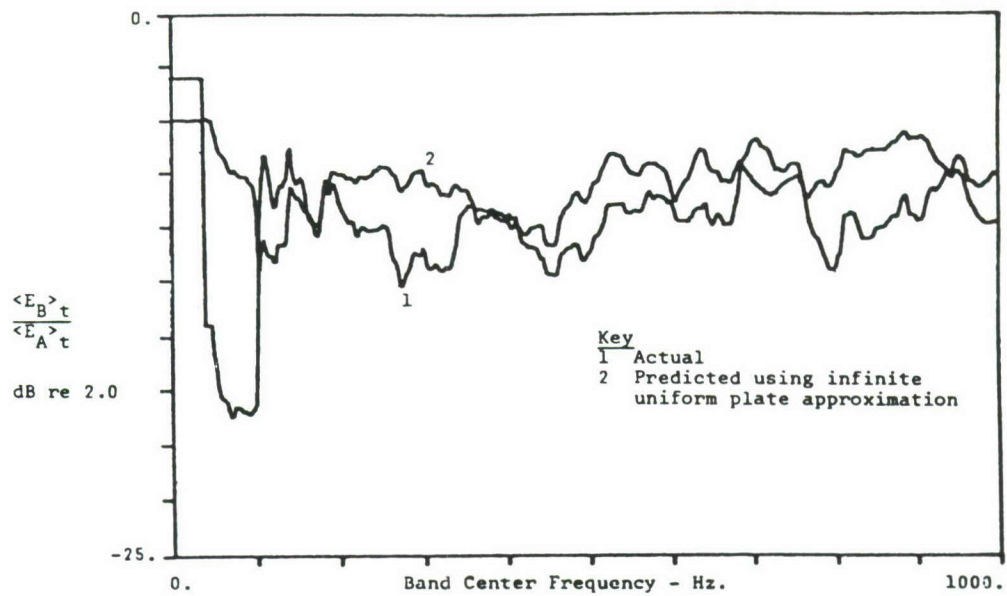


Figure 16 Normalized Energy of Indirectly Excited Body. Actual Damping Used.

spatially-averaged r.m.s. displacement or velocity would be within 41%.

The comparison of angular response, actual vs. estimated from measured energy, is given in terms of smoothed spectral density functions and r.m.s. values in Figure 17.

Predicted r.m.s. angular displacement is within 8% of actual. Most of the discrepancy appears to be due to a single mode around 250 Hz. which contributes disproportionately to mean square angle at the particular point chosen for comparison. Errors of this type are inherent in the SEA procedure if individual mode shapes are not known. This is discussed extensively in Chapter 4 of Lyon [13].

4.6.4.2 Measurement of Internal Loss Factor

Figure 18 shows a comparison of internal loss factor (weighted by band center frequency) as measured by two different methods. The continuous curve is $\langle \pi_{AB} \rangle_t / \langle E_B \rangle_t$ obtained by direct measurement of the numerator and denominator and is, by definition, the correct quantity for SEA modeling. $\langle E_B \rangle_t$ is obtained as a mass-weighted sum of velocity PSD's for forty points arranged over the plate. The symbol \ominus indicates $\omega_c \eta_B$ estimated by curve-fitting to the measured coupling point mobility and using Eq. (4.83). The width of the symbol indicates the frequency range over which $H_B(f)$ was measured using a high-resolution DFT. Natural frequencies and modal damping ratios as estimated by curve fitting are tabulated in Table 2 for the four bands shown.

Agreement between the two measurement methods is excellent. The 0 to 1 dB difference is better than typical repeatability of damping measurements. This result is quite encouraging because measurement by curve fitting was about ten times faster for this simple case. For a more complex structure, the advantage would be even greater due to the effort of assembling a mass matrix which is circumvented in the curve fitting method.

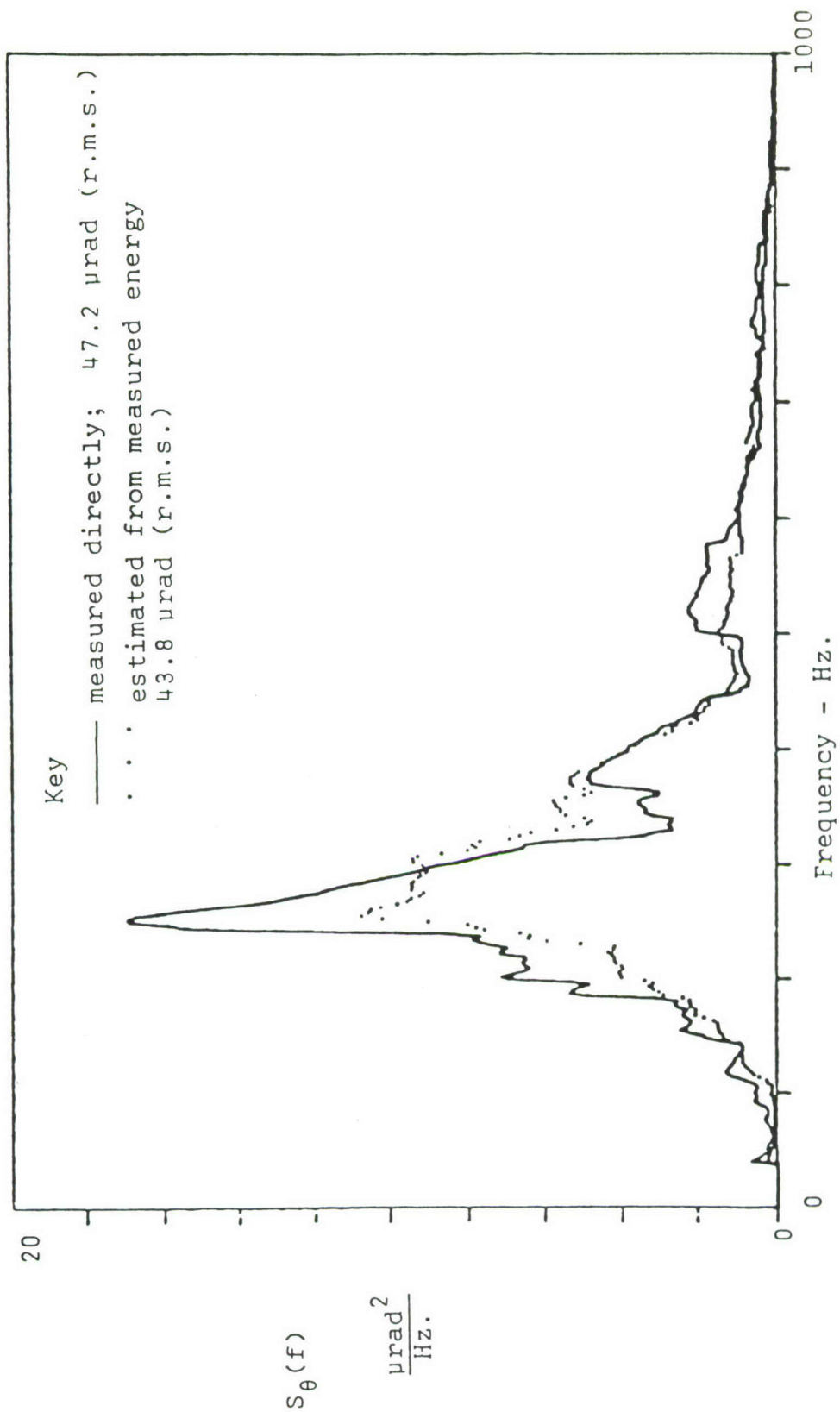


Figure 17 Smoothed Power Spectral Density and R.M.S. Value of High-Frequency Angular Vibration at a Typical Point of a Uniform Rectangular Plate

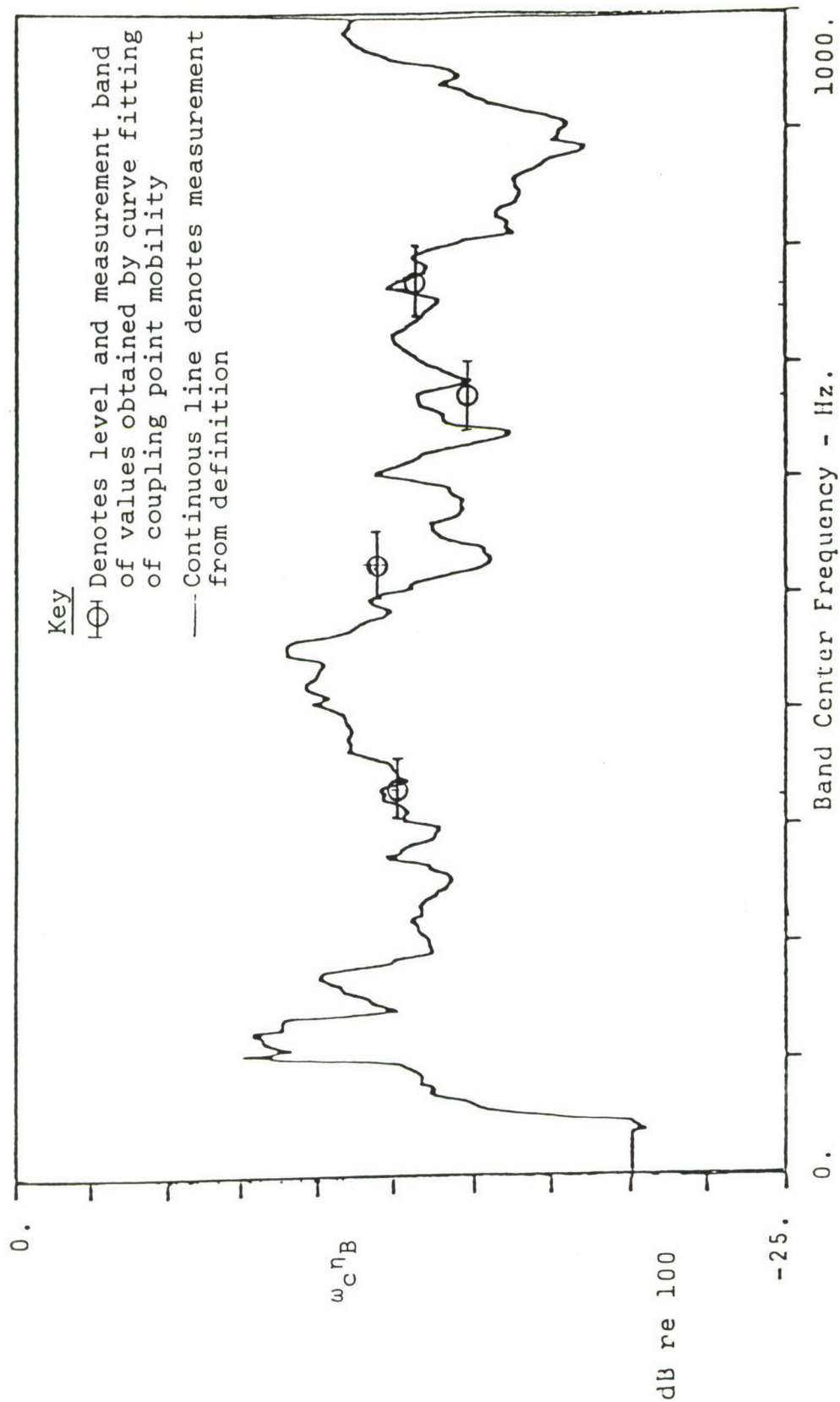


Figure 18 Internal Loss Factor vs. Frequency of Indirectly Excited Plate

TABLE 2

NATURAL FREQUENCIES
AND MODAL DAMPING RATIOS FOR PLATE B

Band 1 300-340 Hz		Band 2 505-540 Hz	
$f^{(i)}$, Hz	$\xi^{(i)}$	$f^{(i)}$, Hz	$\xi^{(i)}$
303.9	.00329	509.0	.00380
306.3	.00497	511.7	.00059
309.8	.00123	521.3	.00059
325.8	.00283	522.9	.00173
326.8	.00125	531.3	.00086
330.3	.00037	533.6	.00134
336.9	.00182	542.4	.00098
Band 3 650-700 Hz		Band 4 750-800 Hz	
$f^{(i)}$, Hz	$\xi^{(i)}$	$f^{(i)}$, Hz	$\xi^{(i)}$
654.3	.00092	756.5	.00035
657.3	.00062	758.6	.00095
659.4	.00106	760.1	.00040
663.5	.00096	766.4	.00035
669.5	.00106	770.6	.00031
676.2	.00007	772.0	.00061
679.9	.00064	775.9	.00043
682.3	.00086	784.3	.00046
686.5	.00032	789.8	.00099
689.2	.00070		

4.6.4.3 Other Results

Figure 19 indicates the probable effect of inadequate DFT resolution in measuring mobilities $H_A(f)$ and $H_B(f)$. The data shown is modulus of the acceleration admittance rather than real part of velocity admittance (mobility) but the implication is still quite clear. Decreasing the frequency resolution will suppress the peaks in $H(f)$ without filling in the valleys. The frequency averaged version of $H(f)$ will thus be reduced and $\omega\eta_{AB}$ predicted from it by Eq. (4.56) will be increased. In Figure 15 the predicted value of $\omega_c \eta_{AB}$ follows the actual but is higher over a broad range of frequencies. It thus appears that a smaller DFT resolution element would result in improved accuracy.

A closer examination of the measured $\bar{H}_B(f_c)$ function shows another peculiarity. It is plotted in Figure 20 in both real-imaginary and modulus formats. In theory, from Reference [14], the real part of $\bar{H}_B(f_c)$, called the conductance, should approach a constant with increasing f_c while the imaginary part, or susceptance, goes to zero. From Figure 20, the modulus approaches a fairly constant value but neither $\text{Re}[\bar{H}_B(f_c)]$ nor $\text{Im}[\bar{H}_B(f_c)]$ behave as predicted. Identical behavior was observed with $\bar{H}_A(f_c)$. It appears that some further work on measurement of averaged mobility of high modal density structures will be in order if wave transmission methods are to be pursued further.

The spectral density of transmitted power and its smoothed version are shown in Figure 21. Spectral densities of total vibration energy for the A and B plates are shown in Figure 22. Two observations may be made:

- (1) The distributions of energy and power with respect to frequency are reasonably even. In testing the SEA model with broadband excitation, one is essentially running several experiments simultaneously, one in each frequency band. Some care must be taken so that a large signal in one band does not cause error in another due to dynamic range limitations.

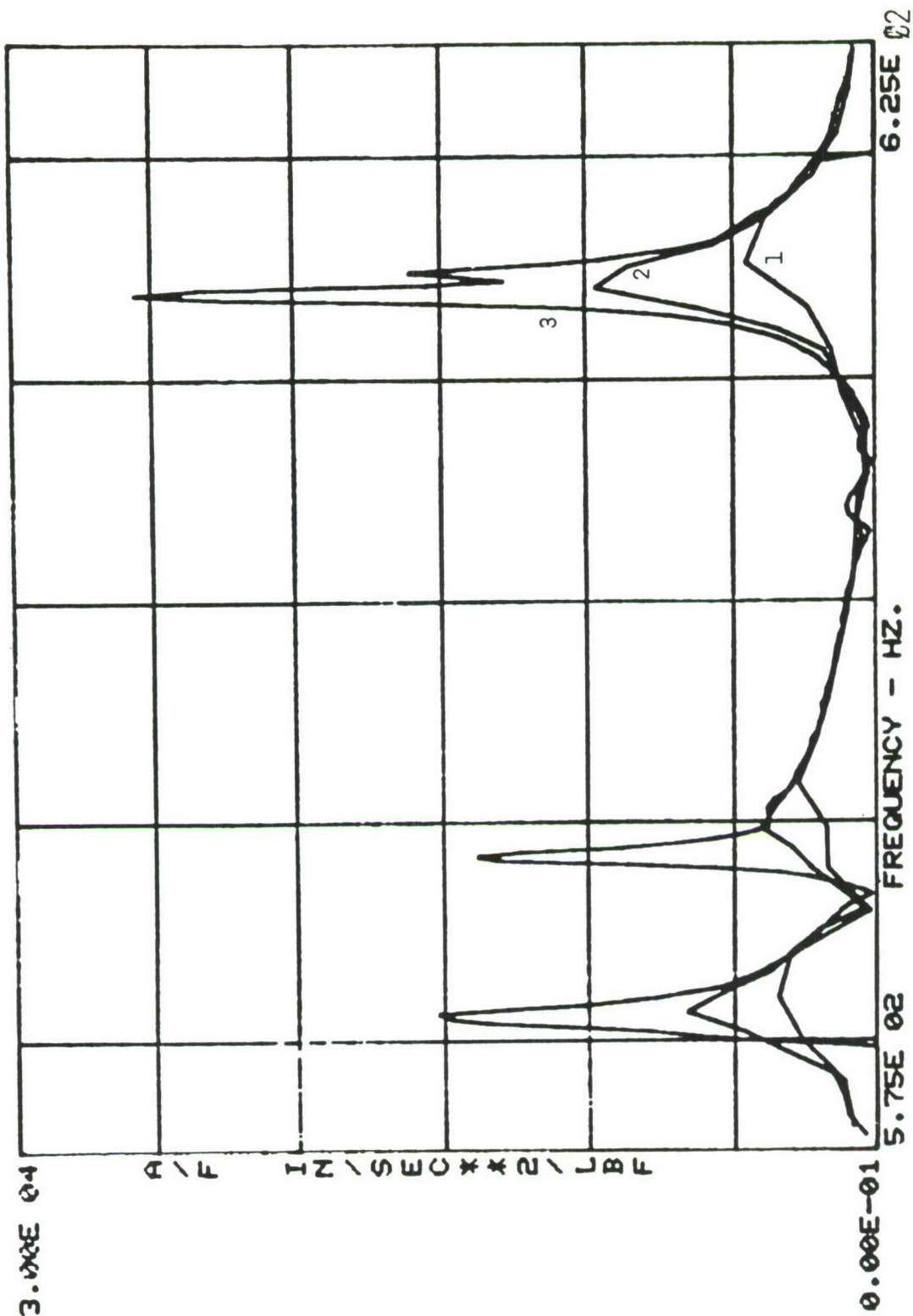
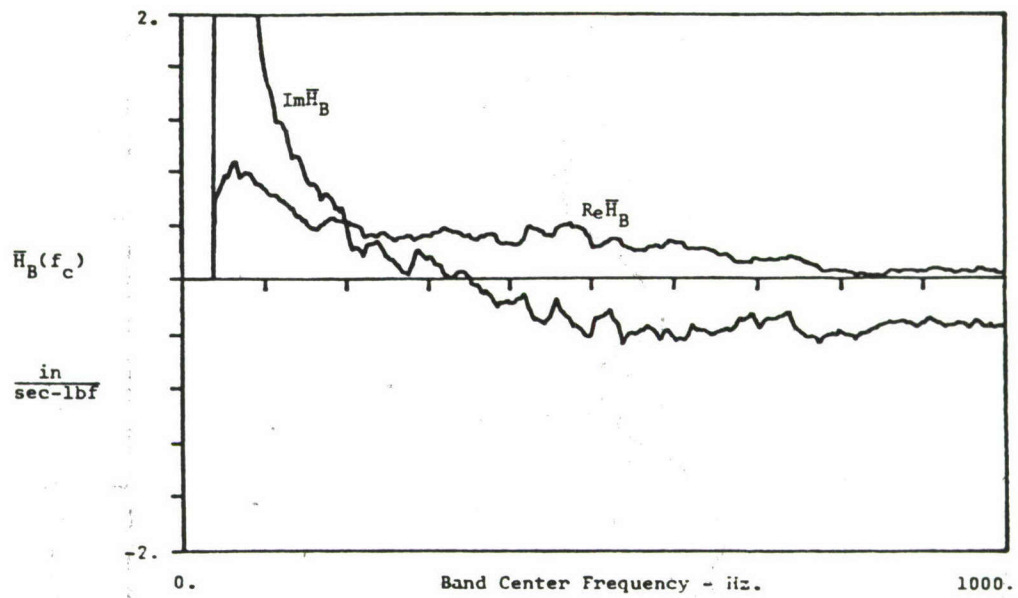
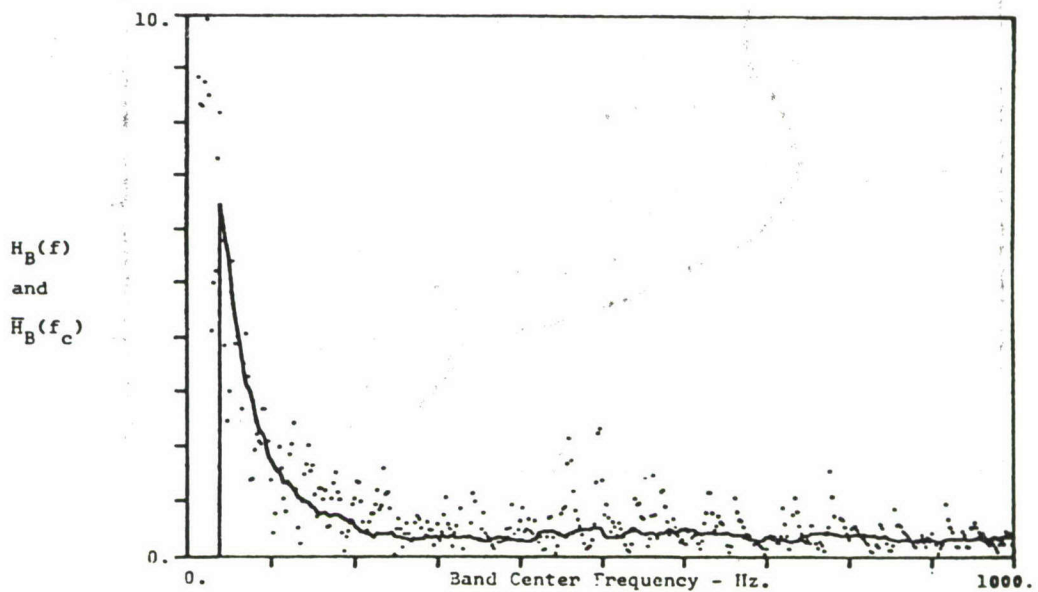


Figure 19 Modulus of Acceleration Admittance at Center Point of Large Plate.
Measured with Frequency Resolution of (1) 1.94 Hz.
(2) 0.94 Hz.
(3) 0.20 Hz.



(real and imaginary part)



(modulus)

Figure 20 Measured Coupling Point Mobility of Small Plate After Smoothing

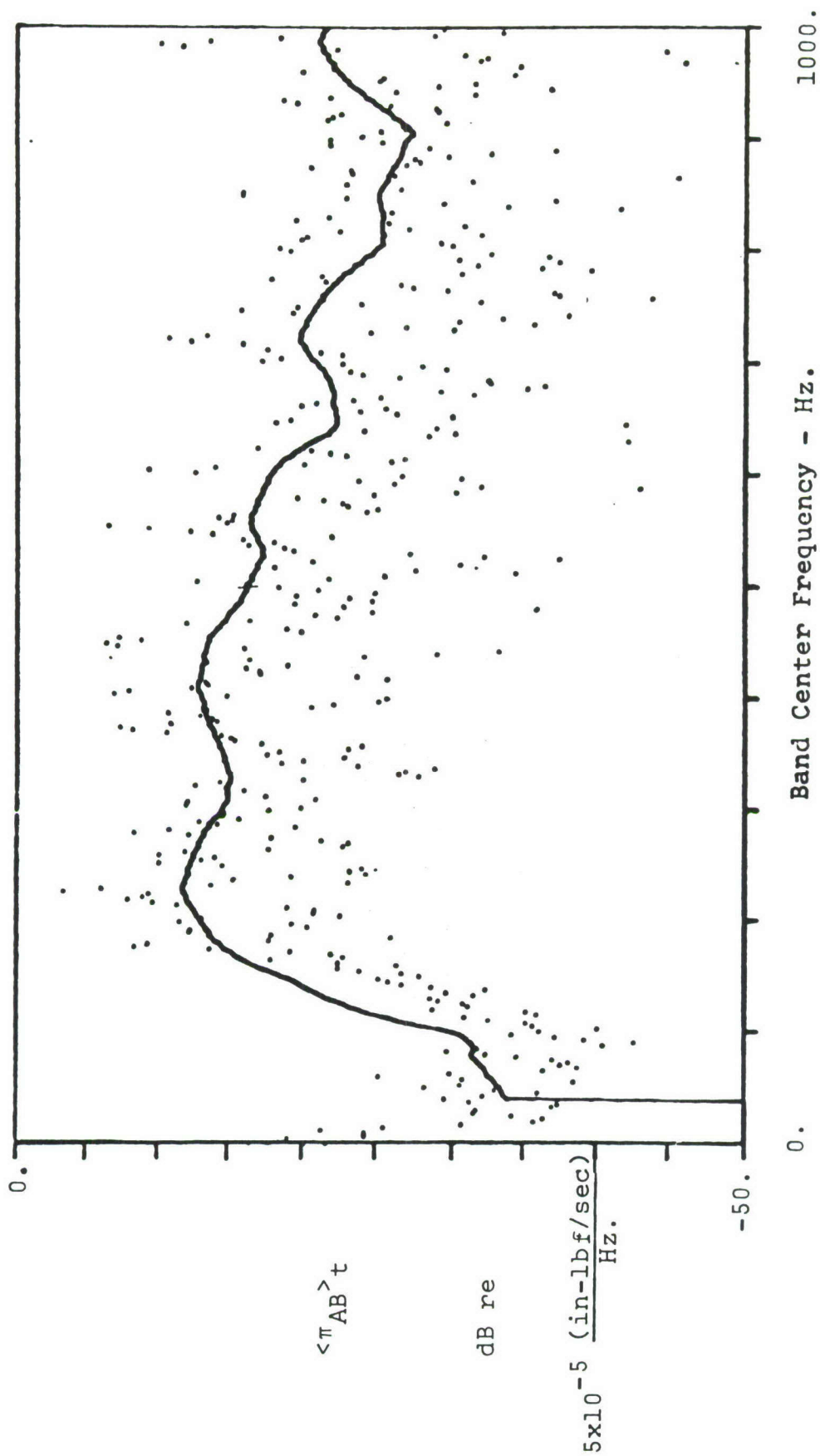


Figure 21 Spectral Density of Transmitted Power, Smoothed and Unsmoothed Versions

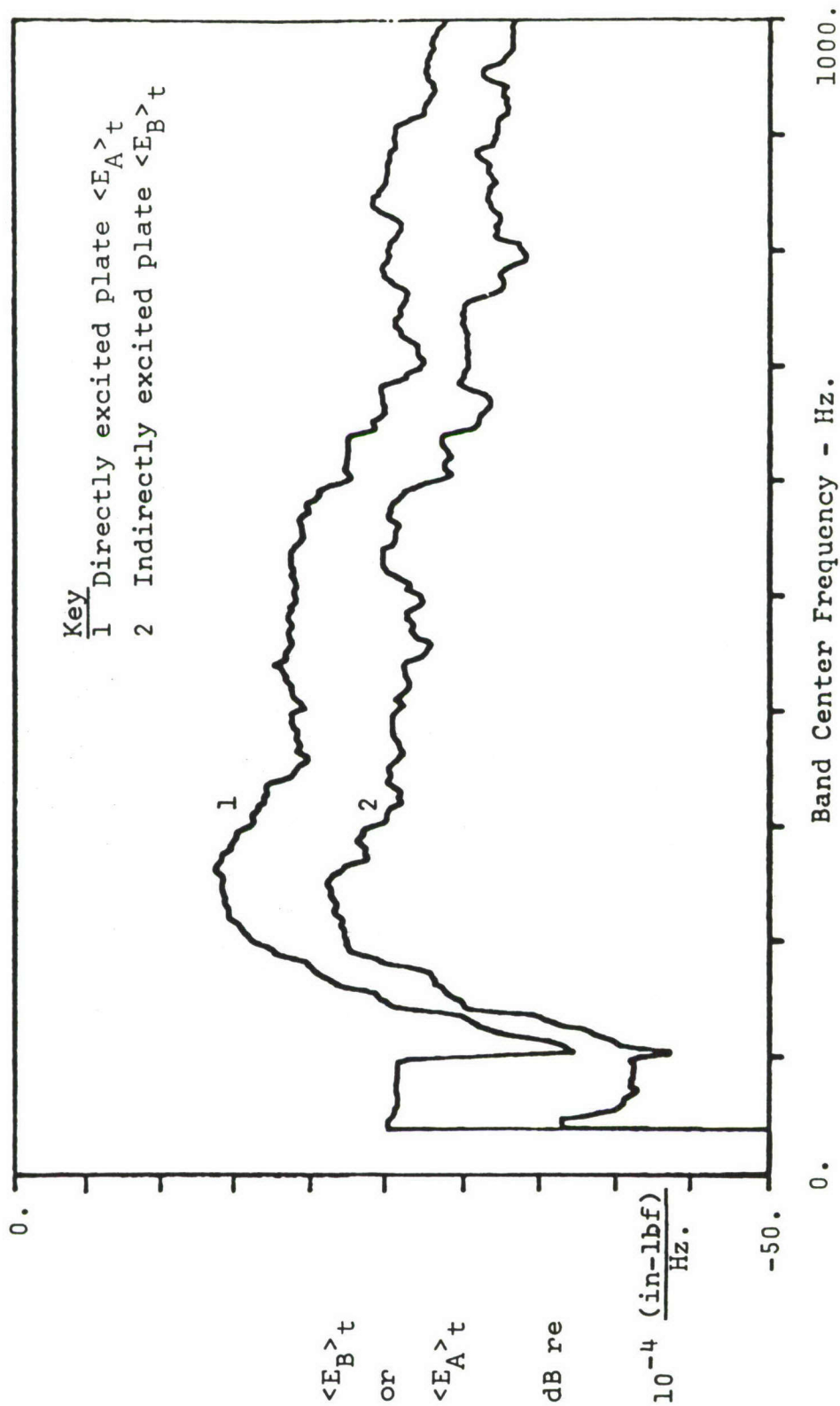


Figure 22 Smoothed Spectral Density of Total Vibration Energy.

This is particularly true with the fixed point arithmetic built into most minicomputer DFT software. Figures 21 and 22 show less than 40 dB of range, which should cause no problems.

- (2) The validity of the assumption $\langle E_A \rangle_t \gg \langle E_B \rangle_t$ may be checked using Figure 22. Over most of the range above 200 Hz the ratio of $\langle E_A \rangle_t / \langle E_B \rangle_t$ is in the range of 4 to 6. Therefore, the assumption is being violated but not grossly.

4.6.4.4 Results for Lightly Coupled Case

An unsuccessful attempt was made to predict coupling loss factors for light coupling between the two plates. A light fiberglass spring of ring shape with $k_c(\text{static}) = 163 \text{ lbf/in.}$ was substituted for the aluminum coupling rod. Figure 23 shows the results. The disparity between curves 1 and 2 is caused by the fact that the actual $H_c(f)$ is much higher than $i2\pi f/k_c$ for frequencies over 200 Hz. In effect, the ring mass is not negligible and the dynamics of the coupling element are not of the simple form assumed in deriving Eq. (4.56).

Curve 3 of Figure 23 resulted from trying to remedy the above problem in a simple way. The B structure was redefined as the series combination of the spring and smaller plate. $H_B(f)$ for this combination was measured in the usual way. The combination was then considered to be attached to plate A by a spring of infinite stiffness and, thus, $H_c = 0$. The poor accuracy of $\omega_c \eta_{AB}$ computed under these assumptions is shown by curve 3. It is probably due to the fact that the coupling point to the new B structure is no longer a "typical" point and, thus, another assumption built into Eq. (4.56) is violated.

Key

1 Actual

2 Predicted from measured plate mobilities and light coupling,
 $k_C = 163 \text{ lbf/in}$

3 Predicted using $H_A =$ measured large plate mobility, $H_B =$
 measured mobility of series combination of smaller plate and
 light spring, $H_C = 0$

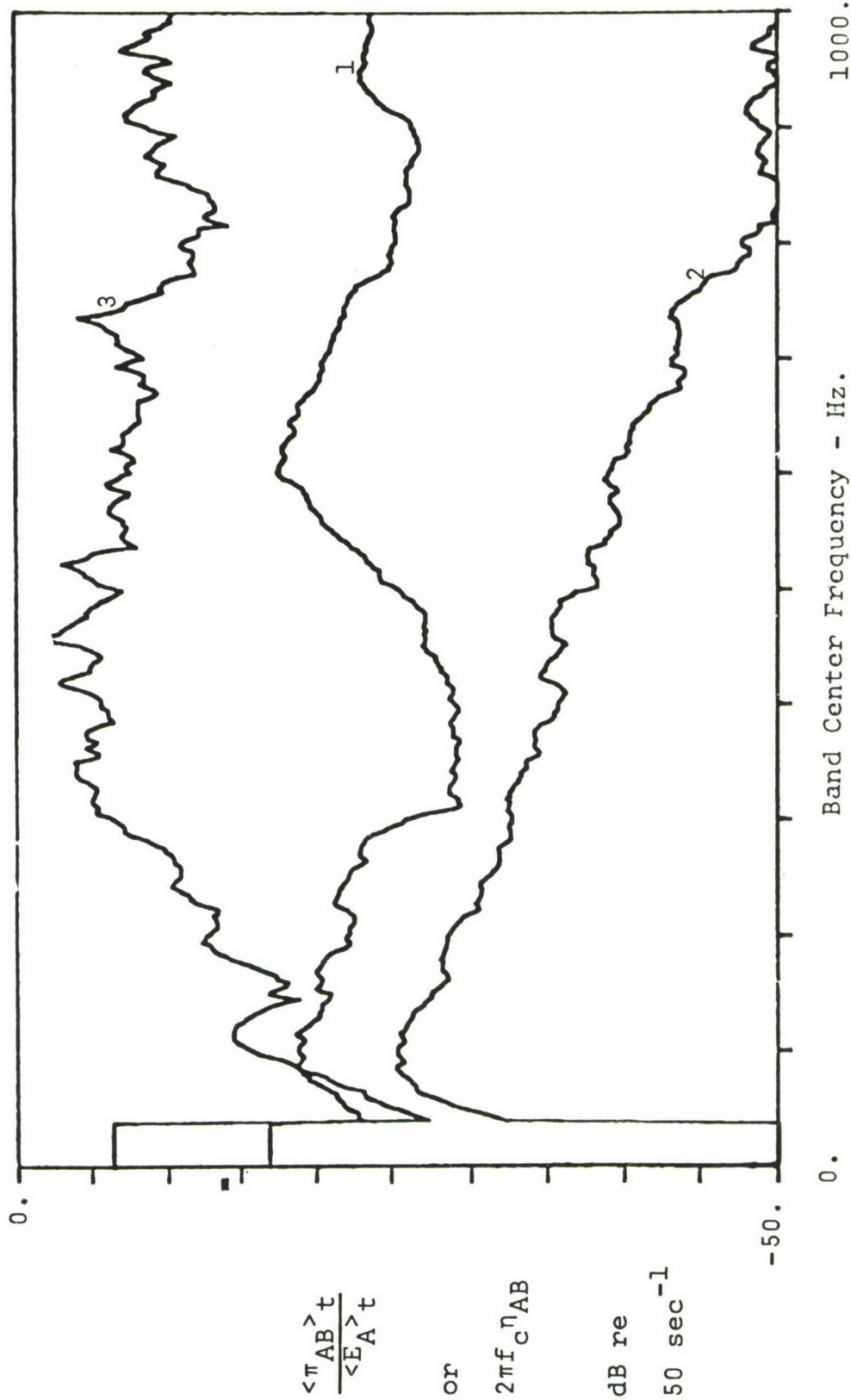


Figure 23 Normalized Transmitted Power With Light Coupling Spring

4.7 CONCLUSION

4.7.1 Summary

The motivation for study of high frequency methods in the current effort has been summarized and the method of Statistical Energy Analysis has been selected for further study. The basic principles of SEA have been reviewed within the context of analysis of airborne optical systems.

The wave transmission method for prediction of coupling loss factors has been reviewed and an alternate derivation of the principal result, without reference to waves, has been presented. It has been applied to a simple system involving two coupled plates of high modal density. Coupling loss factors have been predicted for the case of one plate only receiving direct excitation. Predictions have been made on the basis of modeling each plate either as an infinite uniform extension of the real plate or simply as an infinite uniform structure equivalent to the actual in the sense of having the same frequency-averaged mobility at the coupling point. The basic assumptions of SEA have been used to derive a simple formula for inferring r.m.s. angular response from component energy for the test case of a uniform plate. They have also been used to estimate a global internal loss factor from a single coupling point mobility measurement.

An experiment has been carried out to test various theoretical predictions described above for the case of two coupled uniform plates. Agreement with respect to component energy and transmitted power was generally encouraging, although it is suspected that some improvement in the test procedure itself may be possible. Prediction of angular response from the energy resultant was quite accurate, although it must be noted that this was essentially a demonstration case since the method used is restricted to uniform beams or plates. It was shown that excellent accuracy is possible in obtaining internal loss factor

via curve-fitting of the measured coupling-point mobility function.

At this point it is appropriate to review what has and has not been accomplished to date by the work on high frequency angular vibration. It would not be accurate to say that a first generation method has been developed for prediction of high frequency vibration, angular or otherwise. That goal, in retrospect, was simply unrealistic in view of the complexity and size of the problem relative to the resources allocated to it under this contract. What has been accomplished is, nevertheless, necessary and useful in the pursuit of that end. It has been demonstrated that angular vibration, in the form of a coarsely resolved power spectrum and associated r.m.s. value, can be predicted in a high frequency situation; i.e., one where deterministic modeling would be impractical. Wave transmission, one of the simplest methods for obtaining coupling loss factor for single point connections, has been recast in a format suitable for structures with large but finite modal densities. The new interpretation does not require modeling explicitly in terms of traveling waves and should be easier to relate to practical structures. Finally, data acquisition and processing software has been developed and used for tasks which have a high probability of occurring in the development of more general SEA methods. The usefulness of interactive digital signal processing technology for SEA modeling has thus been demonstrated. On the whole, it is the opinion of the investigators that the original decision to pursue SEA as the most promising method for high frequency prediction has been reinforced by this work.

4.7.2 Suggestions For Future Work

The work performed to date on high frequency methods has been quite useful in identifying specific areas where additional effort could be cost-effective. The ultimate goal is the

development of methods which are practical, though based on a solid theoretical foundation, for prediction of high frequency vibration of airborne optical systems. To this end, the following specific areas are suggested:

- (1) It may be possible to model individual normal modes as SEA components. This seems natural when an optical system with known normal modes is connected to a complex airframe with modes which are numerous and not individually known.
- (2) The basic idea of the wave transmission method might be extended to cover the case where components are connected at more than one degree of freedom. An admittance matrix formulation seems to be called for. It may be anticipated that this effort would be quite software-intensive.
- (3) It might be possible to incorporate finite element results into an SEA model. An interface format suggested by work to date is the mechanical admittance function averaged with respect to frequency [22]. It is suspected on theoretical grounds that this quantity can be predicted by a finite element model which is too coarse to accurately predict individual normal modes.
- (4) The coupled plate experiment suggested numerous improvements in the acquisition and processing of experimental data to obtain coupling loss factors. Several avenues were identified but not pursued due to time restrictions.

SECTION V
RELATIONSHIPS BETWEEN LINEAR AND ANGULAR
VIBRATION IN AIRCRAFT STRUCTURAL COMPONENTS

Because of the predominant availability of linear vibration data as compared to that for angular vibration of structural components, the task of finding a usable and correct relationship between these two vibration quantities appears to be justifiable on the basis that we could use linear vibration data to determine angular vibrations for those portions of the aircraft structure which are of interest. It seemed like an appropriate way to attack the problem would be to investigate the behavior of the simplest two degree-of-freedom (DOF) model of a structural component, which would be a spring-supported rigid bar, subjected to a temporally random concentrated load. In this way one could begin to obtain a feel for how a linear-to-angular relationship would depend upon the load and structural parameters (i.e. mass, stiffness, geometry) for a simple system. The next step was to proceed to more complex structural components and to determine linear-to-angular relationships for them as a function of position on the structure. This was done for a beam with simple-support boundary conditions, a plate with simple-support boundaries, and a beam with free-free boundary conditions all subjected to temporally random, spatially deterministic loading conditions. Some of the results of this study are presented here. However, the reader is referred to Reference [19] for the detailed results. In order to examine a structural form somewhat more representative of an aircraft structural component, a NASTRAN analysis of a stiffened curved panel subjected to a concentrated temporally random load was completed to round out the effort.

5.1 LINEAR-TO-ANGULAR RELATIONSHIP OF SPRING-SUPPORTED RIGID BAR SUBJECTED TO TEMPORALLY RANDOM CONCENTRATED LOAD

The simplest structural system possessing both linear and angular degrees of freedom is a spring-supported rigid bar as pictured below in Figure 24.

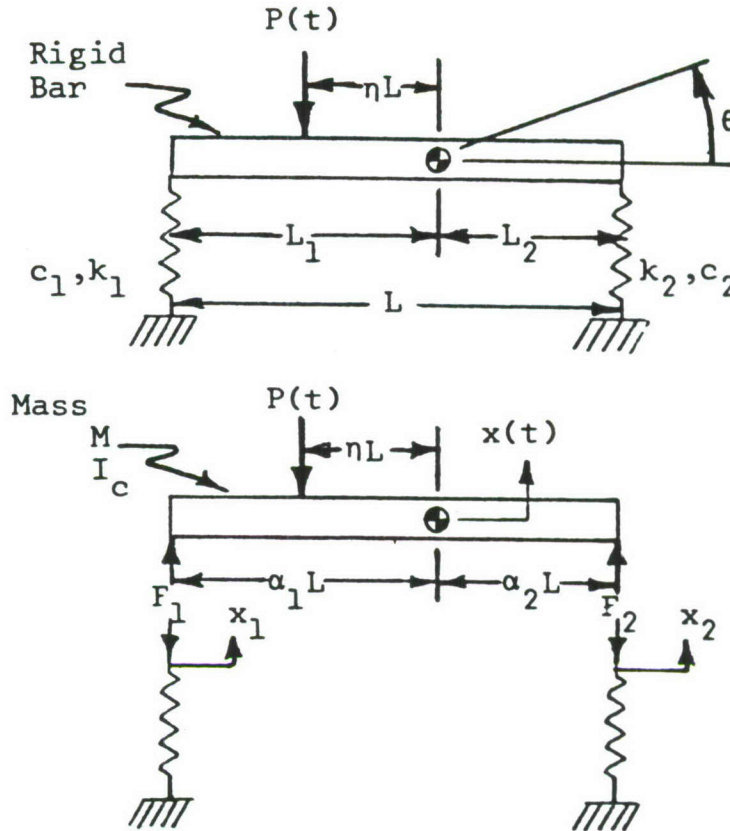


Figure 24 Spring Supported Rigid Bar

The spring and damping constants are defined as follows:

$$\begin{aligned} k_1 &= k ; c_1 = c \\ k_2 &= \gamma k ; c_2 = \beta c \end{aligned} \quad (5.1)$$

where $0 < \beta$ and $0 < \gamma$. The assumptions made in the subsequent analysis of this structure are (1) no gravity forces are present, and (2) angular displacements are small such that $\sin \theta \approx \theta$.

The differential equations of motion governing this system may be written in matrix form as

$$\begin{bmatrix} M & 0 \\ 0 & \frac{I_c}{L^2} \end{bmatrix} \begin{bmatrix} \ddot{x} \\ \ddot{y} \end{bmatrix} + c \begin{bmatrix} (1+\beta) & (\beta\alpha_2 - \alpha_1) \\ (\beta\alpha_2 - \alpha_1)(\alpha_1^2 + \beta\alpha_2^2) & \end{bmatrix} \begin{bmatrix} \dot{x} \\ \dot{y} \end{bmatrix} + k \begin{bmatrix} (1+\gamma) & (\gamma\alpha_2 - \alpha_1) \\ (\gamma\alpha_2 - \alpha_1)(\alpha_1^2 + \gamma\alpha_2^2) & \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix} = \begin{bmatrix} -P \\ -P\eta \end{bmatrix} \quad (5.2)$$

where $y = L\theta$. In order to facilitate the solution of Eq. (5.2) one may employ the linear coordinate transformation from physical to normal coordinates

$$\begin{bmatrix} x \\ y \end{bmatrix} = \begin{bmatrix} u_{11} & u_{12} \\ u_{21} & u_{22} \end{bmatrix} \begin{bmatrix} q_1 \\ q_2 \end{bmatrix} \quad (5.3)$$

which will uncouple the system. In normal coordinates, the set of D.E.'s Eq. (5.2) will take on the form

$$\begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} \ddot{q}_1 \\ \ddot{q}_2 \end{bmatrix} + \begin{bmatrix} C_1 & 0 \\ 0 & C_2 \end{bmatrix} \begin{bmatrix} \dot{q}_1 \\ \dot{q}_2 \end{bmatrix} + \begin{bmatrix} K_1 & 0 \\ 0 & K_2 \end{bmatrix} \begin{bmatrix} q_1 \\ q_2 \end{bmatrix} = \begin{bmatrix} u_{11} & u_{21} \\ u_{12} & u_{22} \end{bmatrix} \begin{bmatrix} -P \\ -P\eta \end{bmatrix} \quad (5.4)$$

If the following substitutions are made

$$2\tilde{\beta}_1 = C_1; \quad 2\tilde{\beta}_2 = C_2; \quad \text{and} \quad \tilde{\omega}_1^2 = K_1; \quad \tilde{\omega}_2^2 = K_2 \quad (5.5a)$$

and

$$a_1 = -(u_{11} + \eta u_{21}); \quad a_2 = -(u_{12} + \eta u_{22}) \quad (5.5b)$$

Equation (5.4) becomes

$$\ddot{q}_1 + 2\tilde{\beta}_1 \dot{q}_1 + \tilde{\omega}_1^2 q_1 = a_1 P(t) \quad (5.6)$$

$$\ddot{q}_2 + 2\tilde{\beta}_2 \dot{q}_2 + \tilde{\omega}_2^2 q_2 = a_2 P(t) \quad (5.7)$$

for which the solutions are

$$q_1(t) = a_1 \int_0^{\infty} h_1(\phi) P(t-\phi) d\phi \quad (5.8)$$

$$q_2(t) = a_2 \int_0^{\infty} h_2(\phi) P(t-\phi) d\phi \quad (5.9)$$

where

$$h_1(t) = \frac{e^{-\tilde{\beta}_1 t}}{\tilde{\omega}_1'} \sin \tilde{\omega}_1' t; \quad h_2(t) = \frac{e^{-\tilde{\beta}_2 t}}{\tilde{\omega}_2'} \sin \tilde{\omega}_2' t$$

$$\text{and } \tilde{\omega}_1' = \sqrt{\tilde{\omega}_1^2 - \tilde{\beta}_1^2} \quad ; \quad \tilde{\omega}_2' = \sqrt{\tilde{\omega}_2^2 - \tilde{\beta}_2^2} \quad (5.10)$$

The next question to answer is what form x and y will take if $P(t)$ is a random function (stationary and ergodic). At this point one might ask about making $\eta = \eta(t)$ thereby allowing the load to be random as to its point of application as well as its magnitude; however, this will not be done here.

Working with $x(t)$ first and substituting Eqs. (5.8) and (5.9) in Eq. (5.3) gives

$$x(t) = \int_0^{\infty} [u_{11} a_1 h_1(\phi) + u_{12} a_2 h_2(\phi)] P(t-\phi) d\phi \quad (5.11)$$

Since $P(t)$ is random and is usually assumed to have a Gaussian distribution about its mean value, $x(t)$ will also be Gaussian, if our structural system is linear. One can calculate the auto correlation of $x(t)$ then to be

$$\begin{aligned} E[x(t)x(t+\tau)] &= E\left[\int_0^{\infty} \int_0^{\infty} \{u_{11} a_1 h_1(\phi_1) + u_{12} a_2 h_2(\phi_1)\} \cdot \right. \\ &\quad \left. \cdot \{u_{11} a_1 h_1(\phi_2) + u_{12} a_2 h_2(\phi_2)\} P(t-\phi_1) P(t+\tau-\phi_2) d\phi_1 d\phi_2\right] \end{aligned} \quad (5.12)$$

which can be written

$$\begin{aligned}
 R_x(\tau) = & \int_0^\infty \int_0^\infty R_p(\tau + \phi_1 - \phi_2) \{ u_{11}^2 a_1^2 h_1(\phi_1) h_1(\phi_2) \\
 & + u_{11} a_1 u_{12} a_2 [h_1(\phi_1) h_2(\phi_2) + h_1(\phi_2) h_2(\phi_1)] \\
 & + u_{12}^2 a_2^2 h_2(\phi_1) h_2(\phi_2) \} d\phi_1 d\phi_2
 \end{aligned} \tag{5.13}$$

In terms of spectral densities, the Wiener-Khintchine relation states that

$$S_x(\omega) = \frac{1}{2\pi} \int_{-\infty}^{\infty} R_x(\tau) e^{-i\omega\tau} d\tau \tag{5.14}$$

Substituting Eq. (5.13) for $R_x(\tau)$ into Eq. (5.14) and integrating yields

$$\begin{aligned}
 S_x(\omega) = & [u_{11}^2 a_1^2 |H_1(\omega)|^2 + u_{12}^2 a_2^2 |H_2(\omega)|^2] S_p(\omega) \\
 & + u_{11} a_1 u_{12} a_2 [H_1(\omega) H_2^*(\omega) + H_2(\omega) H_1^*(\omega)] S_p(\omega)
 \end{aligned} \tag{5.15}$$

where

$$H_1(\omega) = \frac{1}{\tilde{\omega}_1^2 - \omega^2 + 2i\tilde{\beta}_1\omega} \quad ; \quad H_2(\omega) = \frac{1}{\tilde{\omega}_2^2 - \omega^2 + 2i\tilde{\beta}_2\omega} \tag{5.16}$$

and $H_1^*(\omega)$ and $H_2^*(\omega)$ are the complex conjugates of $H_1(\omega)$ and $H_2(\omega)$. Similarly, one would find

$$\begin{aligned}
 S_y(\omega) = & [u_{21}^2 a_1^2 |H_1(\omega)|^2 + u_{22}^2 a_2^2 |H_2(\omega)|^2] S_p(\omega) \\
 & + u_{21} a_1 u_{22} a_2 [H_1(\omega) H_2^*(\omega) + H_2(\omega) H_1^*(\omega)] S_p(\omega)
 \end{aligned} \tag{5.17}$$

For white noise forcing with constant spectral density,

$S_p(\omega) = S_{po}$, the next step would be to determine the mean square responses

$$\overline{x^2} = \int_{-\infty}^{\infty} S_x(\omega) d\omega \quad (5.18)$$

$$\overline{y^2} = \int_{-\infty}^{\infty} S_y(\omega) d\omega \quad (5.19)$$

If one lets

$$\begin{aligned} \hat{H}_{11} &= \int_{-\infty}^{\infty} |H_1(\omega)|^2 d\omega \quad ; \quad \hat{H}_{12} = \int_{-\infty}^{\infty} H_1(\omega) H_2^*(\omega) d\omega \\ \hat{H}_{21} &= \int_{-\infty}^{\infty} H_2(\omega) H_1^*(\omega) d\omega; \quad \hat{H}_{22} = \int_{-\infty}^{\infty} |H_2(\omega)|^2 d\omega \end{aligned} \quad (5.20)$$

Equations (5.18) and (5.19) become

$$\overline{x^2} = [u_{11}^2 a_1^2 \hat{H}_{11} + u_{11} a_1 u_{12} a_2 (\hat{H}_{12} + \hat{H}_{21}) + u_{12}^2 a_2^2 \hat{H}_{22}] S_{po} \quad (5.21)$$

$$\overline{y^2} = [u_{21}^2 a_1^2 \hat{H}_{11} + u_{21} a_1 u_{22} a_2 (\hat{H}_{12} + \hat{H}_{21}) + u_{22}^2 a_2^2 \hat{H}_{22}] S_{po} \quad (5.22)$$

Taking the ratio of mean square angular to mean square linear displacement, one obtains

$$\frac{y_{rms}}{x_{rms}} = \sqrt{\frac{u_{21}^2 a_1^2 \hat{H}_{11} + u_{21} u_{22} a_1 a_2 (\hat{H}_{12} + \hat{H}_{21}) + u_{22}^2 a_2^2 \hat{H}_{22}}{u_{11}^2 a_1^2 \hat{H}_{11} + u_{11} u_{12} a_1 a_2 (\hat{H}_{12} + \hat{H}_{21}) + u_{12}^2 a_2^2 \hat{H}_{22}}} \quad (5.23)$$

where

$$\hat{H}_{11} = \pi/C_1 K_1 \quad ; \quad \hat{H}_{22} = \pi/C_2 K_2 \quad (5.24)$$

$$\hat{H}_{12} = 2\pi(C_1 + C_2) / \{ [K_1 - K_2 - \frac{1}{2}C_1(C_1 + C_2)]^2 + (C_1 + C_2)^2 (K_1 - \frac{C_1}{4}) \} \quad (5.25)$$

$$\hat{H}_{21} = 2\pi(C_1 + C_2) / \{ [K_2 - K_1 - \frac{1}{2}C_2(C_1 + C_2)]^2 + (C_1 + C_2)^2 (K_2 - \frac{C_2}{4}) \} \quad (5.26)$$

Calculating the coefficients of \hat{H}_{ij} in terms of M_1 and M_2 , one arrives at

$$u_{21}^2 a_1^2 = \frac{1}{4M_2} \left[\left(\frac{1}{M_1} - \frac{1}{M_2} \right) \frac{1}{1+a^2} + 2 \left(\frac{\eta}{\sqrt{M_1 M_2}} \sqrt{\frac{1}{1+a^2}} + \frac{1}{M_2} \right) \left(1 - \sqrt{1 - \frac{1}{1+a^2}} \right) \right] \quad (5.27)$$

$$u_{21} u_{22} a_1 a_2 = \frac{1}{4M_2} \left[\left(\frac{1}{M_1} + \frac{\eta^2}{M_2} \right) \frac{1}{1+a^2} + \frac{2\eta}{M_1 M_2} \sqrt{\frac{1}{1+a^2}} \sqrt{1 - \frac{1}{1+a^2}} \right] \quad (5.28)$$

$$u_{22}^2 a_2^2 = \frac{1}{4M_2} \left[\left(\frac{1}{M_1} - \frac{1}{M_2} \right) \frac{1}{1+a^2} - 2 \left(\frac{\eta}{\sqrt{M_1 M_2}} \sqrt{\frac{1}{1+a^2}} - \frac{1}{M_2} \right) \left(1 + \sqrt{1 - \frac{1}{1+a^2}} \right) \right] \quad (5.29)$$

$$u_{11}^2 a_1^2 = \frac{1}{4M_1} \left[\left(\frac{1}{M_2} - \frac{1}{M_1} \right) \frac{1}{1+a^2} + 2 \left(\frac{\eta}{\sqrt{M_1 M_2}} \sqrt{\frac{1}{1+a^2}} + \frac{1}{M_1} \right) \left(1 + \sqrt{1 - \frac{1}{1+a^2}} \right) \right] \quad (5.30)$$

$$u_{11} u_{12} a_1 a_2 = - \frac{1}{4M_1} \left[\left(\frac{1}{M_1} + \frac{\eta^2}{M_2} \right) \frac{1}{1+a^2} + \frac{2\eta}{\sqrt{M_1 M_2}} \sqrt{1 - \frac{1}{1+a^2}} \right] \quad (5.31)$$

$$u_{12}^2 a_2^2 = \frac{1}{4M_1} \left[\left(\frac{1}{M_2} - \frac{1}{M_1} \right) \frac{1}{1+a^2} - 2 \left(\frac{\eta}{\sqrt{M_1 M_2}} \sqrt{\frac{1}{1+a^2}} - \frac{1}{M_1} \right) \left(1 - \sqrt{1 - \frac{1}{1+a^2}} \right) \right] \quad (5.32)$$

where

$$a = \frac{1}{2(\gamma\alpha_2 - \alpha_1)} (\alpha_1^2 + \gamma\alpha_2^2 - \frac{M_2}{M_1} (1 + \gamma)) \quad (5.33)$$

and $M_1 = M$; $M_2 = I_c/L^2$.

If one assumes that the damping is proportional, then

$$C_1 = \frac{\Lambda}{K} K_1 \quad \text{and} \quad C_2 = \frac{\Lambda}{K} K_2 \quad (\Lambda \text{ a constant})$$

where

$$K_1 = \frac{\tilde{k}_1}{2M_1} \left(1 + \sqrt{1 - \frac{1}{1+a^2}}\right) + \frac{\tilde{k}_2}{\sqrt{M_1 M_2}} \sqrt{\frac{1}{1+a^2}} + \frac{\tilde{k}_3}{2M_2} \left(1 - \sqrt{1 - \frac{1}{1+a^2}}\right) \quad (5.34)$$

$$K_2 = \frac{\tilde{k}_1}{2M_1} \left(1 - \sqrt{1 - \frac{1}{1+a^2}}\right) - \frac{\tilde{k}_2}{\sqrt{M_1 M_2}} \sqrt{\frac{1}{1+a^2}} + \frac{\tilde{k}_3}{2M_2} \left(1 + \sqrt{1 - \frac{1}{1+a^2}}\right) \quad (5.35)$$

and

$$\begin{aligned} \tilde{k}_1 &= (1+\gamma)k \\ \tilde{k}_2 &= (\gamma\alpha_2 - \alpha_1)k \\ \tilde{k}_3 &= (\alpha_1^2 + \gamma\alpha_2^2)k \end{aligned} \quad (5.36)$$

There is a special case which one might examine. When $\gamma = \frac{\alpha_1}{\alpha_2}$, which implies that $\beta = \frac{\alpha_1}{\alpha_2}$ for proportional damping, the DE Eq. (5.2) automatically uncouple; and the elements of the transformation matrix reduce to:

$$u_{11} = \sqrt{1/M_1} \quad ; \quad u_{22} = \sqrt{1/M_2} \quad ; \quad u_{12} = u_{21} = 0 \quad (5.37)$$

Also

$$a_1 = -u_{11} \quad ; \quad a_2 = -\eta u_{22} \quad (5.38)$$

$$\hat{H}_{11} = \pi/C_1 K_1 = (\pi/\Lambda k) (1 + \alpha_1/\alpha_2)^2 / M^2 \quad (5.39)$$

$$\hat{H}_{22} = \pi/C_2 K_2 = (\pi/\Lambda k) (\alpha_1^2 + \alpha_1\alpha_2)^2 L^4 / I_c^2 \quad (5.40)$$

Equation (5.23) then reduces to

$$\frac{y_{rms}}{x_{rms}} = \frac{\eta M L^2}{I_c} \sqrt{\frac{I_c^2 (1 + \alpha_1/\alpha_2)^2}{(\alpha_1^2 + \alpha_1\alpha_2)^2 L^4 M^2}} = \frac{\eta (1 + \alpha_1/\alpha_2)}{\alpha_1 \alpha_2 (1 + \alpha_1/\alpha_2)} \quad (5.41)$$

or

$$\frac{y_{rms}}{x_{rms}} = \frac{\eta}{\alpha_1 \alpha_2} \quad ; \quad y_{rms} = L \theta_{rms} \quad (5.42)$$

for the condition $\gamma = \frac{\alpha_1}{\alpha_2}$, which physically represents that case where the spring and damping forces at each end of the bar exactly match the translational inertia reactions at those points.

One might observe that y_{rms}/x_{rms} may vary between zero and a finite number depending upon the location of P. Equation (5.42) also portends a rather significant variance depending upon the values for α_1 and α_2 . The occurrence of such large variances would suggest that spatially averaged values for random locations of P(t) are subject to question.

5.2 BEHAVIOR OF BEAMS SUBJECTED TO TEMPORALLY RANDOM LOADING CONDITIONS

Lee and Whaley had already laid some of the ground work for beams in their paper [23]. One of the major analytical problems with more complex structures is that the loading conditions can be random in space as well as random in time. Lee and Whaley assumed that load $q(x)$ on the beam was random in space such that the Fourier coefficient q_n was completely uncorrelated with q_k , i.e. $\langle q_n q_k \rangle = 0$ for $n \neq k$, which reduced the double summations to single summations (and provided one way to circumvent spatial randomness). An equally realistic approach seemed to be that of assuming that the load would be deterministic in space but random in time, since most experimental investigations of angular vibrations in structures would have to proceed on that path anyway.

In order to assess the effect of various generalized coordinate loading factors q_n , two spatially deterministic, temporally random loading conditions on a beam with simple support were considered and the spatial variation of the ratio of r.m.s. angular to r.m.s. linear displacement along the beam was examined. This spatial variation was compared with Lee and Whaley's "beam averaged r.m.s. amplitudes."

The derivation of the ratio of r.m.s. angular to r.m.s. linear displacement along the beam has been presented in detail in Reference [19], the Phase II Interim report on this contract. For white noise forcing with constant spectral density, this ratio was found to be

$$\frac{\theta_{rms}}{y_{rms}} = \frac{\pi}{L} \sqrt{\frac{\sum_n \sum_k q_n q_k I_{nk} \cos \frac{n\pi x}{L} \cos \frac{k\pi x}{L}}{\sum_n \sum_k q_n q_k I_{nk} \sin \frac{n\pi x}{L} \sin \frac{k\pi x}{L}}} \quad (5.43)$$

where

$$I_{nk} = 8\pi\beta / [(\omega_n^2 - \omega_k^2 - 4\beta^2)^2 + (4\beta\omega_n')^2]$$

$$\text{and } \omega'_n = \sqrt{\omega_n^2 - \beta^2} \quad ; \quad \omega_n^2 = n^4 \frac{EI\pi^4}{mL^4} \quad (5.44)$$

In order to obtain the spatial variation in $\theta_{\text{rms}}/y_{\text{rms}}$, the following parameters for the beam were used:

$$\begin{aligned} L &= 100 \text{ in.} & \rho &= 0.10 \text{ lbm/in.}^3 \\ E &= 10 \times 10^6 \text{ psi.} & g &= 386.04 \text{ lbm-in/lbf-sec.}^2 \\ I &= 5.00 \text{ in.}^4 & A &= 1.25 \text{ in.}^2 \\ \beta &= 38.78 \text{ sec.}^{-1} \end{aligned}$$

It was decided to subject the beam to two spatially determinant loading conditions: (a) The first was

$$q(x) = q_0 \quad (5.45)$$

which when represented by Fourier sine series takes on the form

$$q(x) = \sum_{n=1}^{\infty} q_n \sin \frac{n\pi x}{L} \quad \text{where} \quad q_n = \frac{2q_0}{n\pi} (1 - \cos n\pi) \quad (5.46)$$

(b) The second loading condition was given by

$$q(x) = q_0 \sum_{n=1,3,5}^{\infty} \sin \frac{n\pi x}{L} \quad (5.47)$$

A plot of Eq. (5.43) under these two loading conditions is shown in Figure 25, where $k_{\text{max}} = n_{\text{max}} = 7$.

As a comparison, Lee and Whaley propose a spatially averaged value of

$\theta_{\text{rms}}/y_{\text{rms}} = 1.17 \pi/L = 0.0367 \text{ in.}^{-1}$ which is also shown on Figure 25.

Now consider the special case of a load condition which contains only a single harmonic

$$q(x) = q_{n_0} \sin \frac{n_0 \pi x}{L} \quad (5.48)$$

Eq. (5.43) would then become

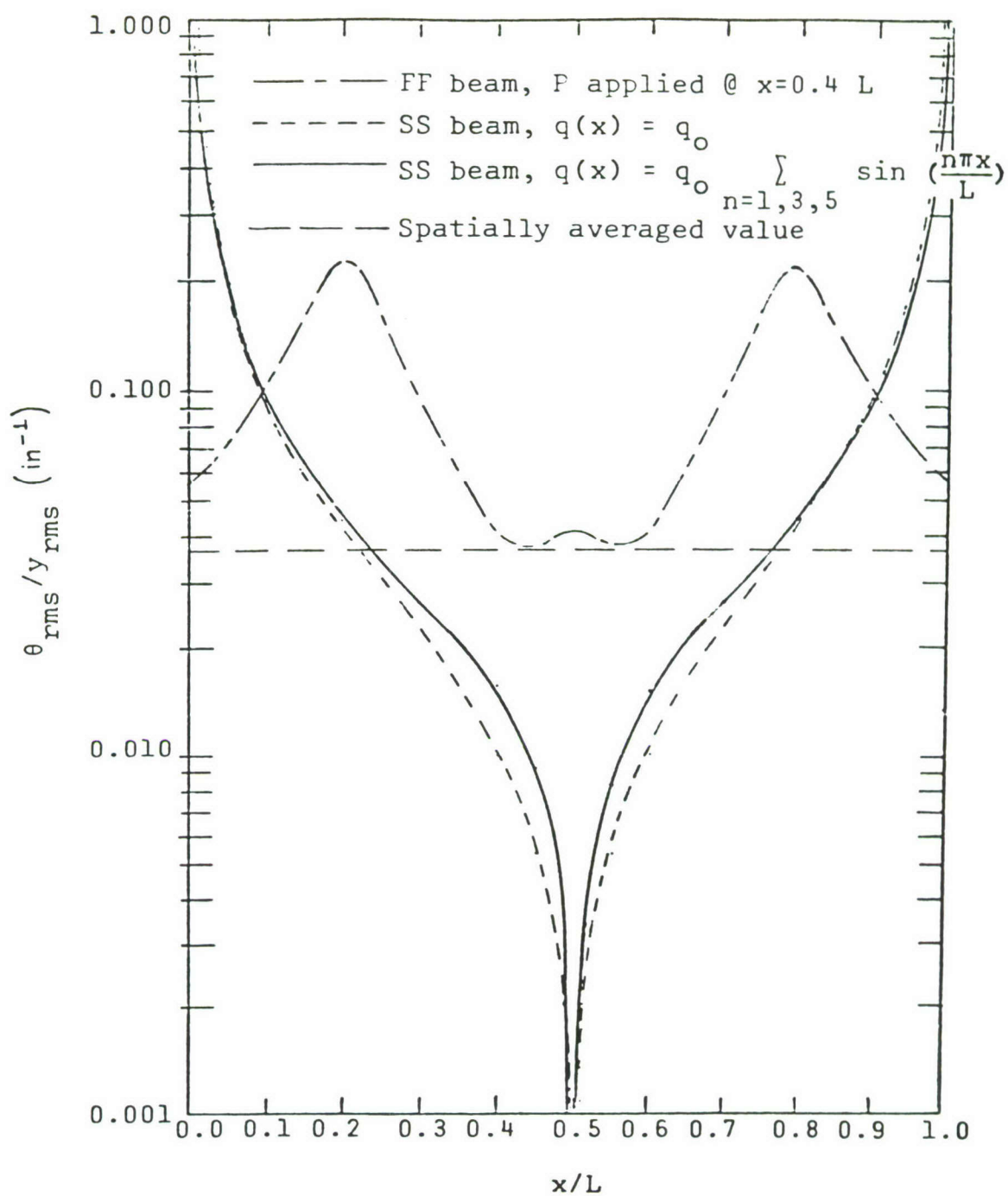


Figure 25 Plot of θ_{rms}/y_{rms} for Simply-Supported and Free-Free Beams Subjected to Temporally Random Loads

$$\frac{\theta_{rms}}{y_{rms}} = \frac{n_o \pi}{L} \cot \frac{n_o \pi x}{L} \quad (5.49)$$

For this loading condition, one sees that θ_{rms}/y_{rms} varies between $\pm\infty$.

As one studies the results of considering the previous loading conditions on a beam with simple support, one may conclude that any load which is symmetrically distributed with respect to the center of the beam will have a variation of θ_{rms}/y_{rms} between, at least, 0 and $+\infty$. Here again, the use of a spatially averaged θ_{rms}/y_{rms} value is subject to question; but, in this case the support condition of $y(0) = y(L) = 0$ and the spatial symmetry of loading are the main source of the extreme variance. For this reason it was thought that a beam in the free-free vibration condition would furnish a more representative variation of θ_{rms}/y_{rms} relative to beam station in comparison to aircraft structural components, which do not normally have simple-support conditions.

One may assess the effects of support conditions by considering the beam in the free-free condition. However, since the panels or beams in an aircraft structure are elastically supported, a question may be raised as to how applicable such analytical results based on free-free conditions might be (notwithstanding the whole airplane which could be in the free-free condition). The justification for using free-free conditions is that the variation in the r.m.s. angular to r.m.s. linear displacement ratio along the beam or over the plate surface is more representative of actual aircraft structural components. Even though there will be discrepancies at the boundaries, they are far less for free-free solutions than they will be for solutions based on simple-support conditions.

For a free-free beam subjected to a concentrated load, P , at any point, $x=a$, the θ_{rms}/y_{rms} ratio is obtained in Reference [19] as

$$\frac{\theta_{rms}}{y_{rms}} = \sqrt{\frac{\sum_{nk} q_n q_k \phi_n' \phi_k' I_{nk}}{\sum_{nk} q_n q_k \phi_n \phi_k I_{nk}}} \quad (5.50)$$

where

$$q_n = \frac{P}{L} \left[\cos \frac{a\lambda_n}{L} + \cosh \frac{a\lambda_n}{L} + \alpha_n \left(\sin \frac{a\lambda_n}{L} + \sinh \frac{a\lambda_n}{L} \right) \right] \quad (5.51)$$

$$\phi_n(x) = \cos \frac{\lambda_n x}{L} + \cosh \frac{\lambda_n x}{L} + \alpha_n \left(\sin \frac{\lambda_n x}{L} + \sinh \frac{\lambda_n x}{L} \right) \quad (5.52)$$

$$\phi_n'(x) = \frac{\lambda_n}{L} \left[-\sin \frac{\lambda_n x}{L} + \sinh \frac{\lambda_n x}{L} + \alpha_n \left(\cos \frac{\lambda_n x}{L} + \cosh \frac{\lambda_n x}{L} \right) \right] \quad (5.53)$$

and

$$\alpha_n = \frac{\sin \lambda_n + \sinh \lambda_n}{\cos \lambda_n - \cosh \lambda_n} \quad (5.54)$$

I_{nk} may be expressed in the form

$$I_{nk} = 8\pi\beta / [\omega_n^2 - \omega_k^2]^2 + 8\beta^2 (\omega_n^2 + \omega_k^2)] \quad (5.55)$$

where

$$\omega_n^2 = \frac{\lambda_n^4}{L^4} \cdot \frac{EI}{m} \quad (5.56)$$

and λ_n represents the roots of the frequency equation

$$1 = \cosh \lambda_n \cos \lambda_n \quad (5.57)$$

The ratios of angular r.m.s. to linear r.m.s., Eq. (5.50) at different stations along the beam, have been calculated for the beam properties given previously and are presented in Table 3 is $k_{max} = n_{max} = 1, 2, 3, 4$, and 6. This ratio is $k_{max} = n_{max}$ is also shown on Figure 25.

As one examines the θ_{rms}/y_{rms} ratios in Table 3 for the one-term approximation one sees that they get quite large as the node points for the first free-free mode shape are approached. With the additional higher frequency mode shapes

TABLE 3
 θ_{rms}/y_{rms} FOR FREE-FREE BEAM SUBJECT TO CONCENTRATED LOAD
(a/L = 0.4)

x/L	θ_{rms}/y_{rms} $n_{max}=k_{max}=1$	θ_{rms}/y_{rms} $n_{max}=k_{max}=2$	θ_{rms}/y_{rms} $n_{max}=k_{max}=3$	θ_{rms}/y_{rms} $n_{max}=k_{max}=4$	θ_{rms}/y_{rms} $n_{max}=k_{max}=6$
0	0.0465	0.0507	0.0518	0.0550	0.0563
0.10	0.0851	0.0961	0.0981	0.1014	0.1014
0.15	0.1415	0.1568	0.1562	0.1504	0.1506
0.16	0.1630	0.1766	0.1741	0.1645	0.1655
0.17	0.1922	0.1999	0.1942	0.1799	0.1820
0.18	0.2345	0.2258	0.2152	0.1961	0.1993
0.19	0.3013	0.2514	0.2342	0.2111	0.2153
0.20	0.4230	0.2697	0.2461	0.2222	0.2267
0.21	0.7159	0.2725	0.2460	0.2262	0.2299
* 0.22	2.4141	0.2572	0.2332	0.2215	0.2235
0.23	1.7015	0.2304	0.2121	0.2090	0.2091
0.24	0.6199	0.2007	0.1833	0.1919	0.1906
0.25	0.3750	0.1733	0.1655	0.1732	0.1712
0.30	0.1166	0.0898	0.0913	0.0992	0.0998
0.35	0.0603	0.0551	0.0583	0.0595	0.0627
0.40	0.0332	0.0389	0.0414	0.0409	0.0416
0.45	0.0151	0.0317	0.0325	0.0374	0.0383
0.50	0.0000	0.0299	0.0296	0.0385	0.0414
0.55	0.0151	0.0321	0.0329	0.0379	0.0387
0.60	0.0332	0.0394	0.0421	0.0415	0.0423
0.65	0.0603	0.0558	0.0591	0.0602	0.0634
0.70	0.1166	0.0906	0.0921	0.0999	0.1005
0.75	0.3749	0.1739	0.1654	0.1734	0.1712
0.76	0.6197	0.2009	0.1876	0.1913	0.1901
0.77	1.6999	0.2297	0.2105	0.2077	0.2079
* 0.78	2.4173	0.2551	0.2305	0.2193	0.2213
0.79	0.7161	0.2689	0.2422	0.2234	0.2271
0.80	0.4231	0.2654	0.2418	0.2191	0.2235
0.81	0.3014	0.2472	0.2302	0.2081	0.2123
0.82	0.2346	0.2223	0.2418	0.1934	0.1967
0.83	0.1922	0.1971	0.1915	0.1778	0.1799
0.84	0.1630	0.1745	0.1721	0.1628	0.1638
0.85	0.1415	0.1552	0.1547	0.1491	0.1493
0.90	0.0851	0.0956	0.0976	0.1010	0.1010
1.00	0.0465	0.0506	0.0518	0.0550	0.0564

considered, the extreme variation in $\theta_{\text{rms}}/y_{\text{rms}}$ as a function of x/L is attenuated, but the fundamental mode shape still retains a significant influence on the shape of the curve. Nevertheless, the variation between the maximum and minimum values of $\theta_{\text{rms}}/y_{\text{rms}}$ along the beam length is finite but of sufficient magnitude to imply that a spatially averaged value could have significant error.

5.3 NASTRAN ANALYSIS OF STIFFENED CURVED PANEL SUBJECTED TO TEMPORALLY RANDOM CONCENTRATED LOAD

In order to assess the linear-to-angular relationship in a structural component of a somewhat more complex geometry, it was decided to analyze a stiffened curved panel. A NASTRAN finite element model of the curved panel was developed and is shown in Figure 26. The model consists of 726 degrees of freedom, 100 quadrilateral bending elements, and 30 beam elements. This panel was assumed to be in the free-free condition with the random concentrated load applied at grid point station 406. The response of the curved panel was evaluated for a 20-2560 Hz. band width and for the seven octaves in between as follows:

- (a) 20 - 40 Hz
- (b) 40 - 80 Hz
- (c) 80 - 160 Hz
- (d) 160 - 320 Hz
- (e) 320 - 640 Hz
- (f) 640 - 1280 Hz
- (g) 1280 - 2560 Hz

The output of the NASTRAN analysis included the following quantities:

- (a) w_{rrms} - temporal average of displacement normal to plate
- (b) $\theta_{\theta rms}$ - temporal average of angular displacement about θ -axis
- (c) $\theta_{z rms}$ - temporal average of angular displacement about z-axis

at each of the following sets of grid points:

- (a) 106, 206, 306, ... 1106
- (b) 401, 402, 403, ... 411
- (c) 104, 204, 304, ... 1104

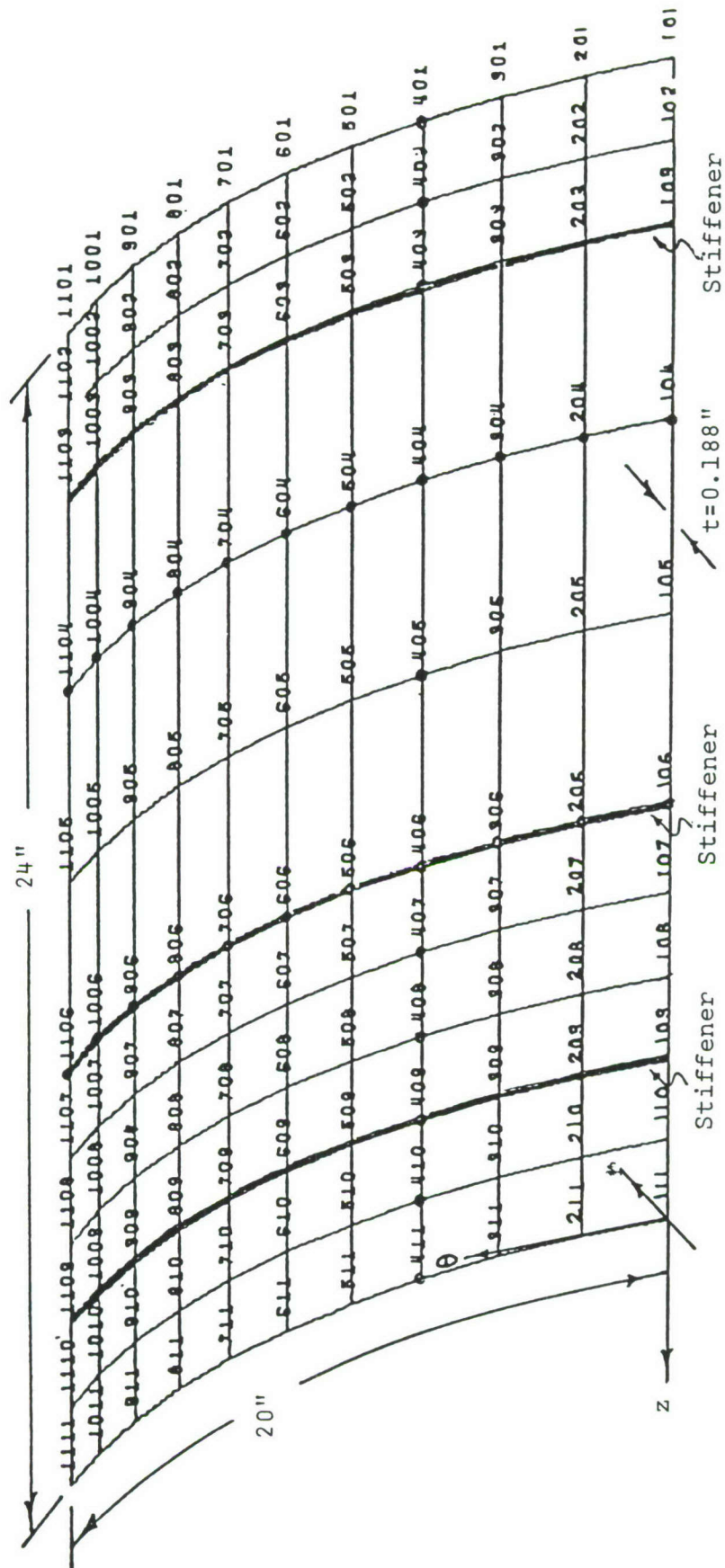


Figure 26 NASTRAN Model of Stiffened Curved Panel

Figure 27 shows the ratio of RMS total resultant angular displacement to RMS linear displacement vs. grid point station for 20-2560 Hz. frequency band for the bare curved panel. In this figure the straight lines plotted are the spatial average as represented by the quantity

$$\frac{\sum_{106}^{1106} \sqrt{\int_{\omega_1}^{\omega_2} S_{\theta_{\theta}} d\omega + \int_{\omega_1}^{\omega_2} S_{\theta_z} d\omega}}{\sum_{106}^{1106} \sqrt{\int_{\omega_1}^{\omega_2} S_{w_r} d\omega}} \quad \text{for Stations 106 to 1106} \quad (5.58)$$

where

$$\begin{aligned} \omega_1 &= 20 \text{ Hz} \\ \omega_2 &= 2560 \text{ Hz} \end{aligned} \quad (5.59)$$

and

$$\overline{w_r^2} = \int_{\omega_1}^{\omega_2} S_{w_r} d\omega \quad ; \quad \overline{\theta_{\theta}^2} = \int_{\omega_1}^{\omega_2} S_{\theta_{\theta}} d\omega \quad ; \quad \overline{\theta_z^2} = \int_{\omega_1}^{\omega_2} S_{\theta_z} d\omega \quad (5.60)$$

This particular method of averaging corresponds to that of Lee and Whaley in their aforementioned report.

Figure 28 presents the ratio of RMS angular displacement component to spatially averaged RMS linear displacement vs. grid point station for 20-2560 Hz. for the curved panel. The straight lines plotted are the spatial averages as given by the expressions

$$\frac{\theta_{\text{rms avg}}}{w_{\text{rrms avg}}} = \frac{\sum_{106}^{1106} \sqrt{\int_{\omega_1}^{\omega_2} S_{\theta_{\theta}} d\omega}}{\sum_{106}^{1106} \sqrt{\int_{\omega_1}^{\omega_2} S_{w_r} d\omega}} \quad (5.61)$$

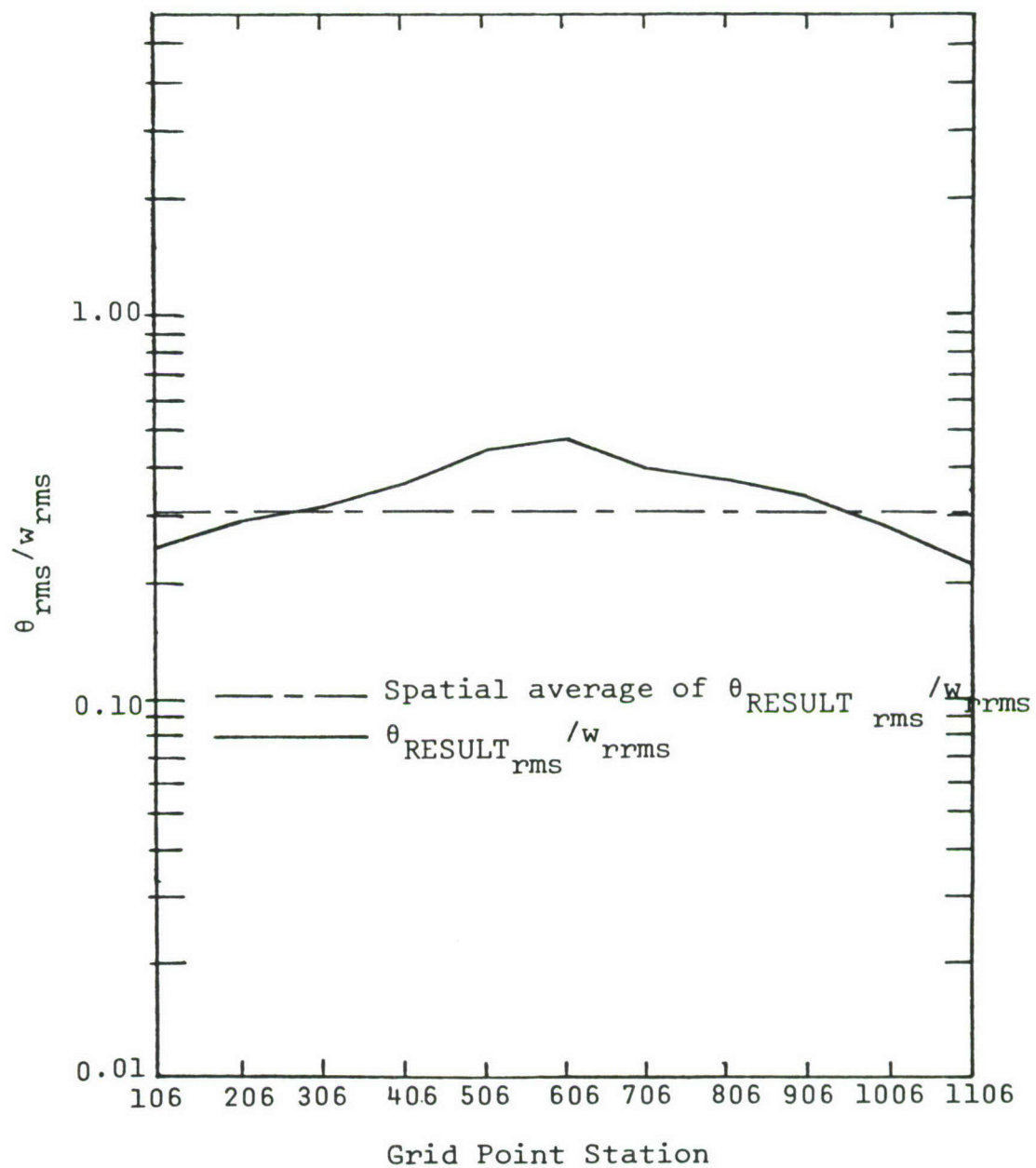


Figure 27 Ratio of RMS Resultant Angular Displacement to RMS Linear Displacement vs. Grid Point Station for Curved Plate with Stiffeners, 20-2560 Hz.

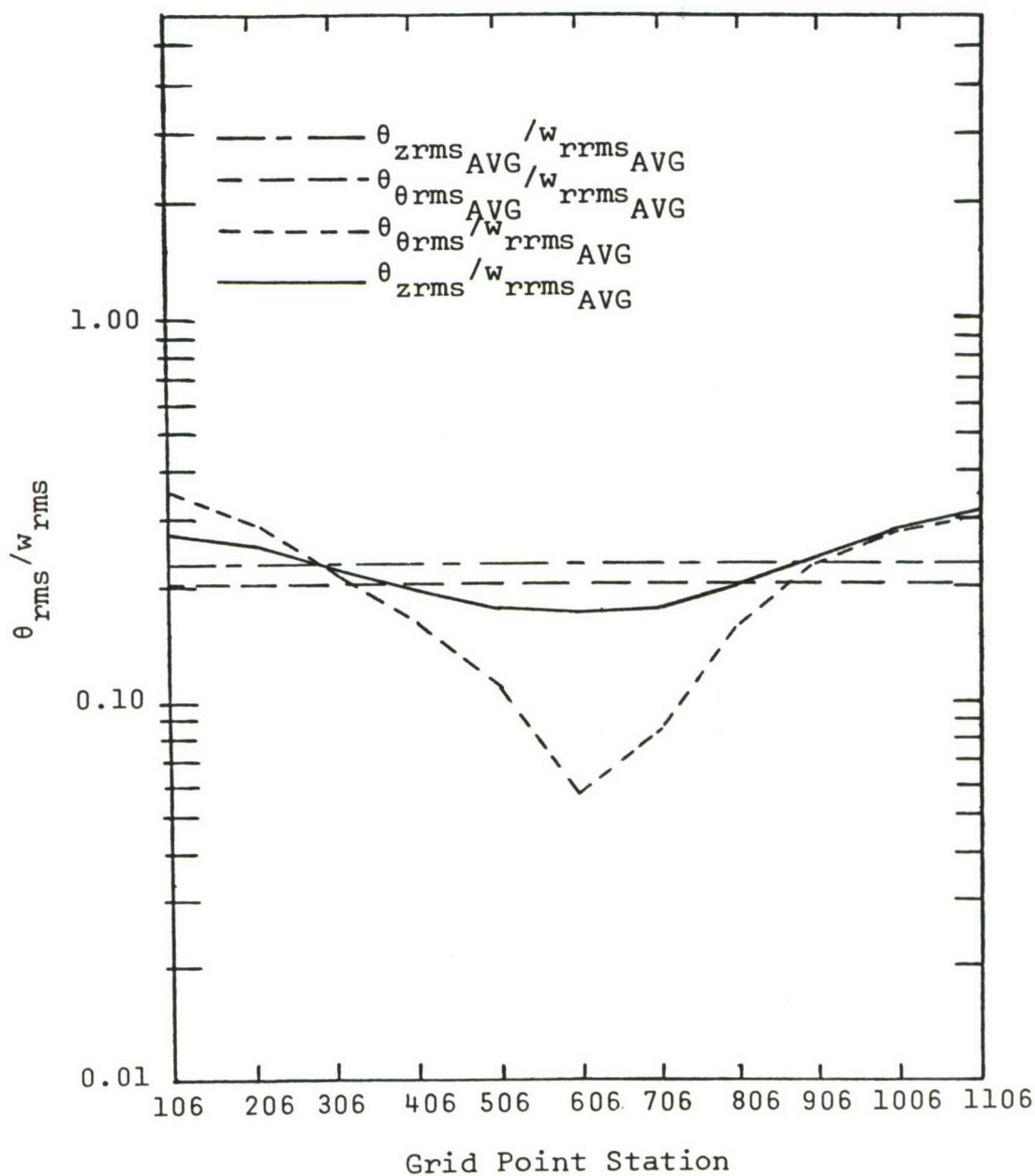


Figure 28 Ratio of RMS Angular Displacement to Spatially Average RMS Linear Displacement vs. Grid Point Station for Curved Plate with Stiffeners, 20-2560 Hz.

$$\frac{\theta_{zrms_avg}}{y_{rms_avg}} = \frac{\sum_{106}^{1106} \sqrt{\int_{\omega_1}^{\omega_2} S_{\theta_z} d\omega}}{\sum_{106}^{1106} \sqrt{\int_{\omega_1}^{\omega_2} S_{w_r} d\omega}} \quad (5.62)$$

for stations 106 to 1106. The variable quantity is actually the angular displacement component at each station except that it has been normalized by the spatial average of the linear displacement in accordance with the ratios

$$\frac{\theta_{\theta rms}}{\sum_{106}^{1106} \sqrt{\int_{\omega_1}^{\omega_2} S_{w_r} d\omega}} \quad \text{and} \quad \frac{\theta_{zrms}}{\sum_{106}^{1106} \sqrt{\int_{\omega_1}^{\omega_2} S_{w_r} d\omega}} \quad (5.63)$$

The subsequent plots, Figures 29 through 36, present the angular-displacement-component-to-linear-displacement ratios of $\theta_{\theta rms}/w_{rrms}$ and θ_{zrms}/w_{rrms} as a function of grid point stations 106 through 1106 for the frequency band width of 20-2560 Hz. and the seven octaves in between. In each of these plots, the straight lines represent the spatial average of the RMS angular component divided by the spatial average of the RMS linear displacement. For the 20-2560 Hz. case, there are two additional lines plotted which represent the spatial average of all grid points considered in this study on the curved panel.

An examination of the plots for the 20-2560 Hz. band widths (Figures 27 through 29) reveal a somewhat mild spatial variation in the linear-to-angular relationships for the stiffened curved panel subjected to the temporally random concentrated load. However, as one studies the response of the panel to the narrow band octaves, e.g. Figure 30, the maximum value for the r.m.s. angular-to-r.m.s. linear displacement ratio (θ_{zrms}/w_{rrms}) was almost 100 times greater than the minimum value. Thus, for narrow band random excitation, the results

definitely preclude the use of a spatial average as a representative value for the stiffened curved panel subjected to the aforementioned loading condition.

Additional plots for stations 401 through 411 and 104 through 1104 may be found in Reference [24]. Also, a study where mass was added to the panel was reported in this reference.

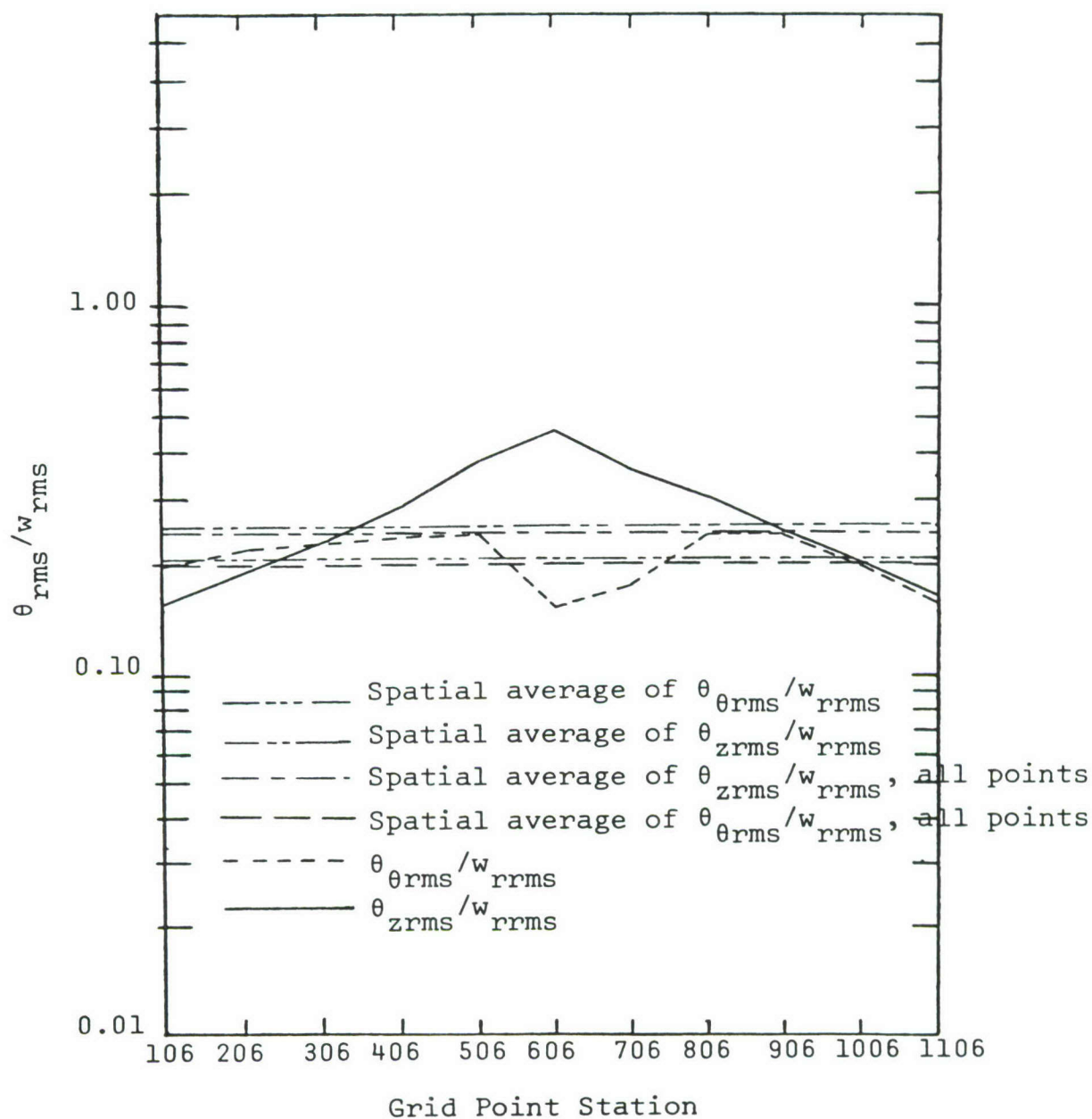


Figure 29 Ratio of RMS Angular Displacement to RMS Linear Displacement vs. Grid Point Station for Curved Plate with Stiffeners, 20-2560 Hz.

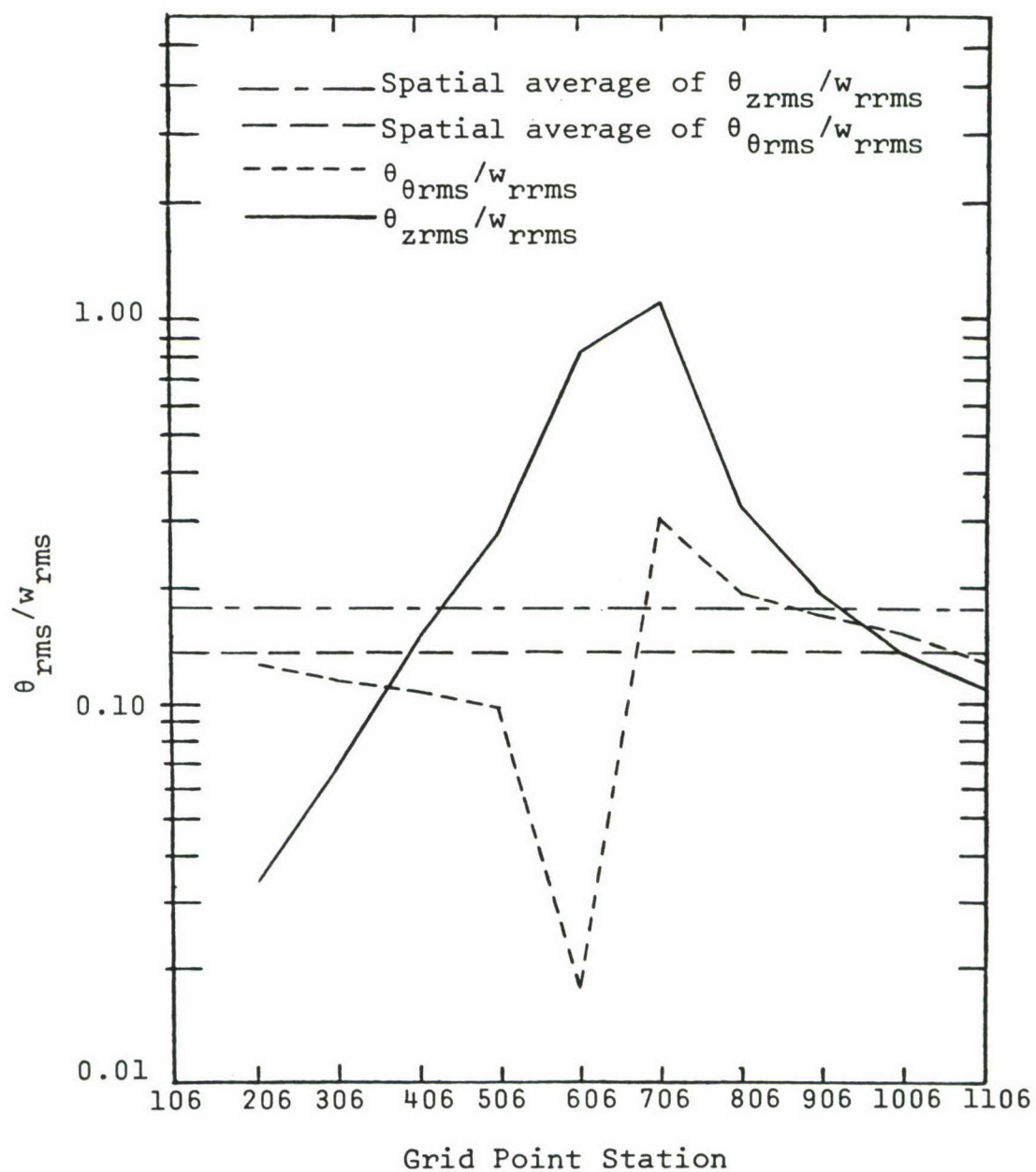


Figure 30 Ratio of RMS Angular Displacement to RMS Linear Displacement vs. Grid Point Station for Curved Plate with Stiffeners, 20-40 Hz.

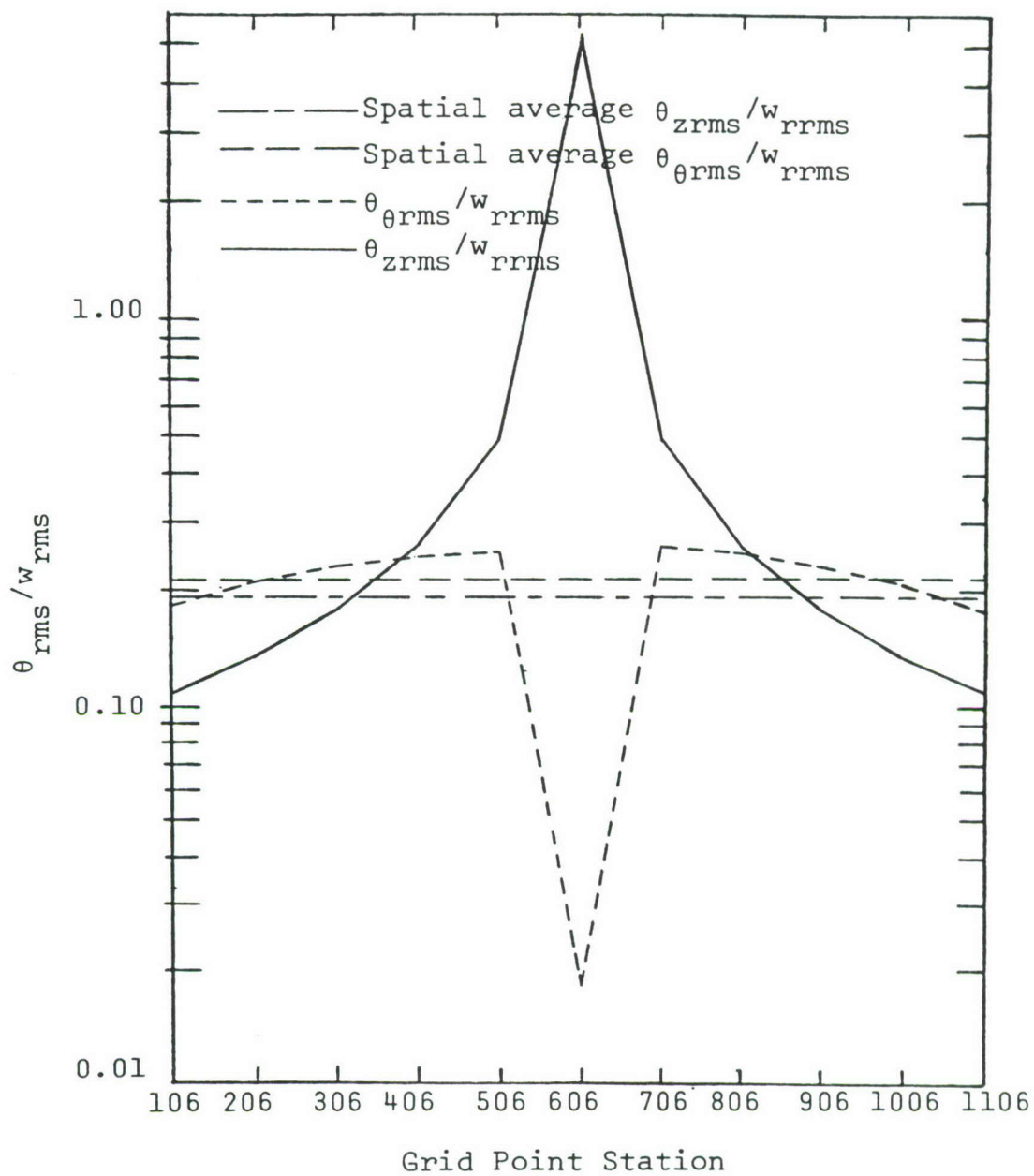


Figure 31 Ratio of RMS Angular Displacement to RMS Linear Displacement vs. Grid Point Station for Curved Plate with Stiffeners, 40-80 Hz.

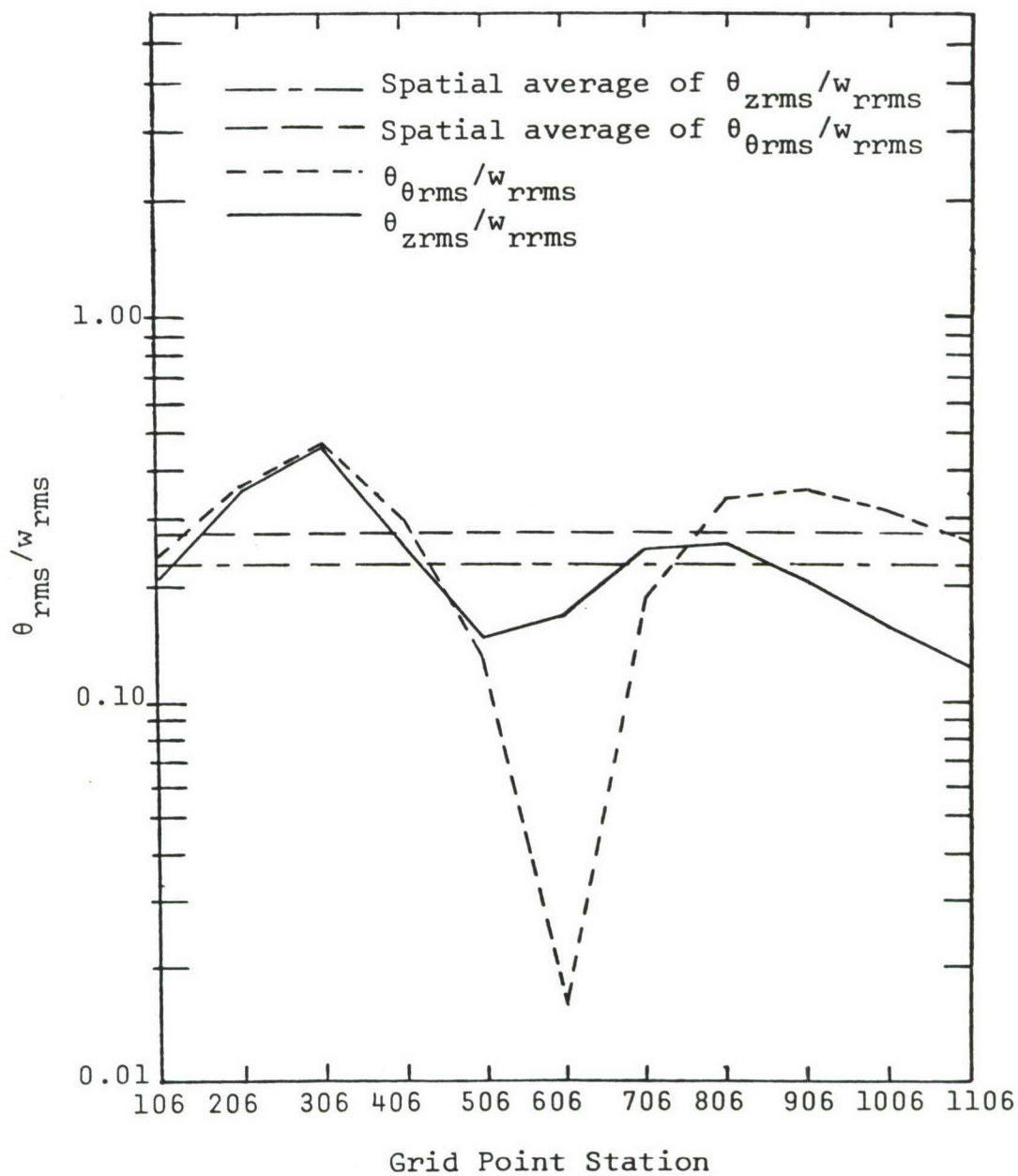


Figure 32 Ratio of RMS Angular Displacement to RMS Linear Displacement vs. Grid Point Station for Curved Plate with Stiffeners, 80-160 Hz.

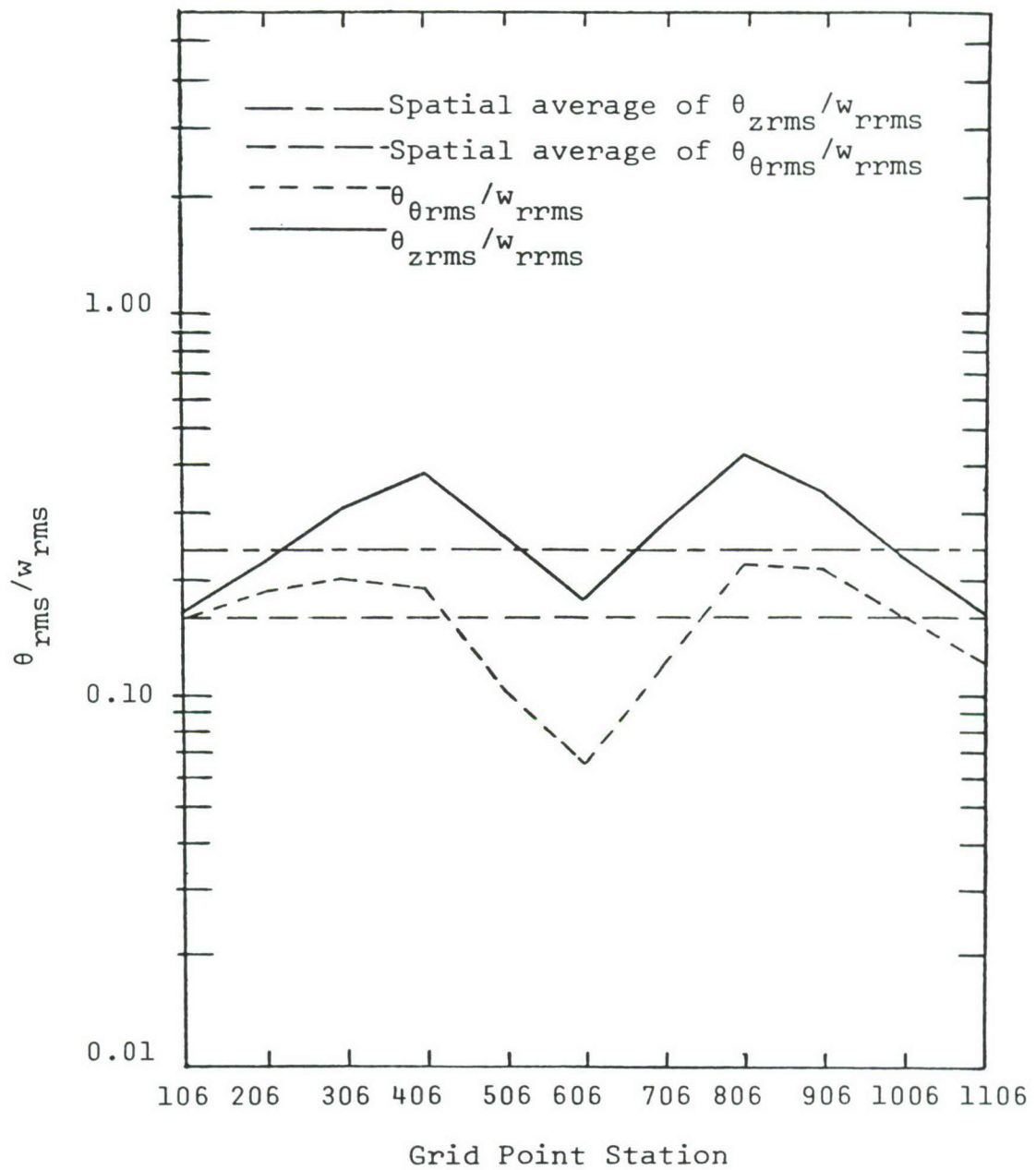


Figure 33 Ratio of RMS Angular Displacement to RMS Linear Displacement vs. Grid Point Station for Curved Plate with Stiffeners, 160-320 Hz.

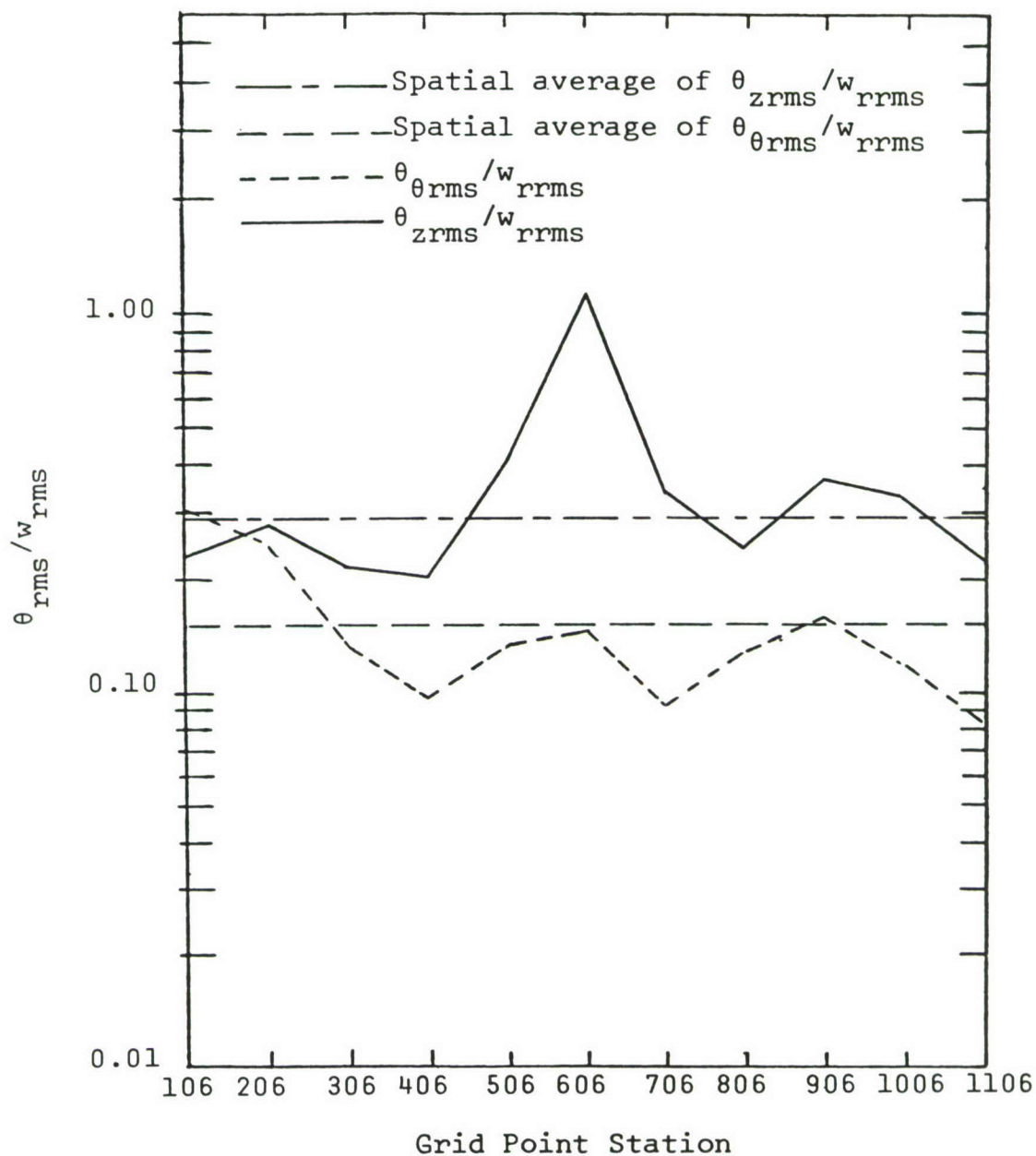


Figure 34 Ratio of RMS Angular Displacement to RMS Linear Displacement vs. Grid Point Station for Curved Plate with Stiffeners, 320-640 Hz.

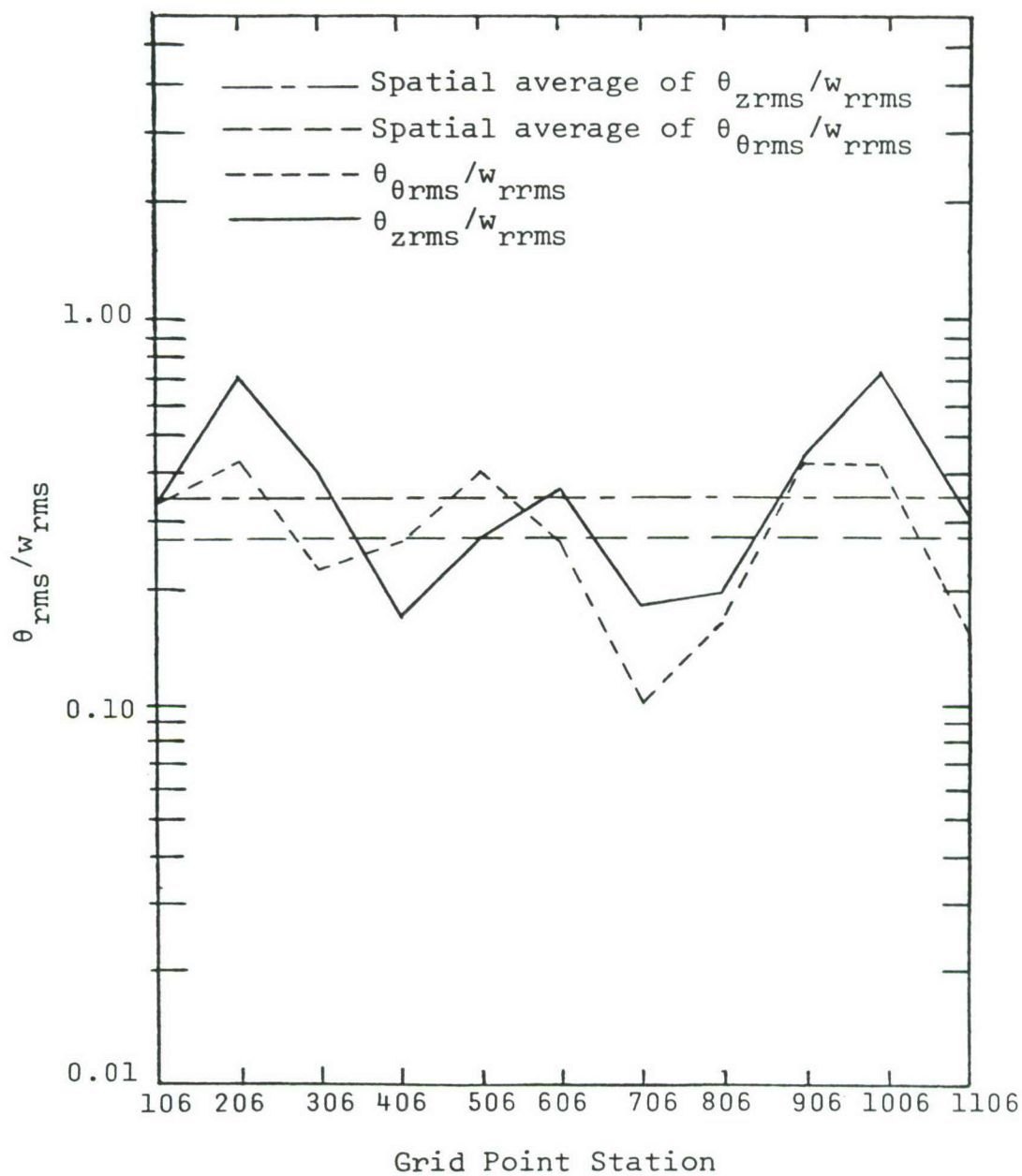


Figure 35 Ratio of RMS Angular Displacement to RMS Linear Displacement vs. Grid Point Station for Curved Plate with Stiffeners, 640-1280 Hz.

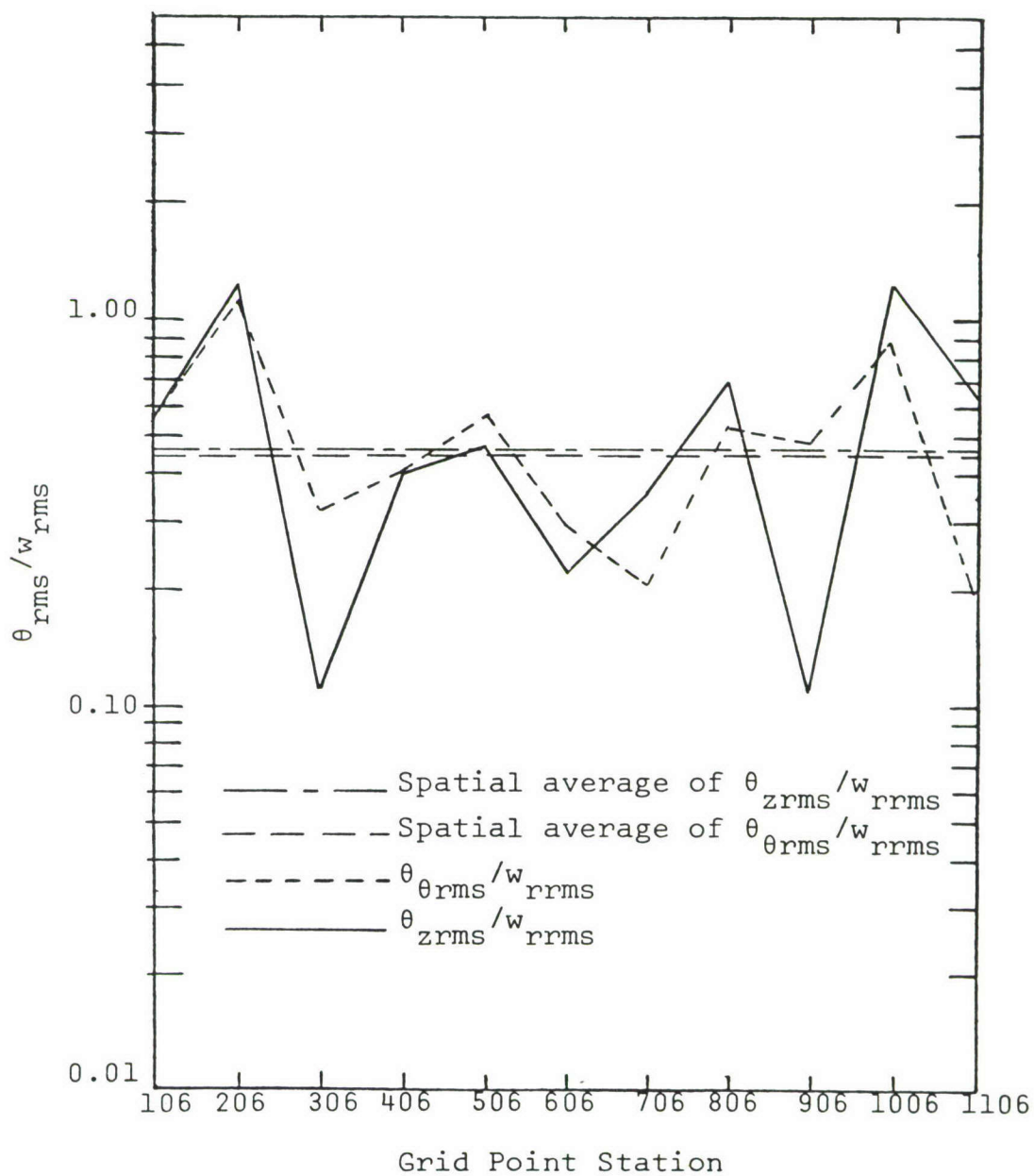


Figure 36 Ratio of RMS Angular Displacement to RMS Linear Displacement vs. Grid Point Station for Curved Plate with Stiffeners, 1280-2560 Hz.

5.4 VARIATION OF ANGULAR TO LINEAR DISPLACEMENT RATIOS IN STRUCTURAL COMPONENTS

As one studies the results of the analytical investigation for random vibrations of a rigid bar on springs, a simple-supported beam and a free-free beam, the conclusion would be made that the spatial variance of the r.m.s. angular-to-r.m.s. linear displacement ratio is significant enough to warrant not using a spatial average for these structures if great accuracy is desired. If large errors (for the free-free beam this would be a factor of 7 between the maximum r.m.s. amplitude occurring and Lee & Whaley's proposed spatial average) can be tolerated, then a single spatial average for the structure is convenient to use; but, it seems the behavior of the structure ought to be understood in some detail so that the applicability of the spatial average is assessed. In the panel example, the ratio was formed by actually computing (via a finite element model for a specific random load case) both r.m.s. angular and linear responses at many points. The conclusion which may be drawn from the example is that even if one has knowledge of the ratio of averages, that quantity is not a particularly good estimation of the ratio at a given point as indicated by those plots encompassing octave bands. The spatial average might be considered an acceptable quantity when considering the 20-2560 Hz. case, but it is felt that the response of different structures under random loading should be investigated before a definite conclusion is drawn. The examples investigated in the previous paragraphs also show cases when the ratio is quite sensitive to structural parameters, load distribution, and the frequency bands.

One means of eliminating the effects of end conditions is by simply considering an infinitely long beam. For this case a traveling wave solution may be written. For sinusoidal excitation at a single point

$$y = Ae^{i(kx-\omega t)} \quad (5.64)$$

where the wave number k is related to frequency by the usual dispersion relation for one-dimensional bending waves:

$$(2\pi f)^2 = k^4 \kappa^2 c_\ell^2 \quad (5.65)$$

In this expression, κ is the radius of gyration and c_ℓ is the extensional wave speed. From (5.64), one may compute $\theta = \frac{dy}{dx}$

$$\theta = ikAe^{i(kx-\omega t)} = iky \quad (5.66)$$

The factor i simply comes from using complex arithmetic to keep track of phase in time and space. It will drop out when temporal averaging is performed to obtain r.m.s. values.

$$\frac{\theta_{\text{rms}}}{y_{\text{rms}}} = \sqrt{\frac{2\pi f}{\kappa c_\ell}} \quad (5.67)$$

For random steady state excitation at a point we note that, by linearity, the portion of response in a frequency band Δf is due only to the excitation in that band so we may write

$$\frac{S_\theta(f)}{S_y(f)} = \frac{2\pi f}{\kappa c_\ell} \quad (5.68)$$

This ratio does not vary in space and requires only intensive beam parameters κ and c_ℓ rather than a global length L . It can be expected to hold for beams of finite length at frequencies sufficiently high that $L \gg \lambda$.

The infinite beam may be thought of as having an infinite number of natural frequencies which are spaced infinitesimally close together. Since for this case, an exact relation exists between $S_y(f)$ and $S_\theta(f)$, one might expect that a long beam with a large but finite modal density might show approximately this ratio. It is shown in Section IV that, for a plate with high modal density, one may indeed exploit this idea to estimate angular response from component energy expressed as a mass-weighted mean square translational velocity.

SECTION VI

ANGULAR VIBRATION MEASUREMENT

The major objective of the contract effort was the development of methods for prediction of angular vibration. However, it was known from the beginning that some attention would have to be given to purely measurement problems. This view was motivated by two factors:

- (1) The scarcity of experimental data suitable for testing of analytical methods meant that Anamet had to be capable of performing its own verification tests. This capability had to be ready when analytical work reached the testing stage in order to produce a reliable and timely product.
- (2) Modern methods of structural response prediction, angular or otherwise, often depend on models built entirely or in part from measured data. Two examples are experimental modal analysis and statistical energy analysis where component SEA parameters are obtained by measurement.

The measurement of dynamic rotations at specific points on an elastic structure is not new although concern about angular displacements at the microradian level seems to be confined to optical system applications. Historically, three basic methods of measurement have been used.

- (1) Outputs of translational motion sensors mounted a known distance apart may be differenced.
- (2) A light beam may be reflected off the point in question and its lateral displacement measured.
- (3) Inertia torque on a suspended seismic mass may be sensed.

The most critical angular measurements on airborne optical or laser systems are generally those where motion of a reflecting surface is transduced. Typical requirements in this case may include:

- (1) Bandwidth. As usual in structural dynamics, the frequency composition of unwanted motion often contains information which is highly useful to the designer attempting to reduce such motion. Frequency components in the 0.1 to 1.0 kHz range may be particularly important since, in general, they cannot be effectively removed by active servo systems.
- (2) Size and Weight. Mirror assemblies as small as 0.1 m (4 in.) and weighing less than 1 kg (2.2 lb.) must be instrumented. Transducers must neither change the dynamic properties of the assembly nor interfere with the intended optical function.
- (3) Flexibility of use. During development of component assemblies, sensors must be installed and removed quickly and easily without the requirement of elaborate fixturing. This tends to favor inertially-referenced methods over optical sensing. Furthermore, the sensing of displacement by optical systems rather than velocity or acceleration may make high frequency measurements difficult.
- (4) Cost. As usual, the use of standard devices made in production quantities and usable for other purposes is desirable.

It quickly became clear that differencing of translational acceleration signals was the most appropriate method for the present effort. Other investigators [25, 26] have used the method and have concluded that it is quite practical. However, it appeared that a better quantitative understanding of limitations and error sources was desirable. The work described in this chapter was intended to develop that understanding.

Theoretical derivations are presented for estimating errors introduced by noise in individual channels, frequency dependent gain and phase mismatching between channels, and flexure of the mounting surface. Effects of the first two error sources are demonstrated by experiment and it is demonstrated that mismatch error can be reduced by appropriate data processing. Finally, an expression is derived for coherence of a measured angular frequency response when angular response is obtained by differencing.

6.1 THEORETICAL ERROR ANALYSIS

Figure 37 shows a typical arrangement for measuring dynamic rotation. The linear sensors must often be spaced quite closely, either to obtain a good discrete approximation to the angle of rotation at a point, or simply due to space constraints. The difference signal may be much smaller than either of the individual signals. The question then arises, "How 'good' must the individual linear sensing channels be in order that their difference can provide a meaningful estimate of angular motion?"

In attempting to understand the dynamic behavior of a structure, spectral measurements obtained through discrete Fourier transforms are often used. The two most common types of frequency domain data are the power spectral densities of response variables and complex frequency response functions measured for spatially fixed excitations. The intent of this section is to present some quantitative methods for estimating the reliability of these two data types for the case where the response quantity is an angular motion obtained by differencing of signals from translational motion sensors.

The figures of merit used are narrow band signal power/noise power for PSD estimates and the ordinary coherence function for frequency response estimates. Both are derived for the special case of differential sensing in terms of the corresponding quantities for the individual channels.

6.1.1 Narrow-Band Signal-to-Noise Ratio

In Figure 37 two single axis translation accelerometers are mounted a distance Δx apart and their output signals are conditioned and subtracted. The pair then form a differential acceleration sensing system. Each leg is assumed to be a linear, time-invariant system with impulse response $h_1(\tau)$ and frequency response $H_1(f)$. In general, $H_1(f)$ and $H_2(f)$ will be slightly different and will both show a weak dependence on frequency in the range of interest. Noise in each channel is simulated by $n_1(t)$ and $n_2(t)$ as shown. Noise sources are assumed to be

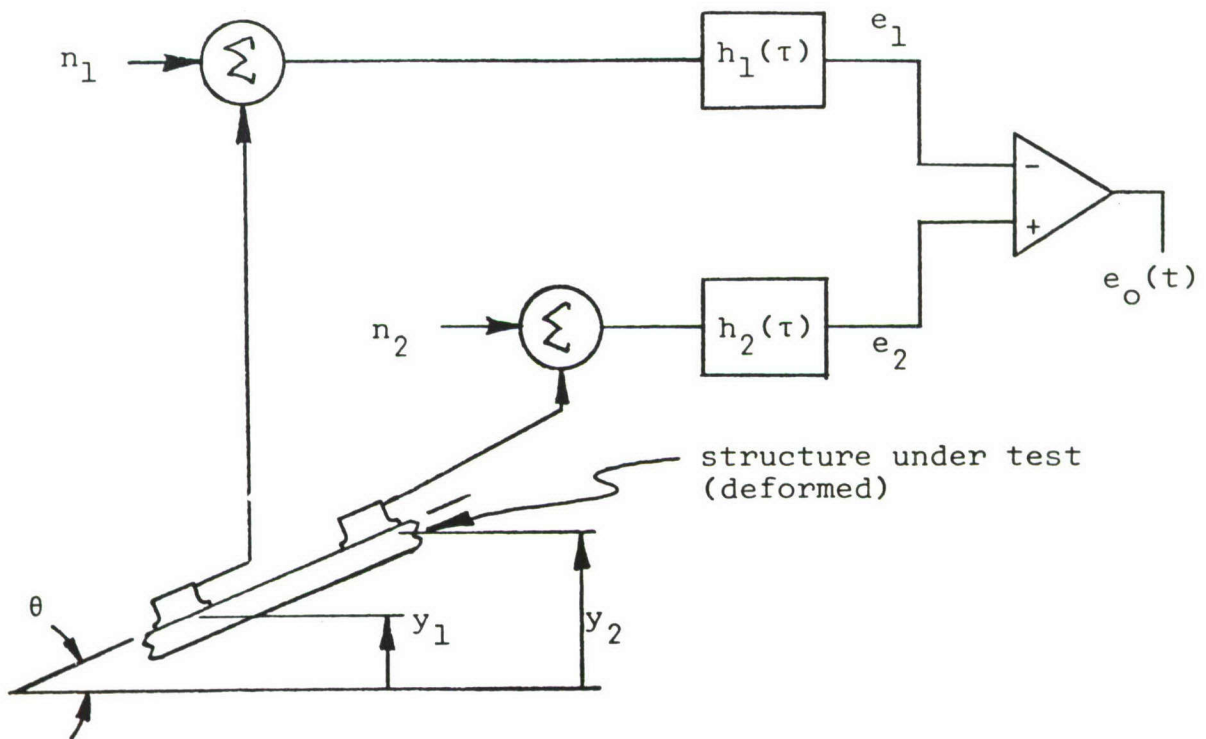
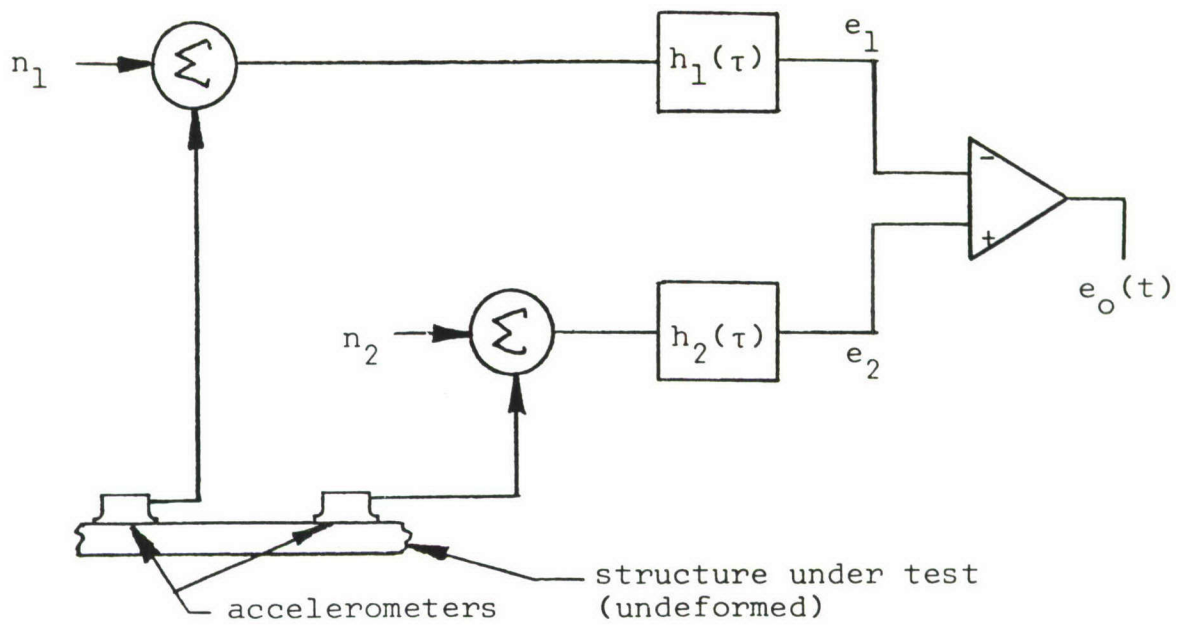


Figure 37 Measurement of Angular Acceleration by Differencing

correlated neither with the measurands \ddot{y}_1 and \ddot{y}_2 nor with each other.

One component of the angular acceleration at the point midway between the transducers can be estimated as

$$\hat{\hat{\theta}}(t) = e_0/c\Delta x \quad (6.1)$$

where c is constant and equal to some nominal value of $|H_1|$ or $|H_2|$. If $\ddot{\theta}(t)$ is the true value of angular acceleration, the estimate $\hat{\hat{\theta}}(t)$ which one obtains from $e_0(t)$ will deviate from $\ddot{\theta}(t)$ due to imperfections introduced by noise sources n_1 and n_2 as well as differences between c and H_{1_actual} and H_{2_actual} . We therefore define an ideal transducing system which has $n_1 = n_2 = 0$ and $H_{1_ideal} = H_{2_ideal} = c$ where c is close to H_{1_actual} and H_{2_actual} . This definition will allow calculation of a figure of merit for the actual system.

The figure of merit will be called narrow-band signal-to-noise ratio and is a function of frequency defined as

$$S_{NR}(f)_{dB} = 10 \log \left[\frac{S_{\hat{\hat{\theta}}_I}(f)}{S_{(\hat{\hat{\theta}}_A - \hat{\hat{\theta}}_I)}(f)} \right] \quad (6.2)$$

where

$S_{NR}(f)$ = narrow-band signal-to-noise ratio

$\hat{\hat{\theta}}_I(t)$ = estimate of $\ddot{\theta}(t)$ which would be obtained with an ideal system

$S_{\hat{\hat{\theta}}_I}(f)$ = power spectral density of $\hat{\hat{\theta}}_I$

$\hat{\hat{\theta}}_A(t)$ = estimate of $\ddot{\theta}(t)$ obtained with the actual system

$S_{(\hat{\hat{\theta}}_A - \hat{\hat{\theta}}_I)}(f)$ = power spectral density of $\hat{\hat{\theta}}_A - \hat{\hat{\theta}}_I$

The quantity $\hat{\hat{\theta}}_A - \hat{\hat{\theta}}_I$ is a residual random variable similar to those used to define partial and multiple coherence functions (Reference [27]).

It should be clear that S_{NR} will be degraded only by mismatching between H_1 and H_2 . If they are perfectly matched and both equal to c , a small shift in c will have no effect. Thus, there is no loss of generality in allowing c to be a complex function of frequency although Eq. (6.1) should properly be written in terms of Fourier transforms.

In all of the following, the power spectral density of a time history will be represented simply by the ensemble averaged value of the squared magnitude of its Fourier transform. This ignores a number of important mathematical details [28] but is appropriate since virtually all applied work will be carried out using digital FFT methods. The normalizing factor which is needed to compensate for finite record length when using a discrete Fourier transform is also omitted. This simplifies the algebra and is allowable since the final results are arranged in dimensionless form where this factor would cancel out.

If the impulse response of each leg of the ideal system is $h(\tau)$, then

$$h(\tau) = \mathcal{F}^{-1}(c(f)) \quad (6.3)$$

and

$$e_i(t) = h_i(t) * \ddot{y}_i(t) \quad , \quad i = 1, 2 \quad (6.4)$$

where $*$ is the convolution operator and $\mathcal{F}^{-1}()$ is the inverse Fourier transform operator. Taking the Fourier transform of Eq. (6.4)

$$E_i(f) = -(2\pi f)^2 c(f) Y_i(f), \quad i = 1, 2 \quad (6.5)$$

where

$$E_i(f) = \mathcal{F}[(e_i(t))]$$

$$Y_i(f) = \mathcal{F}[y_i(t)]$$

From Figure 37, for either the actual or ideal system

$$(\Delta x) \ddot{\theta}(t) = \ddot{y}_2 - \ddot{y}_1 \quad (6.6)$$

or

$$(\Delta x) \theta(f) = Y_2(f) - Y_1(f) \quad (6.7)$$

For the ideal system, Eqs. (6.5) and (6.7) may be combined

$$-(2\pi f)^2 \hat{\theta}_I = \frac{E_2 - E_1}{c\Delta x} \quad (6.8)$$

where

$$\hat{\theta}_I = \mathcal{F}[\hat{\theta}_I]$$

Forming the conjugate square of Eq. (6.8)

$$|-4\pi^2 f^2 \hat{\theta}|^2 = \frac{|E_2|^2 + |E_1|^2 - 2\text{Re}[E_1^* E_2]}{|c|^2 (\Delta x)^2} \quad (6.9)$$

where

* = complex conjugate

Ensemble averaging Eq. (6.8) and using the derivative relation for power spectral density

$$S_{\ddot{u}}(f) = (2\pi f)^4 S_u(f) \quad (6.10)$$

gives

$$S_{\hat{\theta}_I} = \frac{S_{e_2} + S_{e_1} - 2\text{Re}[S_{e_1 e_2}]}{|c|^2 (\Delta x)^2} \quad (6.11)$$

An expression for $S_{(\hat{\theta}_A - \hat{\theta}_I)}$ is obtained in a similar way. For the actual system

$$e_i = h_i * (\ddot{y}_i + n_i) \quad , \quad i = 1, 2 \quad (6.12)$$

Transforming Eq. (6.12)

$$E_i = H_i [-(2\pi f)^2 Y_i + N_i] \quad , \quad i = 1, 2 \quad (6.13)$$

Combining Eqs. (6.7) and (6.13)

$$-(2\pi f)^2 \Delta x \hat{\theta}_A = \left[\frac{E_2}{H_2} - N_2 \right] - \left[\frac{E_1}{H_1} - N_1 \right] \quad (6.14)$$

By linearity of the Fourier transform

$$\mathcal{F}(\hat{\theta}_A - \hat{\theta}_I) = \hat{\theta}_A - \hat{\theta}_I \quad (6.15)$$

From Eqs. (6.8) and (6.14)

$$-(2\pi f)^2 \Delta x (\hat{\theta}_A - \hat{\theta}_I) = E_2 \left[\frac{1}{H_2} - \frac{1}{c} \right] - E_1 \left[\frac{1}{H_1} - \frac{1}{c} \right] - N_2 + N_1 \quad (6.16)$$

Now define complex functions of frequency $C_1(f)$ and $C_2(f)$ where

$$C_i(f) \triangleq \left(\frac{1}{H_i(f)} - \frac{1}{c(f)} \right)^{-1}, \quad i = 1, 2 \quad (6.17)$$

Substituting Eq. (6.17) into (6.16) and taking the conjugate square of both sides

$$\begin{aligned} (2\pi f)^4 (\Delta x)^2 |\hat{\theta}_A - \hat{\theta}_I|^2 &= \left| \frac{E_2}{C_2} \right|^2 + \left| \frac{E_1}{C_1} \right|^2 + |N_2|^2 + |N_1|^2 \\ &- \left[\frac{E_1}{C_1} \right]^* \frac{E_2}{C_2} - \frac{E_2 N_2^*}{C_2} + \frac{E_2 N_1^*}{C_2} - \frac{E_1 E_2^*}{C_1 C_2} + \frac{E_1 N_2^*}{C_1} - \frac{E_1 N_1^*}{C_1} \\ &- \frac{N_2 E_2^*}{C_2} + \frac{N_2 E_1^*}{C_1} - N_1^* N_2 + \frac{N_1 E_2^*}{C_2} - \frac{N_1 E_1^*}{C_1} - N_1 N_2^* \quad (6.18) \end{aligned}$$

After ensemble averaging, dropping uncorrelated products, and using the derivative relations for power spectra

$$(\Delta x)^2 S_{(\hat{\theta}_A - \hat{\theta}_I)} = \frac{S_{e_1}}{|C_1|^2} + \frac{S_{e_2}}{|C_2|^2} - 2 \operatorname{Re} \left[\frac{S_{e_1 e_2}}{C_1 C_2} \right] + 2 S_n \quad (6.19)$$

Combining Eqs. (6.2), (6.11), and (6.19)

$$S_{NR} = \frac{S_{e_1} + S_{e_2} - 2 \operatorname{Re}[S_{e_1 e_2}]}{\frac{|c|^2}{|C_1|^2} S_{e_1} + \frac{|c|^2}{|C_2|^2} S_{e_2} - 2 |c|^2 \operatorname{Re} \left[\frac{S_{e_1 e_2}}{C_1^* C_2} \right] + 2 |c|^2 S_n} \quad (6.20)$$

Equation (6.20) is the desired quantity but it can be simplified considerably. Let the frequency response of each leg of the ideal system be chosen as the average of the two legs of the actual system.

$$c(f) = \frac{H_1(f) + H_2(f)}{2} \quad (6.21)$$

Also define the average of the single channel power spectra as S_e where

$$S_e \triangleq \frac{S_{e1} + S_{e2}}{2} \quad (6.22)$$

Finally, let the channel mismatch be specified in terms of a single complex function of frequency $\delta(f)$ where

$$\delta(f) = \frac{H_2(f)}{H_1(f)} - 1 \quad (6.23)$$

From Eqs. (6.17), (6.21), and (6.23)

$$\frac{c}{C_1} = \frac{\delta}{2}$$

$$\frac{c}{C_2} = -\frac{\delta}{2} [1 - \delta + \delta^2 - \dots] \quad (6.24)$$

where the expansion for c/C_2 can be shown to converge for δ anywhere inside the circle $|\delta| = 1$. Using Eq. (6.22) and Eq. (6.24) to simplify Eq. (6.20) gives, after dropping third and higher order terms in δ

$$S_{NR} = \frac{4[1 - \frac{Re(S_{e1}e_2)}{S_e}]}{[|\delta|^2(1 + \frac{Re(S_{e1}e_2)}{S_e}) + 4\frac{|c|^2 S_n}{S_e}]} \quad (6.25)$$

The quantity $\text{Re}(S_{e_1 e_2})/S_e^-$ is a dimensionless function which indicates the amount of angular acceleration which may be measured by finite spatial differencing at a given frequency. It occurs in other developments and will be called the spectral discrete difference and given the symbol α .

$$\alpha(f) \triangleq \text{Re}[S_{e_1 e_2}(f)]/S_e^-(f) \quad (6.26)$$

Combining Eqs. (6.25) and (6.26)

$$S_{NR} = \frac{4(1 - \alpha)}{[|\delta|^2(1 + \alpha) + 4 \frac{|c|^2 S_n}{S_e^-}]} \quad (6.27)$$

This is then the final result. The narrow-band discrete angle signal-to-noise ratio is expressed in terms of a corresponding single channel quantity $S_e^-/|c|^2 S_n$, the complex channel mismatch parameter δ , and a new quantity called the spectral discrete difference. All variables in Eq. (6.27) are functions of frequency.

A simple interpretation of the quantity $\alpha(f)$ may be obtained from the following identity. For any two stationary random signals e_1 and e_2

$$1 - \alpha(f) = \frac{S_{(e_1 - e_2)}}{S_{e_1} + S_{e_2}} \quad (6.28)$$

The quantity $1 - \alpha$ is thus simply the nondimensional size of the PSD of a difference. In the present case $1 - \alpha$ will tend to zero as either Δx or the true angle go to zero. As $1 - \alpha$ goes to zero we may expect poor accuracy in our estimate of $S_{(e_1 - e_2)}$.

If the numerator and denominator in Eq. (6.27) are thought of as dimensionless signal power and dimensionless noise power, then the noise power can be divided. It is composed of a signal-

correlated noise term $|\delta|^2 (1 + \alpha)$ induced by mismatching of channel gain and phase characteristics and an uncorrelated noise term $c^2 S_n / S_e$ which depends only on single channel noise performance.

Some numerical examples are useful to compare the degradation caused by incoherent single channel noise, gain mismatching between channels, and phase mismatching. Suppose one is attempting to measure $S_{\ddot{g}}(f)$ at a frequency where single channel S/N is quite good, say 60 dB. Then

$$\frac{|c|^2 S_n}{S_e} = 10^{-60/10} = 10^{-6}$$

Suppose further that $\alpha = 0.98$. It is shown later that values closer to unity (i.e. worse) than this can be expected to occur routinely in practice. Now consider four situations regarding channel mismatch.

- 1) Perfect matching. Differential narrow-band S/N is

$$S_{NR}|_{dB} = 10 \log \left[\frac{4(1-.98)}{4 \times 10^{-6}} \right] = 43.0 \text{ dB}$$

(17.0 dB worse than single channel)

- 2) Perfect phase matching but 1.5% gain mismatch. In this case

$$\frac{H_2}{H_1} = 1.015 + i0, \quad \delta = \frac{H_2}{H_1} - 1 = 0.015 + i0$$

$$|\delta| = 0.015$$

$$S_{NR}|_{dB} = 10 \log \left[\frac{4(1-.98)}{.015^2(1+.98) + 4 \times 10^{-6}} \right]$$

$$= 22.5 \text{ dB} \quad (37.5 \text{ dB worse than single channel})$$

3) Perfect gain matching but 2° of phase error.

$$\frac{H_2}{H_1} = 1 \angle 2^\circ = .9994 + i .0349$$

$$\delta = -.0006 + i .0349 = .0349 \angle 91.0^\circ$$

$$S_{NR}|_{dB} = 10 \log \left[\frac{4(1-.98)}{.0349^2(1+.98) + 4 \times 10^{-6}} \right]$$

$$= 15.2 \text{ dB (44.8 dB worse than single channel)}$$

4) Phase mismatch of 2° and gain mismatch of 1.5%.

$$\frac{H_2}{H_1} = 1.015 \angle 2^\circ = 1.0144 + i .0354$$

$$\delta = .0144 + i .0354 = .0382 \angle 67^\circ$$

$$S_{NR}|_{dB} = 10 \log \left[\frac{4(1-.98)}{.0382^2(1+.98) + 4 \times 10^{-6}} \right]$$

$$= 14.4 \text{ dB (45.6 dB worse than single channel)}$$

The latter values for noise and mismatch parameters are probably typical for a differential sensing system where each leg consists of a piezoelectric accelerometer, charge converter, voltage amplifier, anti-aliasing filter, and A/D converter. If, in addition, individual channels are passed through analog tape, the situation could be considerably worse.

Eq. (6.27) has been used to compute a set of plots which may be used for quick reference in setting up an experiment where the PSD of a difference quantity is to be measured. These are displayed as Figures 38 through 41.

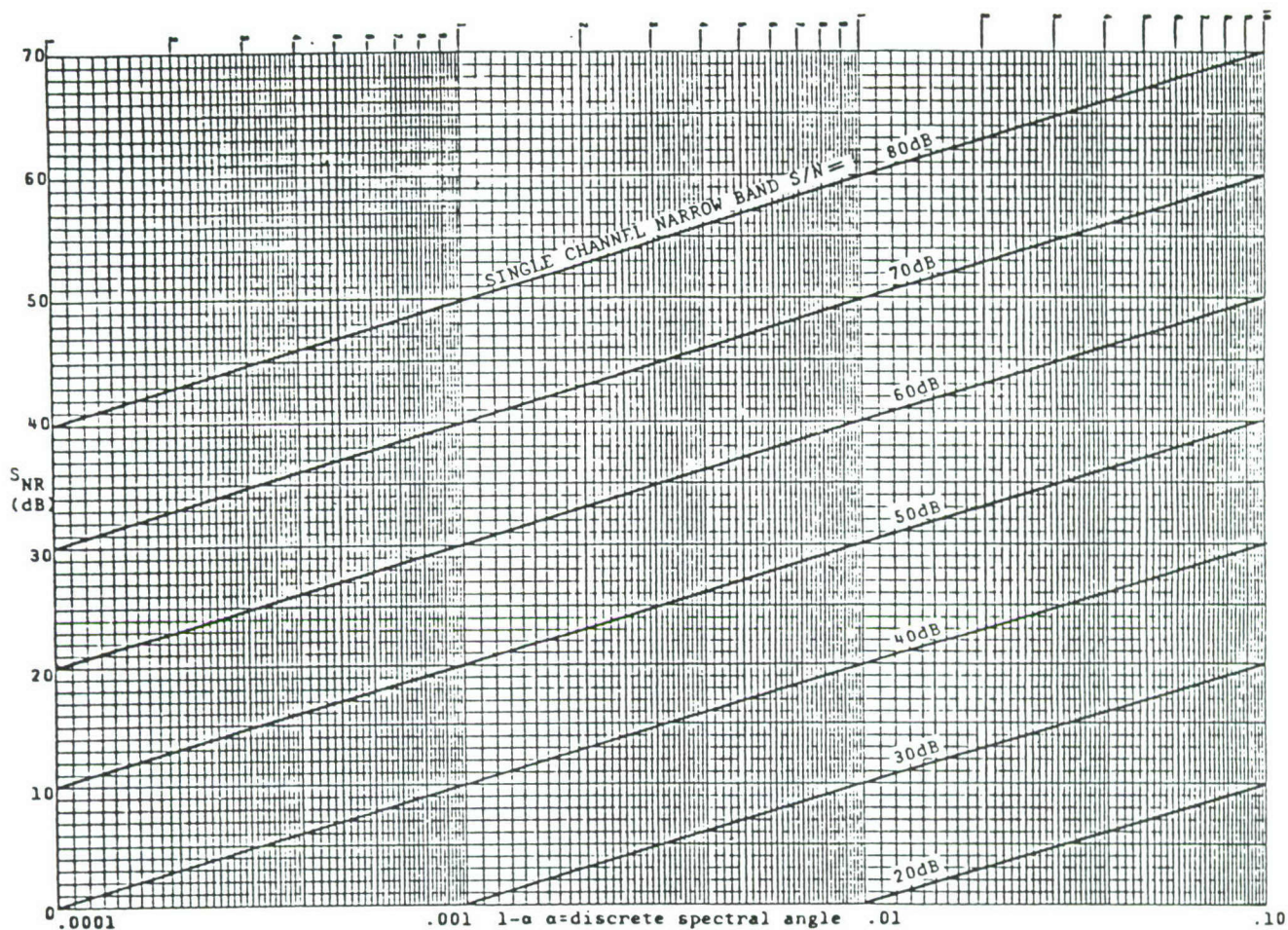


Figure 38 Narrow-Band Signal-to-Noise Ratio of Angular Acceleration Estimate Obtained by Differencing of Linear Accelerations
 $|\delta| = 0.000$

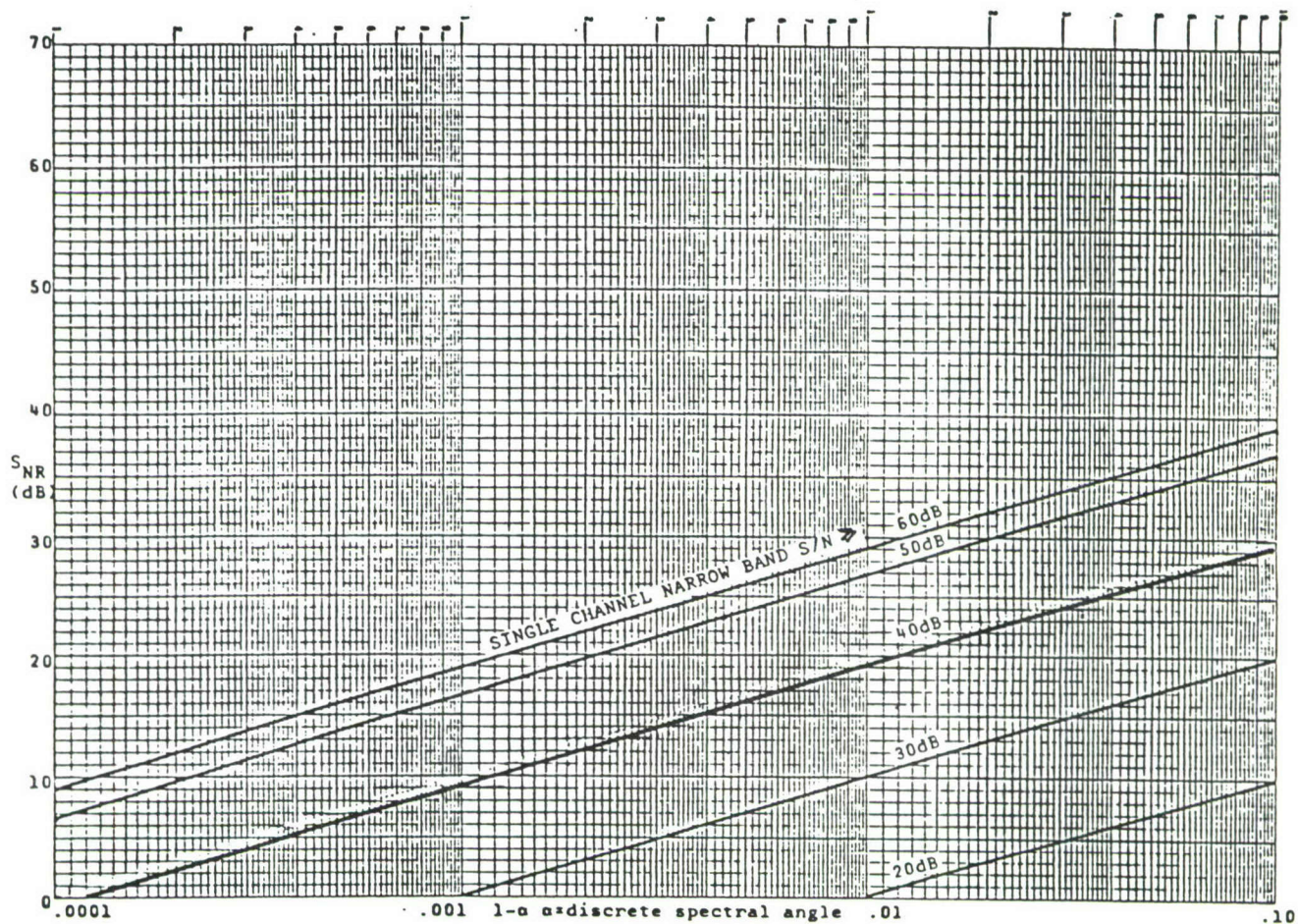


Figure 39 Narrow-Band Signal-to-Noise Ratio of Angular Acceleration Estimate Obtained by Differencing of Linear Accelerations
 $|\delta| = 0.005$

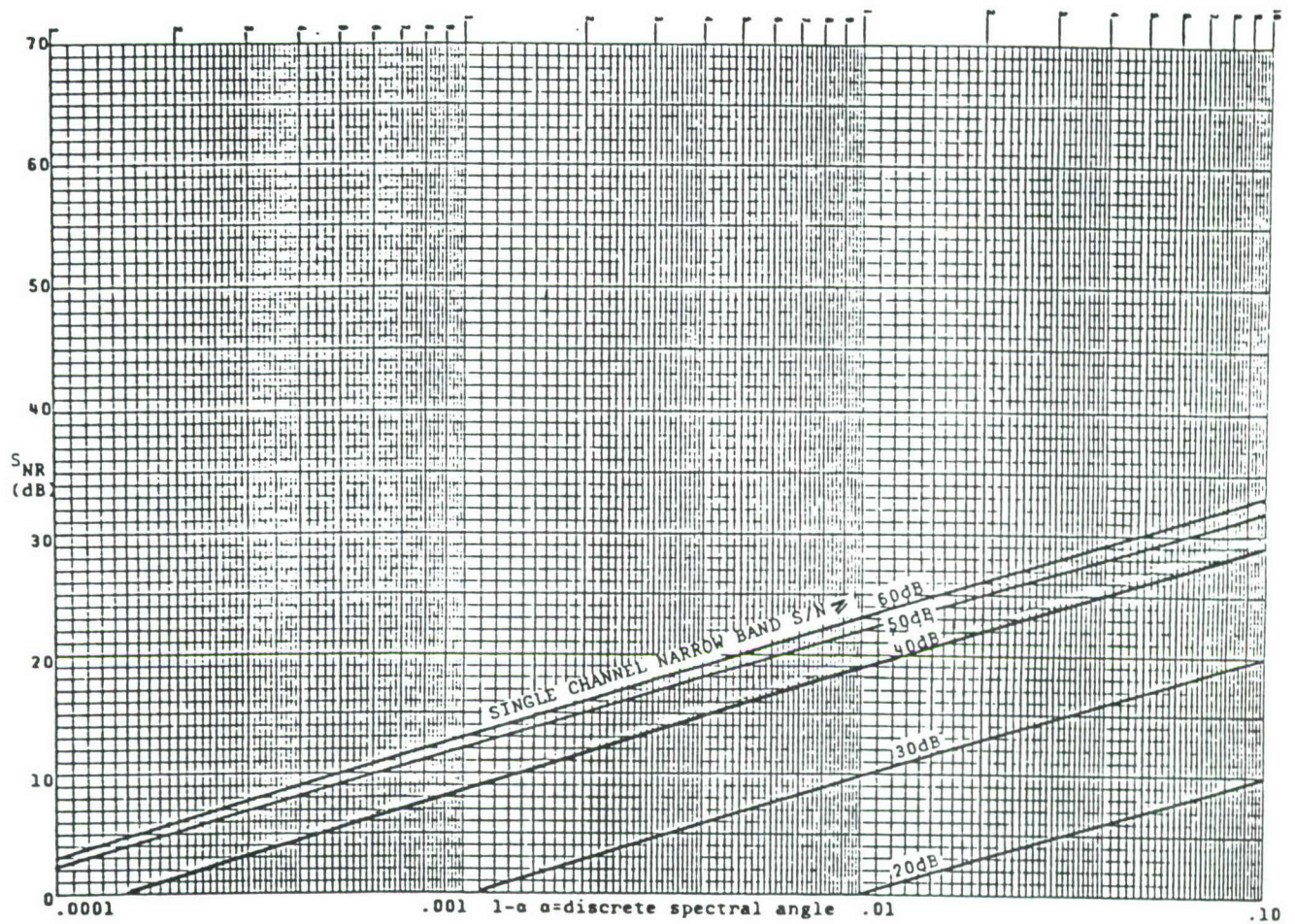


Figure 40 Narrow-Band Signal-to-Noise Ratio of Angular Acceleration Estimate Obtained by Differencing of Linear Accelerations
 $|\delta| = 0.01$

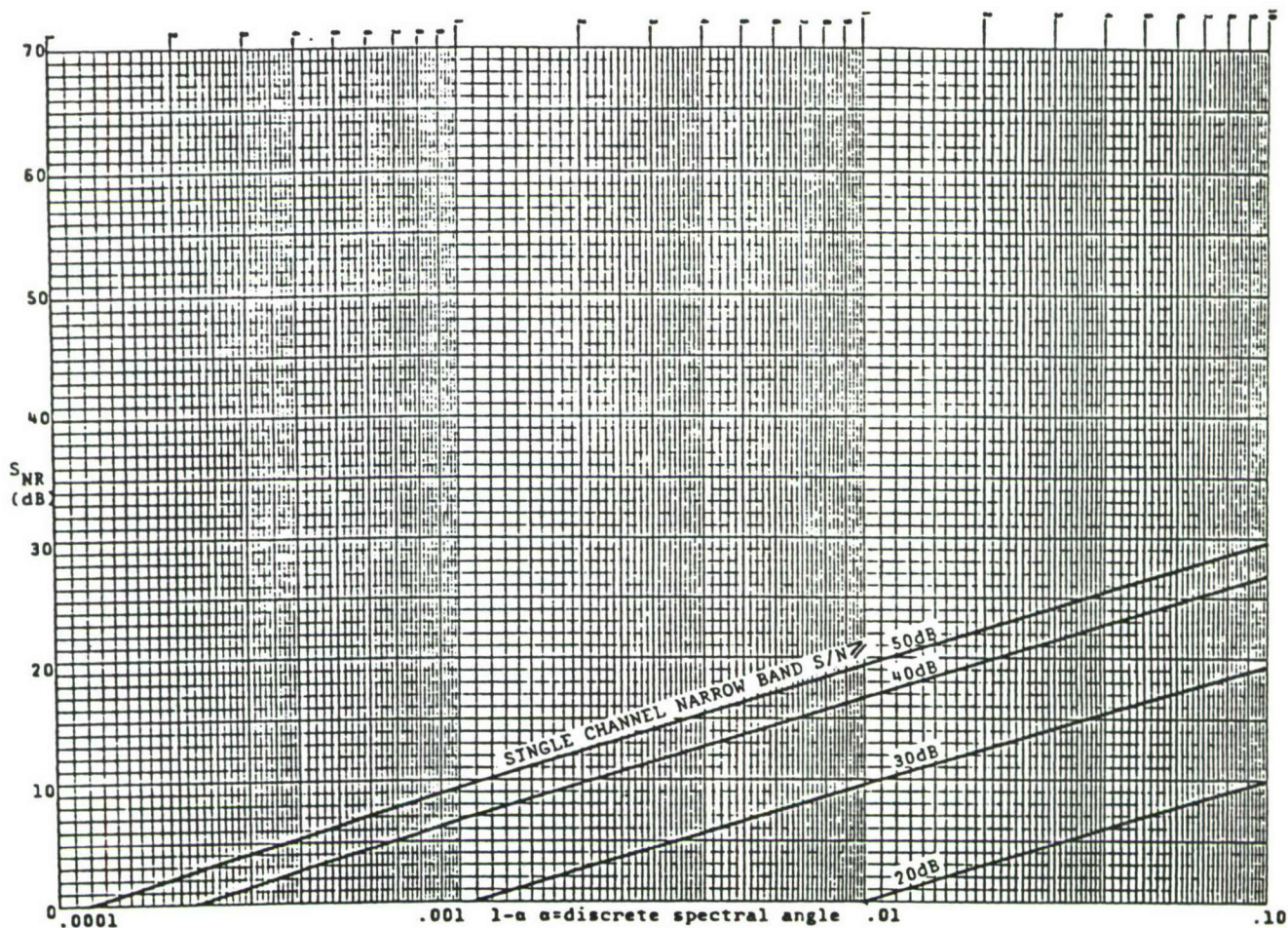


Figure 41 Narrow-Band Signal-to-Noise Ratio of Angular Acceleration Estimate Obtained by Differencing of Linear Accelerations
 $|\delta| = 0.015$

The transducer separation distance Δx does not appear explicitly in Eq. (6.27). Its effect is included in the spectral discrete difference α . As Δx is increased, a difference appears between e_1 and e_2 which causes α to drop below unity. It is shown in a later paragraph that α may be very close to unity for practical cases. It should be noted that while $\alpha(f)$ is a useful concept for quantifying errors, its measurements is not necessarily straightforward. Experiments have indicated that accurate determination of $\alpha(f)$ is at least as difficult as the measurement of $S_\theta(f)$ itself.

One additional plot is provided as Figure 42 to aid in estimating narrowband signal-to-noise ratio at a given frequency. It is intended for use with piezoelectric accelerometers where essentially all of the incoherent noise power is below 100 Hz and a high pass filter is incorporated to suppress very low frequency noise below about 1 Hz. A typical noise power spectrum for this type of instrument is of the form

$$S_n(f) = A f^{-m} \quad (6.29)$$

If the total noise power σ_n^2 between $f_{\min} = 1$ Hz and $f_{\max} = 100$ Hz is known, the quantity A in Eq. (6.29) can be obtained in terms of σ_n^2 and m from

$$\sigma_n^2 = \int_{f_{\min}}^{f_{\max}} S_n(f) df \quad (6.30)$$

From Eqs. (6.29) and (6.30)

$$A = \frac{(1 - m) \sigma_n^2}{(f_{\max}^{1-m} - f_{\min}^{1-m})} \quad (6.31)$$

Measurements of noise spectra for a variety of piezoelectric transducers have indicated that $m = 1.32$ is a representative value. This was used in the calculation of Figure 42.

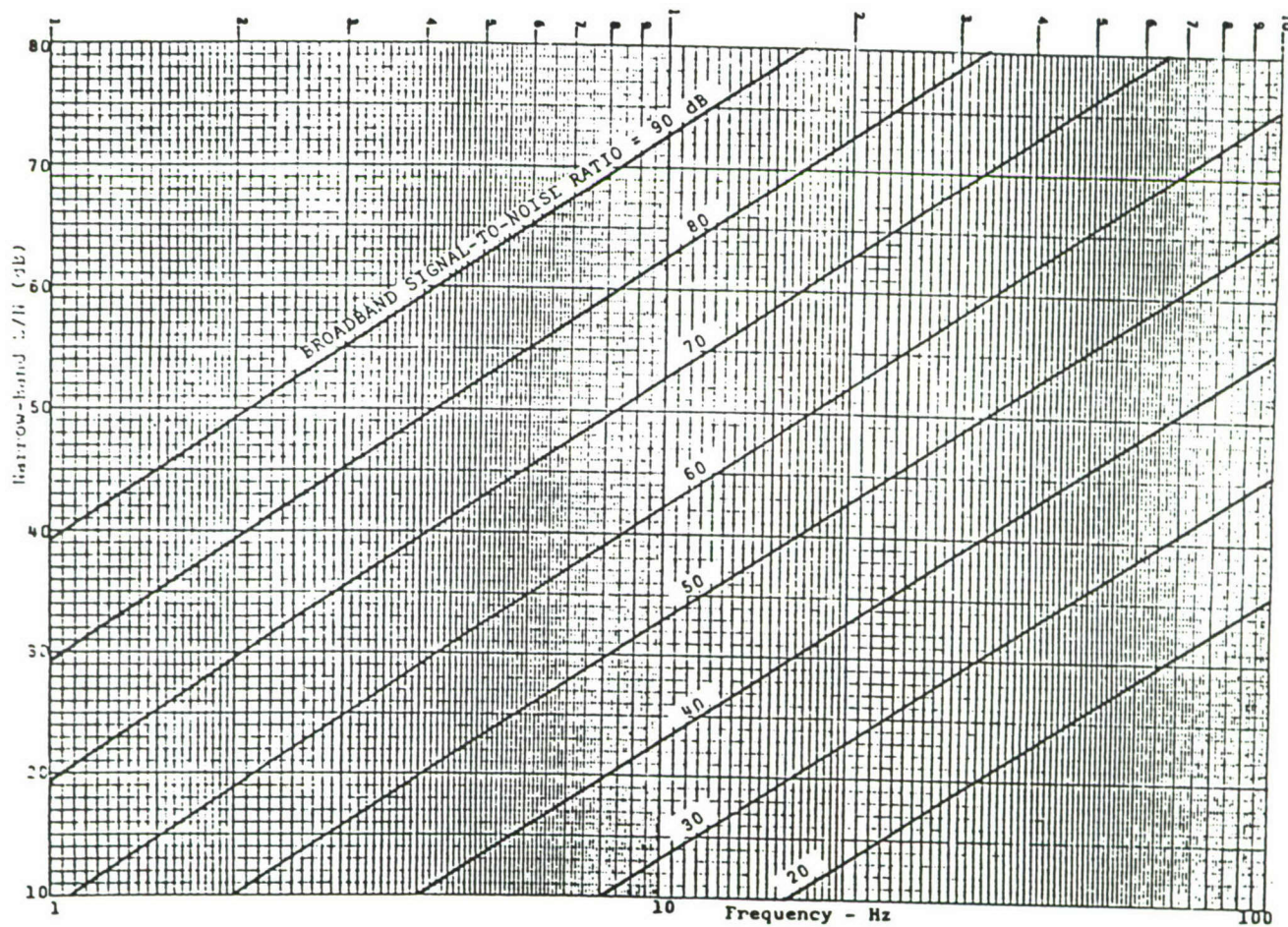


Figure 42 Typical Narrow-Band Signal-to-Noise for Piezoelectric Accelerometers

For an infinite plate, the mechanical impedance (ratio of input force to response velocity at a given frequency) can be shown to be independent of frequency [21]. Thus, for a structure built up of plate elements, it is reasonable to expect the envelope of the PSD acceleration response under random excitation to increase with frequency as ω^2 once the region of reasonably high modal density has been reached. The signal power spectrum will therefore be approximated as

$$S_{e_i} = B f^k, \quad i = 1, 2 \quad (6.32)$$

with $k = 2$. Then if σ_e^2 is the total signal power in the band of interest, B can be found from

$$\sigma_{e_i}^2 = \int_{f_{\min}}^{f_{\max}} B f^k \quad (6.33)$$

or

$$B = \frac{(k + 1) \sigma_{e_i}^2}{(f_{\max}^{k+1} - f_{\min}^{k+1})} \quad (6.34)$$

In the current case $f_{\min} \approx 1$ Hz and $f_{\max} = 100$ Hz, so with negligible error

$$B = \frac{(k + 1) \sigma_{e_i}^2}{f_{\max}^{k+1}} \quad (6.35)$$

In addition, the use to which $c^2 S_n / S_{e_i}$ will be put does not require extreme accuracy so A will be approximated as

$$A \approx \frac{(m - 1) \sigma_n^2}{f_{\min}^{1-m}} \quad (6.36)$$

Combining Eqs. (6.29), (6.31), (6.32), and (6.34)

$$\frac{S_e}{c^2 S_n} = \left(\frac{k+1}{m-1}\right) \left(\frac{f_{\min}}{f_{\max}}\right) \left(\frac{f}{f_{\min}}\right)^m \left(\frac{f}{f_{\max}}\right)^k \left(\frac{\sigma_e^2}{c^2 \sigma_n^2}\right) \quad (6.37)$$

The quantity $\sigma_e^2/c^2\sigma_n^2$ is the single channel broadband signal-to-noise ratio which may be quickly measured or estimated from experience in setting up an experiment. Eq. (6.37) simply quantifies the effect of the noise power spectrum being shaded towards low frequencies and the signal power spectrum towards high frequencies. From Figure 42 it may be observed that, for the values of m , k , f_{\min} , and f_{\max} used, the narrowband signal-to-noise ratio is worse than the broadband below about 34 Hz. At 5 Hz, which is still high enough to be important for airborne optical system applications, the narrowband S/N is approximately 27 dB worse.

6.1.2 Coherence Function for Angular Frequency Response

In establishing a mathematical model from test data or for verifying an analytical model, some form of standardized, measurable input-output relationship is needed. For linear structures this is usually taken as the frequency response. If angular response can be reliably measured by differencing of linear signals then one important use of these angular data would be to compute the frequency response of angular degrees of freedom to spatially fixed force inputs. The ordinary coherence function [29] is a standard method of checking the consistency of frequency responses measured by the DFT cross-spectral averaging method. In this section, an expression is obtained for the coherence function of a differentially measured angular frequency response in terms of the coherence functions associated with the individual channels.

Suppose $\ell(t)$ with Fourier transform $L(f)$ is a force input producing two acceleration outputs $\ddot{y}_1(t)$ and $\ddot{y}_2(t)$ which are differenced to estimate $\ddot{\theta}(t)$ according to Eq. (6.6). The definition of coherence for this output quantity is

$$\gamma_{\ell\ddot{\theta}} = \frac{|S_{\ell\ddot{\theta}}|^2}{S_{\ddot{\theta}} S_{\ell}} \quad (6.38)$$

Relating this to single channel quantities

$$(\Delta x) S_{\ell\ddot{\theta}} = \overline{L^*(\ddot{Y}_2 - \ddot{Y}_1)} = \overline{L^* \ddot{Y}_2} - \overline{L^* \ddot{Y}_1} = S_{\ell\ddot{y}_2} - S_{\ell\ddot{y}_1} \quad (6.39)$$

$$(\Delta x)^2 S_{\ddot{\theta}} = \overline{(\ddot{Y}_2 - \ddot{Y}_1)^* (\ddot{Y}_2 - \ddot{Y}_1)} = S_{\ddot{y}_1} + S_{\ddot{y}_2} - 2\text{Re}(S_{\ddot{y}_1\ddot{y}_2}) \quad (6.40)$$

Combining Eqs. (6.38), (6.39), and (6.40)

$$\gamma_{\ell\ddot{\theta}} = \frac{|S_{\ell\ddot{y}_2}|^2 + |S_{\ell\ddot{y}_1}|^2 - 2\text{Re}(S_{\ell\ddot{y}_2}^* S_{\ell\ddot{y}_1})}{S_{\ell}[S_{\ddot{y}_1} + S_{\ddot{y}_2} - 2\text{Re}(S_{\ddot{y}_1\ddot{y}_2})]} \quad (6.41)$$

The effect of channel sensitivity mismatching is to introduce signal-correlated noise. Since this will not be detected by the coherence function there is nothing to be gained by considering $H_1 \neq H_2$ in simplifying Eq. (6.41). So let $H_1 = H_2 = c$. Then since

$$|S_{\ell y_i}|^2 = |c|^2 |S_{\ell e_i}|^2 \quad (6.42)$$

$$S_{\ell\ddot{y}_i} = c S_{\ell e_i}$$

etc.

we have from Eq. (6.41) after canceling $|c|^2/|c|^2$

$$\gamma_{\ell\ddot{\theta}}^2 = \frac{|S_{\ell e_1}|^2 + |S_{\ell e_2}|^2 - 2\text{Re}(S_{\ell e_2}^* S_{\ell e_1})}{S_{\ell}[S_{e_1} + S_{e_2} - 2\text{Re}(S_{e_1 e_2})]} \quad (6.43)$$

Introducing the definitions of average signal power spectrum S_e^- from Eq. (6.22) and spectral discrete difference α from Eq. (6.26) into Eq. (6.43)

$$\gamma_{l\theta}^{2..} = \frac{1}{2(1-\alpha)} \left[\frac{|S_{le_1}|^2}{S_l S_e^-} + \frac{|S_{le_2}|^2}{S_l S_e^-} - \frac{2\text{Re}(S_{le_2}^* S_{le_1})}{S_l S_e^-} \right] \quad (6.44)$$

Combining this with the definition of the single channel coherence function

$$\gamma_{le_i}^2 = \frac{|S_{le_i}|^2}{S_l S_{e_i}} \quad (6.45)$$

$$\gamma_{l\theta}^{2..} = \frac{1}{2(1-\alpha)} \left[\frac{S_{e_2}}{S_e^-} \gamma_{le_2}^2 + \frac{S_{e_1}}{S_e^-} \gamma_{le_1}^2 - \frac{2\text{Re}(S_{le_2}^* S_{le_1})}{S_l S_e^-} \right] \quad (6.46)$$

Since there is no reason to expect the individual channel coherences to be different any more than the individual channel noises, they will be assumed equal, that is

$$\gamma_{le_1}^2 = \gamma_{le_2}^2 = \gamma_{le}^2 \quad (6.47)$$

Then

$$\gamma_{l\theta}^{2..} = \frac{1}{(1-\alpha)} \left[\gamma_{le}^2 - \frac{\text{Re}(S_{le_1} S_{le_2}^*)}{S_l S_e^-} \right] \quad (6.48)$$

The second term in the brackets is a real, dimensionless function of frequency with both numerator and denominator being products of two power spectra (except for the $\text{Re}(\cdot)$ operator). In these respects, it resembles an ordinary coherence function. It will therefore be given the symbol Γ

$$\Gamma^2 \triangleq \frac{\text{Re}(S_{le_1} S_{le_2}^*)}{S_l S_e^-} \quad (6.49)$$

With this notation

$$\gamma_{\ell\theta}^2 = \frac{\gamma_{\ell e}^2 - \Gamma^2}{1-\alpha} \quad (6.50)$$

This is the desired result although it will likely be a poorly conditioned calculation because $\gamma_{\ell e}^2$, Γ^2 and α will all be close to unity for small transducer separation.

6.1.3 Typical Value of Spectral Discrete Difference

The narrowband signal-to-noise ratio S_{NR} and input-output coherence function $\gamma_{\ell\theta}$ for angular responses obtained by differencing have been shown to depend on the spectral discrete difference α . Both of these quantities are of particular interest for the frequency range around a structural resonance. In this section a simple expression is obtained for α in terms of the mode shape associated with a particular resonant frequency. This is useful in order to gain some idea of the difficulty to be expected in measuring α , or more specifically $1-\alpha$.

Suppose a structure is excited at a single degree of freedom by a load $\ell(t)$ with a Fourier transform $L(f)$. Let L be broad band with unit amplitude in the frequency range around an isolated resonance at $f_r = \omega_r/2\pi$. Assume perfect coherence between ℓ and each of the responses \ddot{y}_1 and \ddot{y}_2 . Then

$$\begin{aligned} Y_1 &= H_{01} L \\ Y_2 &= H_{02} L \end{aligned} \quad (6.51)$$

The subscript 0 in Eq. (6.51) refers to the driving point. Then

$$S_{y_1 y_2} = \overline{Y_1^* Y_2} = H_{01}^* H_{02} S_\ell \quad (6.52)$$

By the assumption of perfect coherence

$$S_{y_1} = |H_{01}|^2 S_\ell$$

$$S_{y_2} = |H_{02}|^2 S_\ell$$
(6.53)

Then from Eqs. (6.26), (6.52), and (6.53)

$$\alpha = \frac{2\text{Re}(H_{01}^* H_{02})}{|H_{01}|^2 + |H_{02}|^2}$$
(6.54)

If f_r is the resonant frequency associated with a normal mode shape vector $\psi^{(r)}$, the frequency response has the form

$$H_{0i}(\omega) = \sum_{r=1}^N \frac{\psi_0^{(r)} \psi_i^{(r)}}{(\omega^2 - \omega_r^2) + j 2\zeta_r \omega \omega_r}$$
(6.55)

where

$$\psi_0^{(r)}, \psi_i^{(r)} = 0 \text{ and } i\text{'th entries in mode vector which is}$$

assumed to be real and normalized to unit

modal mass

$$\zeta_r = \text{modal damping ratio}$$

For very light damping H_{0i} can be approximated by the r th term of its modal series for $\omega \approx \omega_r$. Using this approximation along with Eq. (6.54)

$$\alpha = \frac{2\left(\frac{\psi_2}{\psi_1}\right)}{1 + \left(\frac{\psi_2}{\psi_1}\right)^2}$$
(6.56)

$$1 - \alpha = \frac{\left(1 - \frac{\psi_2}{\psi_1}\right)^2}{1 + \left(\frac{\psi_2}{\psi_1}\right)^2}$$
(6.57)

Figure 43 shows a plot of dimensionless signal $1-\alpha$ versus $\frac{\psi_2}{\psi_1} - 1$. This result shows the problem inherent in angular measurement since α departs from unity much more slowly than does ψ_2/ψ_1 . To obtain $1-\alpha$ of even .01 requires $\psi_2/\psi_1 = 1.153$. Unless points 1 and 2 on the mode shape ψ fall close to a node point, one can expect the differential signal-to-noise ratio to be markedly worse than that for individual channels.

6.1.4 Flexural Error

At higher frequencies, the magnitude of uncorrelated noise will diminish to insignificance. However, a new limitation will arise. An error limit will be established by the discrete angle approximation itself. In this section, a chart is presented for choosing a separation distance for a given error.

At high frequencies, the bending mode shapes of plates or beams are affected by boundary conditions only near the edges. For a time-random load applied well inside the boundary the high frequency behavior can be interpreted in terms of an infinite plate or beam. The power input at a frequency f will be propagated outward by waves of length λ where λ and f are related by the dispersion relation for bending of beams or plates. If the acceleration distribution is one dimensional (straight crested waves in a plate) with wavelength λ

$$\ddot{y}(x) = D \sin \frac{2\pi x}{\lambda} \quad (6.58)$$

The finite difference approximation to angular acceleration is

$$\hat{\theta} = \frac{\ddot{y}(x_2) - \ddot{y}(x_1)}{\Delta x} = \frac{D}{\Delta x} \left[\sin\left(\frac{2\pi x_2}{\lambda}\right) - \sin\left(\frac{2\pi x_1}{\lambda}\right) \right] \quad (6.59)$$

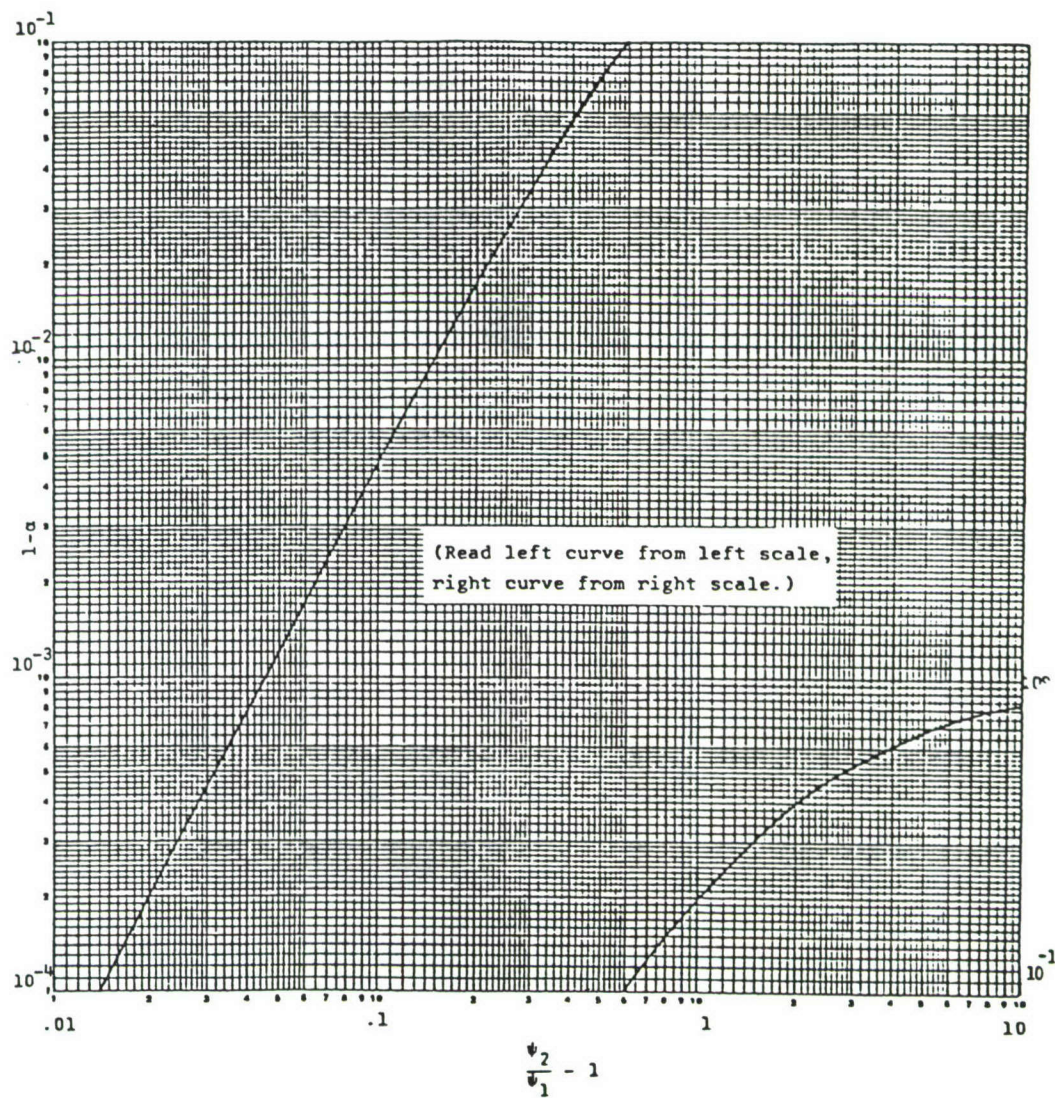


Figure 43 Dimensionless Difference Signal $1-\alpha$ vs. Mode Shape ψ

then if

$$\begin{aligned}x_2 &= x_c + \frac{\Delta x}{2} \\x_1 &= x_c - \frac{\Delta x}{2}\end{aligned}\tag{6.60}$$

Equation (6.59) can be rearranged as

$$\hat{\theta} = \frac{2D}{\Delta x} \cos\left(\frac{2\pi x_c}{\lambda}\right) \sin\left(\frac{\pi \Delta x}{\lambda}\right)\tag{6.61}$$

The true angular acceleration at $x = x_c$ is

$$\ddot{\theta}(x_c) = \frac{d\ddot{y}}{dx} = \frac{2\pi D}{\lambda} \cos\frac{2\pi x_c}{\lambda}\tag{6.62}$$

From Eqs. (6.61) and (6.62)

$$\hat{\theta} = \left[\frac{\sin \frac{\pi \Delta x}{\lambda}}{(\frac{\pi \Delta x}{\lambda})} \right] \ddot{\theta}\tag{6.63}$$

The quantity in brackets is the familiar modulation function which occurs when sampling is performed by averaging over a finite interval. Call it $M(\cdot)$.

$$M(\cdot) \triangleq \frac{\sin(\cdot)}{(\cdot)}\tag{6.64}$$

For straight crested waves in a plate, the dispersion relation is the same as for a simple beam if the Poisson effect is neglected [30].

$$c_b = \sqrt{2\pi f \kappa c_\ell}\tag{6.65}$$

c_b = phase velocity of bending waves

f = frequency

κ = radius of gyration of section

c_ℓ = extensional wave speed $\sqrt{E/\rho}$

If M_0 is the minimum allowable value of the modulation function ($1 - M_0$ = fractional error), then Eqs. (6.63), (6.64), and (6.65) may be combined to find the frequency above which this error will be exceeded.

$$f = \frac{2}{\pi} c_\ell \kappa \left[\frac{M^{-1}(M_0)}{\Delta x} \right]^2 \quad (6.66)$$

Arranging in dimensionless form

$$\left(\frac{f \kappa}{c_\ell} \right) = \frac{2}{\pi} \left(\frac{\Delta x}{\kappa} \right)^{-2} [M^{-1}(M_0)]^2 \quad (6.67)$$

This relation is plotted in Figure 44 for various dimensionless separation distance ($\Delta x/\kappa$). For airborne optical system work the areas where one would be likely to place differential transducers are usually on fairly thick, stiff components. Thus, the upper half of Figure 44 is probably the most relevant.

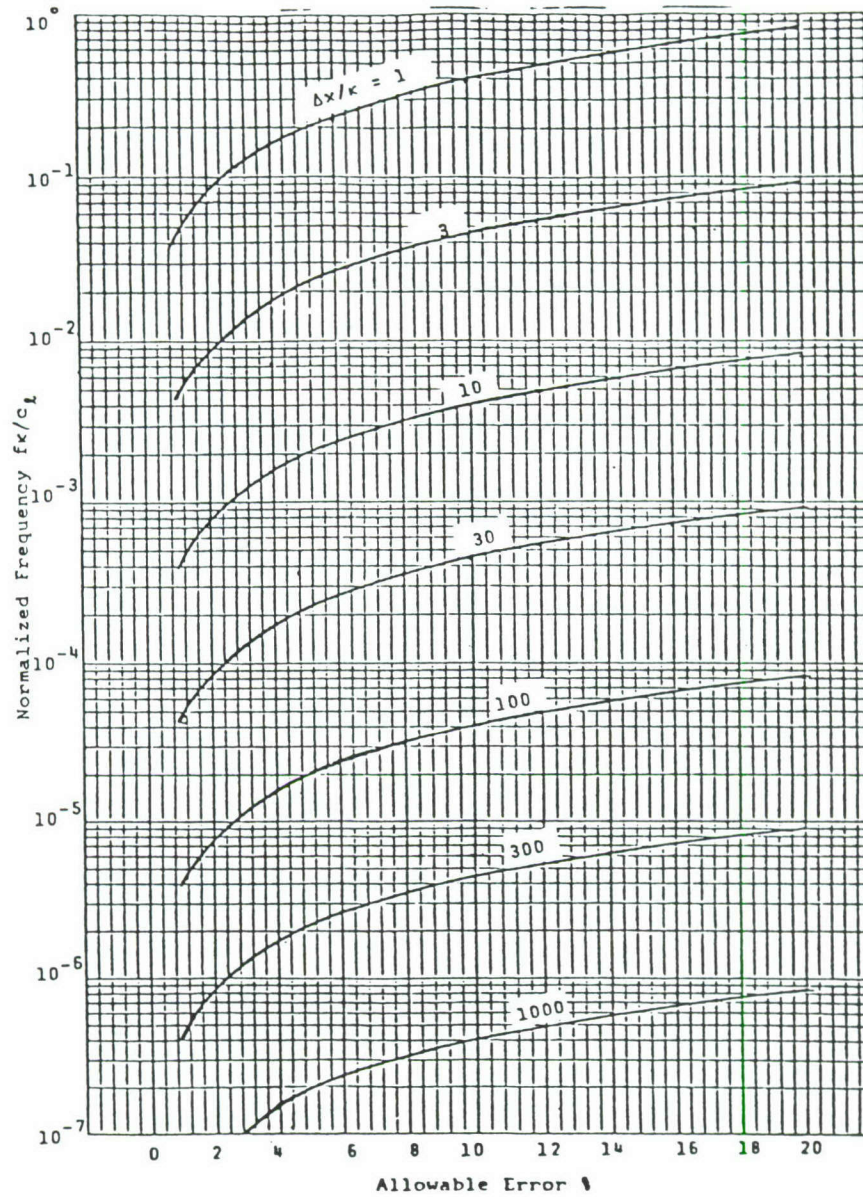


Figure 44 Upper Frequency Limit for Differential Sensing of Angular Vibration

6.2 EXAMPLES

In this section some examples are presented to illustrate the practical meaning of the error sources described in Section 6.1. In addition, a method is demonstrated for reducing the effect of one error source, the interchannel gain and phase mismatch.

6.2.1 Use of Error Formulas

In Section 4.6 of this report, an experiment is described which involves measurement of the PSD of angular acceleration at an interior point on a uniform plate. Error bands for this measurement situation are calculated here as an example. The following data are known:

Plate characteristics

material	aluminum
thickness	0.090 in.
Young's modulus	10×10^6 lbf/in. ²
mass density	2.6×10^{-4} lbf-sec. ² /in. ⁴
Poisson's ratio	0.3

Measurement requirements

minimum frequency	100 Hz.
maximum frequency	1000 Hz.
transducer separation	1.40 in.

It is desired to estimate the differential narrowband signal/noise ratio at the low end of the frequency band and the flexural error at the high end.

The noise spectrum of the accelerometer/signal conditioning system was not measured. However, the transducer sensitivity was less than that of transducer A (see Figure 45) by about a factor of 100. If the curve for transducer A is extrapolated to 100 Hz. and noise power/Hz. is assumed to be proportional to (sensitivity)⁻², then we may estimate the noise power spectrum at 100 Hz. to be about $2.0 \times 10^{-9} \frac{G^2}{Hz.}$ or $3.0 \times 10^{-4} (in/sec.)^2/Hz.$ The r.m.s. force input was about

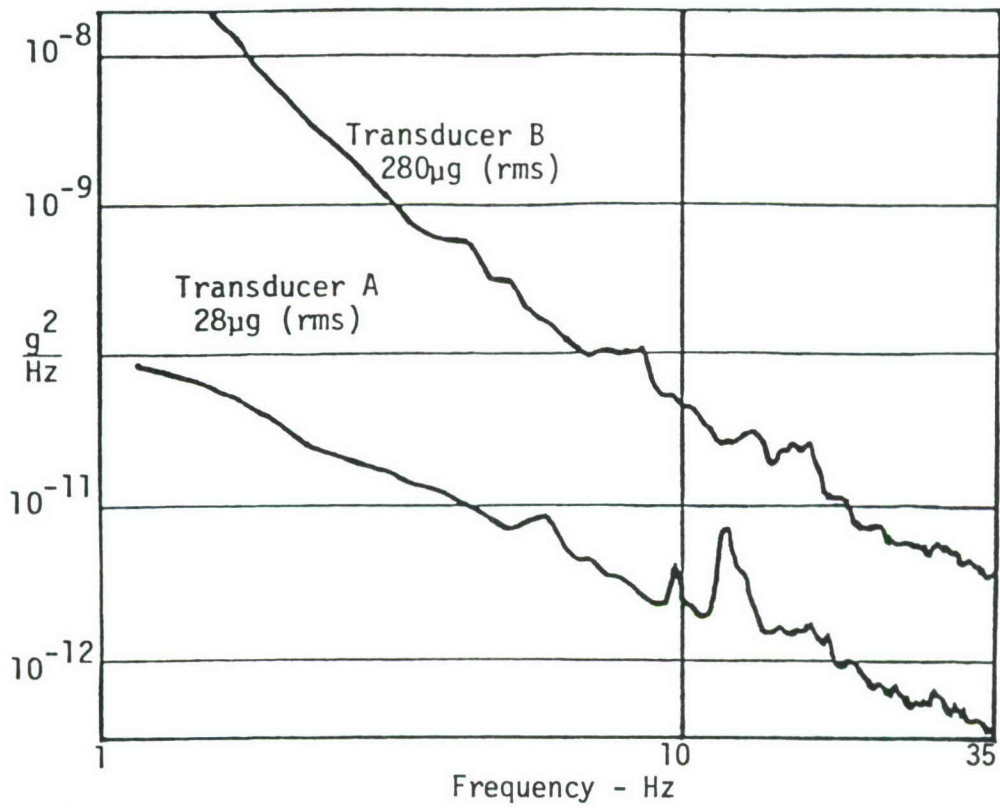


Figure 45 Noise PSD for Two Different Piezoelectric Accelerometers

0.3 lbf. If the force p.s.d. is taken to be constant over 0-1000 Hz. and the plate is approximated as infinite in extent with the thickness and material properties of the actual plate, then the velocity p.s.d. at the drive point is [21]

$$S_{\dot{y}} = \frac{S_f}{(8\rho h \kappa c_\ell)^2}$$

or

$$S_{\ddot{y}} = \frac{S_f \omega^2}{(8\rho h \kappa c_\ell)^2}$$

Computing the signal power/Hz. at 100 Hz.

$$\kappa = \frac{h}{2\sqrt{3}} = \frac{0.090}{2\sqrt{3}} = 0.0260 \text{ in.}$$

$$c_\ell = \sqrt{\frac{E}{\rho}} = \sqrt{\frac{10 \times 10^6}{2.6 \times 10^{-4}}} = 1.96 \times 10^5 \text{ in/sec.}$$

$$S_f = \frac{(0.3)^2 \text{ lbf}^2}{1000 \text{ Hz.}} = 9.0 \times 10^{-5} \frac{\text{lbf}^2}{\text{Hz.}}$$

$$\begin{aligned} S_{\ddot{y}} &= \frac{S_f \omega^2}{(8\rho h \kappa c_\ell)^2} = \frac{(9.0)(10^{-5})[(2\pi)(100)]^2}{[(8)(2.6)(10^{-4})(0.090)(0.026)(1.96)(10^5)]^2} \\ &= 39.0 \frac{(\text{in/sec.}^2)^2}{\text{Hz.}} \end{aligned}$$

The worst case single channel narrowband noise/signal ratio is thus estimated as

$$\frac{|c^2| S_n}{S_e} = \frac{3.0 \times 10^{-4}}{39.0} = 7.69 \times 10^{-6} \text{ (-52 dB)}$$

The sensing channels were gain-matched at 300 Hz. prior to use. so residual mismatch will be taken as 0.5% gain and 1.5° phase. The mismatch parameter δ is then

$$\begin{aligned}\delta &= 1.005 \angle 1.5^\circ - 1.000 \\ &= 0.0267 \angle 80.0^\circ \\ |\delta|^2 &= 7.14 \times 10^{-4}\end{aligned}$$

The spectral discrete difference α could be estimated as anything between -1 and +1 since there is no way of knowing where the accelerometers will be placed relative to the shapes of modes which are resonant around 100 Hz. In deriving Eq. (4.78) (which is the prediction this measurement was intended to verify) it was assumed that mode shapes would vary sinusoidally in space in directions parallel to the plate edges. We therefore assume that the mid-point between the accelerometers is halfway between the peak and zero of a sine wave. Some modes will be worse (closer to the peak) and some will be better. The most important modes, i.e. those contributing the most to $\langle \theta^2 \rangle_t$, will be those where the transducers are closer to a node. The midway assumption should therefore be conservative. The appropriate wave-length λ can be found from the dispersion relation, Eq. (6.65), and the elementary relation $c_b = f\lambda$.

$$\lambda = \frac{2\pi c_b}{f}$$

$$\lambda = \frac{(2\pi)(0.026)(1.96)(10^5)}{100}$$

$$\lambda = 17.9 \text{ in.}$$

The values of k_{xi} (see Section 4.4) corresponding to the transducer locations are thus $\frac{\pi}{8} \pm \frac{2\pi\Delta x}{2\lambda}$ or

$$\frac{\pi}{8} \pm \frac{2\pi\Delta x}{2\lambda} = \frac{\pi}{8} \pm \frac{(2\pi)(1.40)}{(2)(17.9)} = 0.146, 0.638 \text{ radians}$$

The representative value of α is then, by Eq. (6.56)

$$\alpha = \frac{2 \left[\frac{\sin(.638)}{\sin(.146)} \right]}{1 + \left[\frac{\sin(.638)}{\sin(.146)} \right]^2} = 0.461$$

The worst case differential narrowband signal/noise ratio is, by Eq. (6.27)

$$\begin{aligned} S_{NR} &= \frac{4(1-\alpha)}{|\delta|^2 (1+\alpha) + 4 \left(\frac{|c|^2 S_n}{S_e} \right)} \\ &= \frac{4(1 - 0.461)}{[(7.14)(10^{-4})(1 + 0.461) + 4(7.69)(10^{-6})]} \\ &= 2.01 \times 10^3 \text{ (33.0 dB)} \end{aligned}$$

This level of signal fidelity should be quite adequate for the intended purpose.

To compute worst-case flexural error, we calculate the wavelength associated with the highest frequency of interest, 1000 Hz.

$$\lambda = \frac{2\pi c \ell}{f} = \frac{(2\pi)(0.026)(1.96)(10^5)}{1000} = 5.66 \text{ in.}$$

Then, using Eq. (6.63)

$$\frac{\hat{\theta}}{\theta} = \frac{\sin\left(\frac{\pi \Delta x}{\lambda}\right)}{\frac{\pi \Delta x}{\lambda}} = \frac{\sin\left[\frac{\pi(1.40)}{5.66}\right]}{\left[\frac{\pi(1.40)}{5.66}\right]} = 0.902$$

The effective reduction will actually be even less than 10% since the dominant contributions to $\theta_{r.m.s.}$ come from frequencies well below 1000 Hz. (see Figure 17).

6.2.2 Effect of Single Channel Noise

To illustrate the effect of uncorrelated noise in individual accelerometer channels, a simple experiment was performed.

A slender aluminum I-beam was suspended in a horizontal position by long cords to simulate a free-free condition. It was vibrated by a small shaker driving in the beam's weak direction at a point slightly off midspan. A time-random force

signal was input and rotations at various points along the beam's length were sensed by analog differencing of accelerometer outputs. Frequency responses between angular acceleration output and point force input were measured and curve-fitted to estimate the first few normal mode shapes in slope coordinates [16]. The experiment was repeated with two different accelerometer pairs of substantially different noise performance. Measured noise power spectra for both are shown in Figure 45. The force input was identical in both cases and was much smaller than one would normally use in modal survey work in order to emphasize the effect of random noise in the sensor channels. The shape of the second bending mode as measured by each transducer pair is shown in Figure 46. Since the beam was uniform, any deviation from a smooth curve indicates error in the measurement. The improvement with lower noise transducers is evident.

6.2.3 Effect of Interchannel Gain and Phase Mismatch

The derivation leading to Eq. (6.27) suggests a method for removing, or at least reducing, signal-correlated error due to mismatching. The method is basically to measure the complex (i.e., both gain and phase) mismatch between the two sensing channels as a function of frequency and use it to apply a small correction to the data coming from one channel prior to differencing. Mathematically, if we define a correction function H_c as $H_c(f) = H_2(f)/H_1(f)$, then the corrected PSD of $(e_2 - e_1)$ may be calculated either as

$$S_{(e_2 - e_1)} = S_{e_2} + |H_c|^2 S_{e_1} - 2 \operatorname{Re}[H_c S_{e_2 e_1}] \quad (6.68)$$

or as

$$S_{(e_2 - e_1)} = \overline{|E_2 - H_c E_1|^2} \quad (6.69)$$

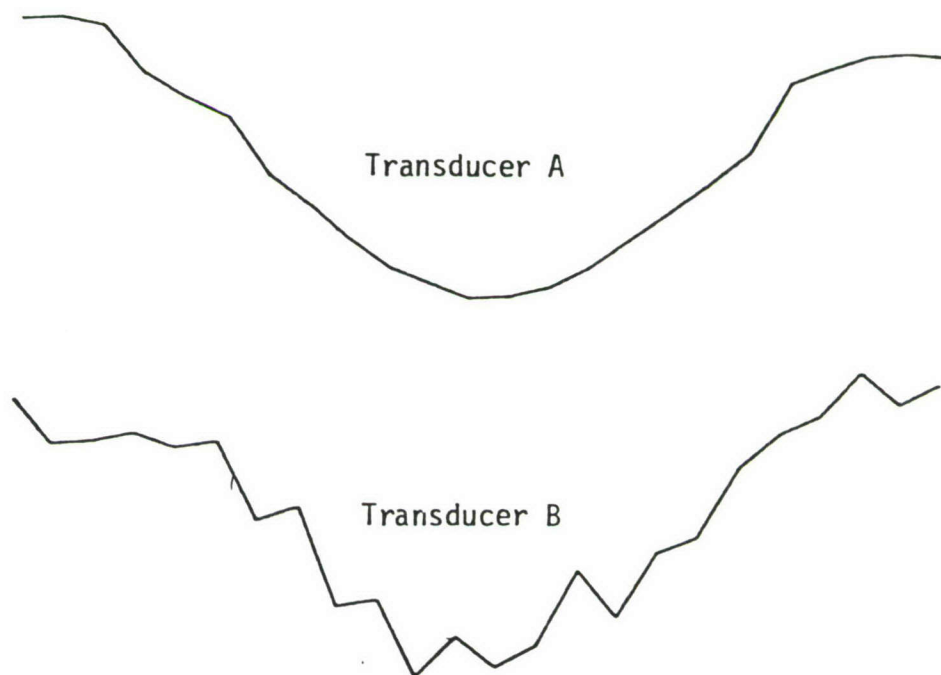


Figure 46 Second Bending Mode in Slope Coordinates of
Free-Free Beam Measured with Accelerometers
of Different Noise Performance
Beam Length = 610 cm.
Section Radius of Gyration = 1.47 cm.
Mode Frequency = 18.7 Hz.

The former implementation is computationally faster since only ensemble averages are corrected rather than the discrete Fourier transform of each frame of data. However, the latter is far less sensitive to digital round-off error. It has been found to be superior for use with the 10 or 12 bit A/D converters and 16 bit word length generally found in minicomputer-based laboratory data systems.

The measurement of H_c is of course quite critical in the method. It is accomplished by using a shaker to input exactly the same broadband random acceleration to both transducers. H_c is then computed simply as $S_{e_1 e_2} / S_{e_1}$ for this excitation. It should be noted that the correction function measured in this way will include mismatch between the low-pass anti-aliasing filters which must be inserted prior to digitizing. Their mismatch will generally be larger than the combined mismatch of all other elements in the channels. Thus, it is not correct to use H_c measured by this method to compute $\delta(f)$ in Eq. (6.27) if analog differencing prior to digitization is to be used. It is correct if differencing is to be done after digitization, although this would clearly be an inferior method.

A second experiment was performed to illustrate the effects of interchannel mismatch and to test the correction method described above.

A short section of aluminum I-beam was mounted as shown in Figure 47. Dimensions and mounting were chosen so as to approximate pinned-free end conditions. It was driven in the stiff direction by a shaker attached at the point calculated to be a node of the first flexural mode. Frequency responses between acceleration outputs at various points along the beam and force input were measured over 0 - 1.5k Hz. It was established that by careful adjustment of the mounting, it was possible to obtain a frequency response which was flat within a few percent up to 700 Hz. The amplitude of the function was also found to be directly proportional to the distance from the

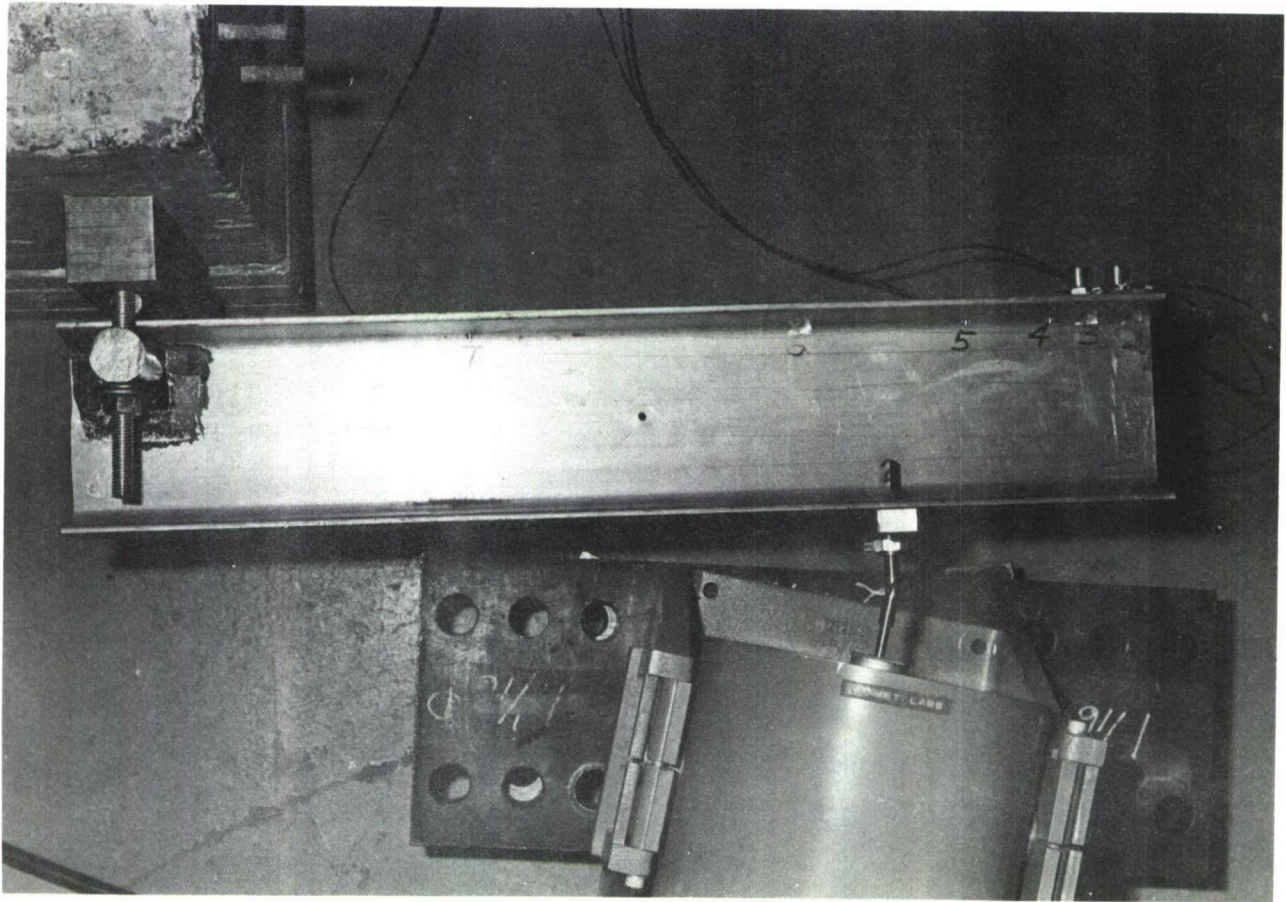


Figure 47 Test Structure for Evaluating Angular Acceleration Measurement Methods. Short I-Beam is Mounted in a Fixed-Free Condition and Driven by a Small Shaker.

pivot point to the acceleration sensor. Thus, it could be concluded that over this frequency range the beam behaves as a rigid body with a fixed axis of rotation. This implied that angular acceleration and its PSD could be accurately measured without differencing using only a single translational sensor. This measurement could then be used to check measurements made by differencing. In addition, dimensionless difference signal $1 - \alpha$ for any separation distance could be computed from Eq. (6.57) by taking ψ_2 and ψ_1 as the distances from the pivot point to the sensors. The experiment was set up and a time-random drive signal with a highly repeatable spectrum was established. Angular acceleration and its PSD were then measured by four methods.

1. Acceleration at a single point was measured and scaled by $1/x_p$ where x_p was the distance from pivot to sensor.

2. Accelerations at two points along the beam were transduced, analog differenced, and the result scaled by $1/\Delta x$ where Δx was the sensor separation distance. Channel gains upstream from the differencing amplifier were set according to the transducer manufacturers calibration data.

3. Procedure was identical to 2 except that the gain of one signal channel (transducer plus charge amplifier) was adjusted to exactly match the other at one frequency (300 Hz.). This required a gain adjustment of 0.4% relative to that used in 2.

4. A correction function H_c was measured for the two sensing channels at their nominal gains. Data from both channels were digitized and the PSD of their difference, corrected for mismatch error, was computed per Eq. (6.69). It was then scaled by $1/(\Delta x)^2$.

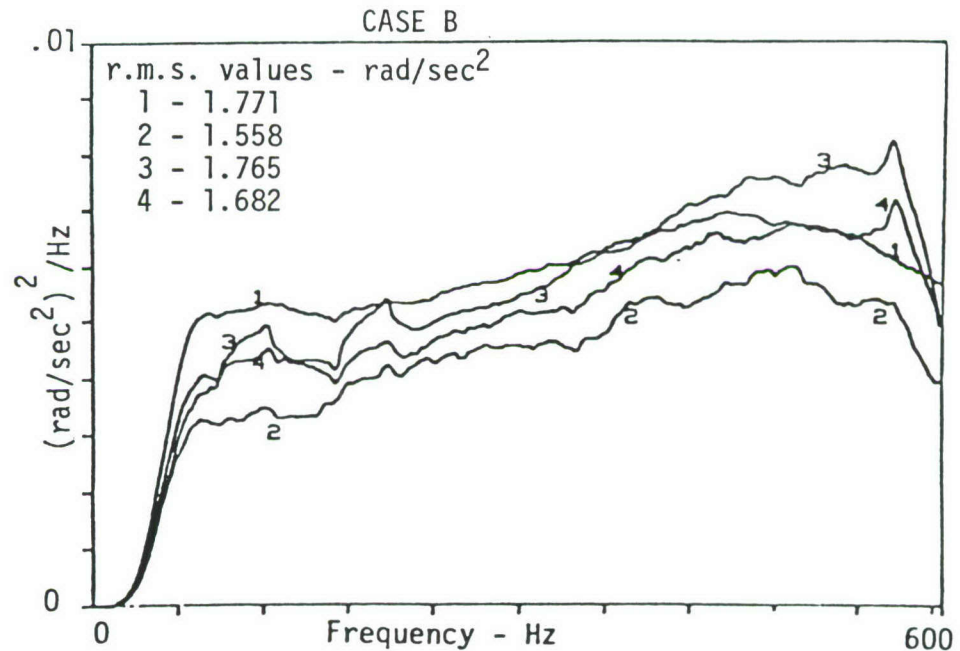
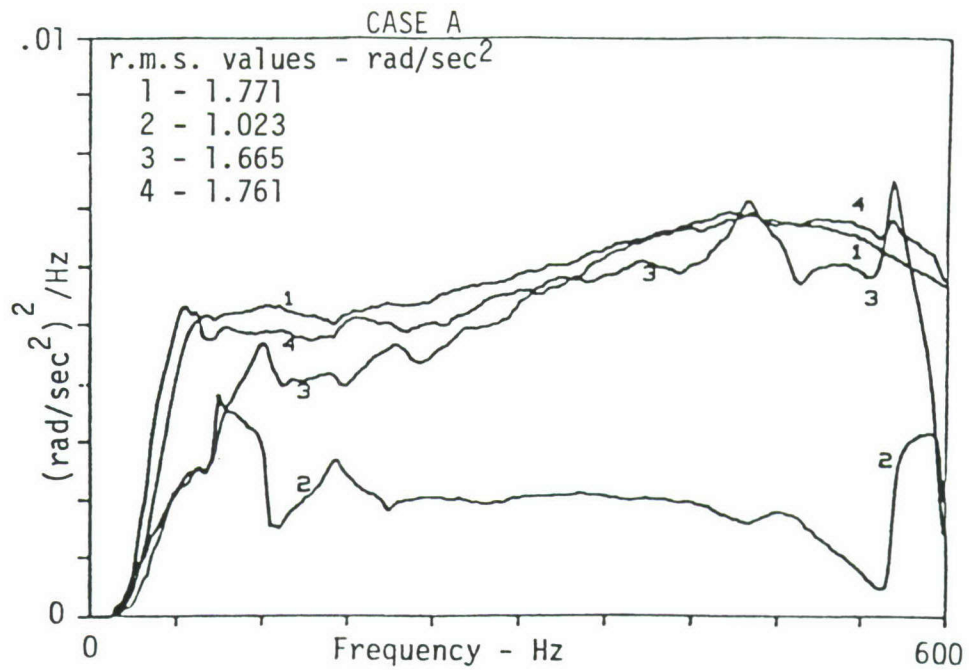
Two different separation distances were used for methods 2, 3, and 4. Data from method 1 is considered to be the most accurate and may be used to evaluate the other methods. Some details of the setup are given in Table 4.

TABLE 4
TEST CONDITIONS FOR SHORT BEAM EXPERIMENT
USED TO INVESTIGATE MISMATCHING ERROR

Case	ψ_1 cm.(in.)	ψ_2 cm.(in.)	$1 - \alpha$
A	58.57(23.061)	60.96(24.000)	0.000796
B	50.55(19.900)	60.96(24.000)	0.0174

Case A is a fairly demanding measurement situation because of the combination of moderately high frequency and small dimensionless difference signal. Case B is probably more typical of day-to-day applications in optical systems. Results are shown in Figure 48 for Case A ($1-\alpha = 0.0008$) and Case B ($1-\alpha = 0.0174$). For Case B, analog differencing with no special attention to balancing is satisfactory. For Case A, a drastic improvement is realized by more careful analog balancing but some error is still evident. Digital frequency compensation shows the best performance in this difficult case. All PSD's were measured by ensemble averaging of 300 frames with 2.25 Hz. frequency resolution. These were then smoothed by averaging in frequency over 16 spectral lines. The resultant estimate is thus based on 9600 statistical degrees of freedom which provides the smooth curves desirable for comparison.

In summary, it appears that the extra setup effort and data processing required by the frequency compensation method are justified only in severe cases. Such cases may be identified as those involving small transducer separations (a few cm. or less), high frequencies (200 Hz. or more), or high common mode accelerations. It may also be worthwhile if sensors of different types or low resonant frequencies (less than ten times the desired data bandwidth) are to be paired. Attention to accurate gain matching, however, is always worthwhile if analog differencing is used.



Key

- 1 = scaled from single point acceleration
- 2 = analog differenced with nominal channel gains
- 3 = analog differenced with gains matched at 300 Hz.
- 4 = digital differenced in frequency domain after correction

Figure 48 Measured PSD of Angular Acceleration of a Rigid Body with One Axis Fixed

6.3 CONCLUSION

Theoretical expressions have been developed for use in evaluating or predicting the validity of two types of angular acceleration spectral data obtained by differencing of linear transducer outputs. These are response power spectra and single input-single output frequency responses. A non-dimensional difference signal quantity α and a coherence-like quantity Γ have been defined. They are indices which may be used to compute figures of merit in the process of measuring power spectra and frequency responses of differential quantities by digital Fourier transform methods.

Experiments have been carried out to demonstrate the undesirable effects of single channel noise and interchannel mismatch as well as to test a method of reducing errors due to the latter.

The work described in this chapter has led to three principal conclusions regarding angular vibration measurement:

- (1) Measurement by analog differencing of accelerometer signals is, in general, an excellent method for day-to-day development work provided frequencies are not below 30-50 Hz., transducers with good noise performance are used, and some attention is paid to analog gain matching.
- (2) For certain critical measurements such as high-frequency angular vibration of mirror surfaces, differencing of acceleration signals may be the only practical method.
- (3) Errors in differential measurement due to channel mismatching may be reduced by digital frequency domain processing although this is warranted only in difficult measurement situations.

SECTION VII

TEST AND EVALUATION OF PREDICTION METHOD

During Phase III of this contract, a test was performed to assess the accuracy and usability of prediction techniques for angular vibration which were developed during Phases I and II. Initially, it was expected that tests of both low and high frequency methods would be performed during Phase III, and that the procedure would involve forced vibration tests of a fairly complex airframe-like structure. The test plan was later reduced in scope to concentrate on evaluation of Semi-Loof/NASTRAN methods for low frequency prediction. This decision was made at the end of Phase II after consultation with AFFDL. It was motivated by several factors.

- (1) Since the Semi-Loof implementation had nearly reached the status of a deliverable product, its testing and evaluation were given priority over that of high frequency methods where the work was still relatively basic and exploratory.
- (2) The SEA work had not yet reached the stage where application to a realistic airframe structure would be meaningful. It was felt that, at a minimum, an SEA model should be able to account for connections between components at multiple degrees of freedom in order to be worth testing on a simulated airframe/optical system.
- (3) The time and cost of completing the Semi-Loof development, modeling a reasonable airframe section, and testing to evaluate the model were thought to be fairly predictable. Together with reporting requirements, they accounted for about 85% of the available Phase III man-hours. Thus, there was no real possibility of significant further SEA development, let alone its testing, during Phase III.

For these reasons, the Phase III test was confined to measurement and prediction of low frequency response of a section of an aircraft fuselage under known excitation.

Section 7.1 describes the test object and its preparation. Section 7.2 describes the Semi-Loof finite element model and demonstrates the use of the pre- and post-processors with NASTRAN. In Section 7.3, the experimental procedure and equipment are described, and Section 7.4 presents the comparison of predicted and measured responses.

7.1 DESCRIPTION OF THE TEST OBJECT

The structure chosen as a test case for the Semi-Loof element was a rear fuselage section of a 1950 vintage Marine fighter aircraft. It was located at the Alameda Naval Property Disposal Yard. With the cooperation of AFFDL/FBG, it was acquired and transported to Anamet's laboratory in San Carlos, California. It is shown, as received, in Figures 49 and 50, and after preparation in Figures 51 - 53. The preparation consisted of removing all wiring and hydraulics, general cleaning, and removal of several feet from the front, as well as the vertical stabilizer, in order to reduce the specimen to manageable size. The prepared structure weighed 334 pounds.

The structure is made up primarily of flat and singly curved panels with extensive stiffeners. There are a number of major ring stiffeners which probably carried concentrated loads, as well as numerous short minor stiffeners attached to individual panels. A box section built into the floor appears to have been an electronics bay. The top of the box also served as the rear engine mount.

While the test structure is by definition aircraft-like in construction, it may not be entirely typical. The fuselage section contains heavy forgings which carried concentrated loads from the vertical stabilizer, the carrier arresting hook, and the jet engine. However, these forgings are tied together by shell structure, which should make the overall object a good test case for Semi-Loof modeling. In addition, the nearly-rigid forgings make excellent response points for measuring angular vibration.

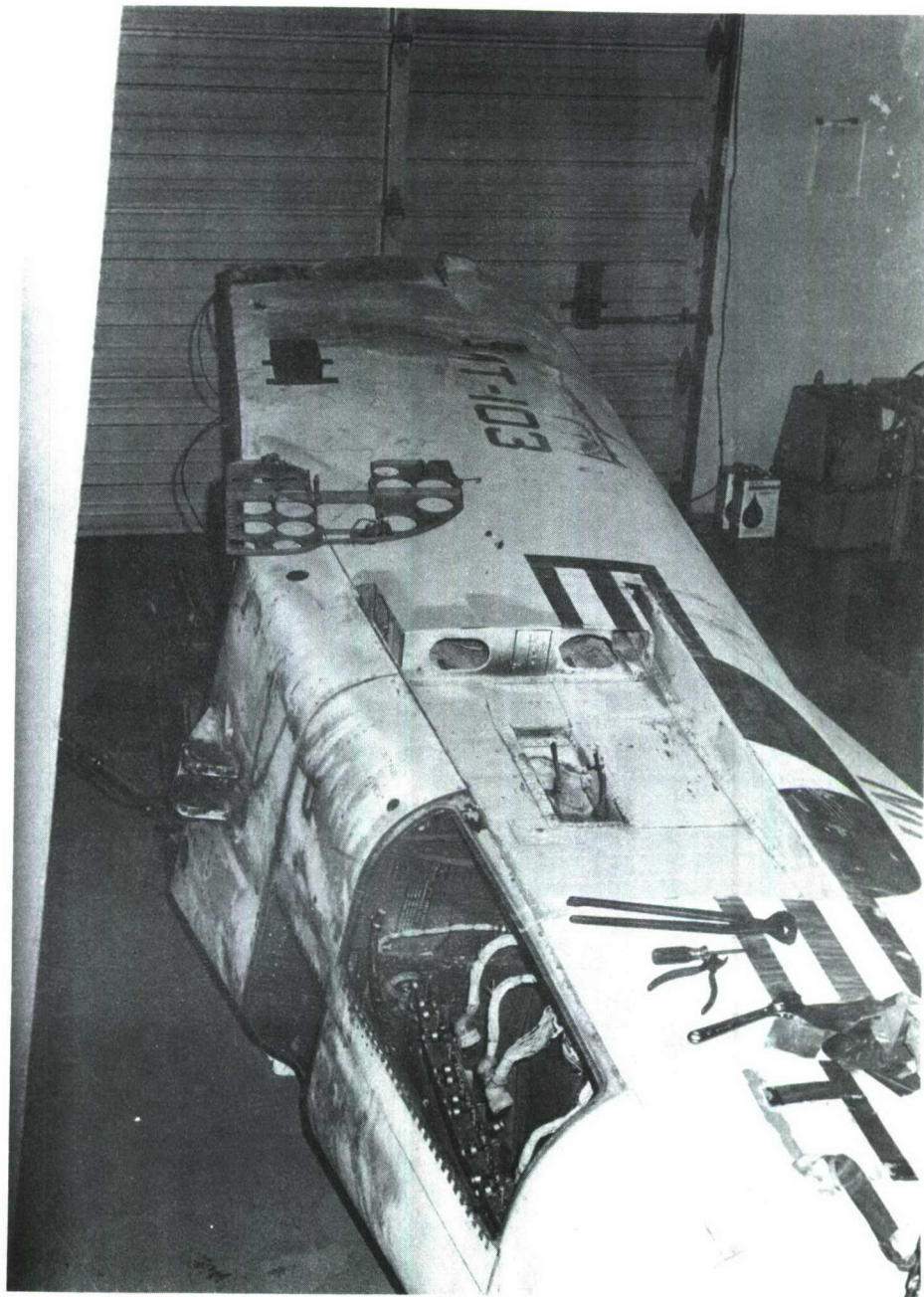


Figure 49 Exterior View of Test Fuselage in As-received Condition

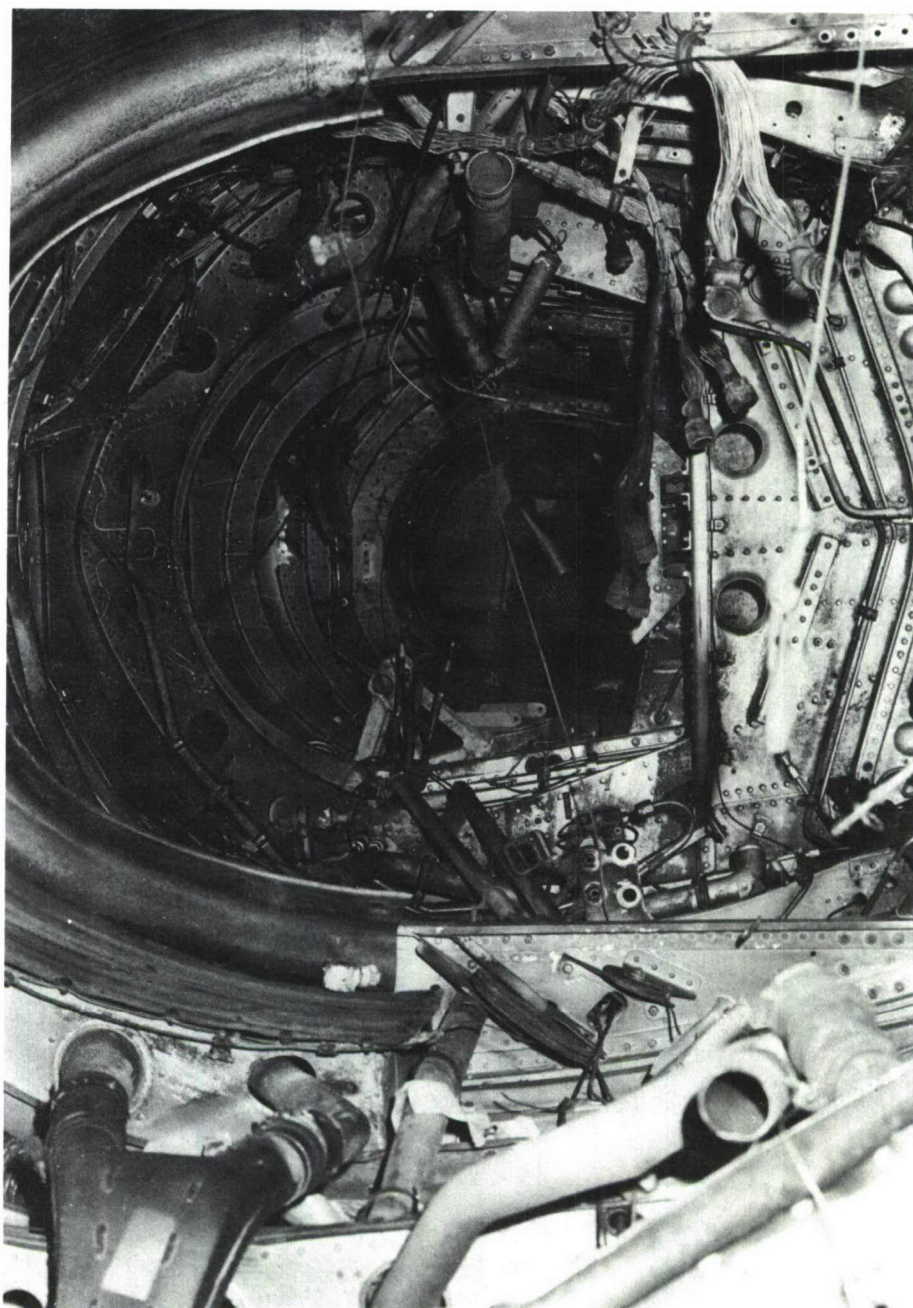


Figure 50 Interior View of Test Fuselage in As-received Condition

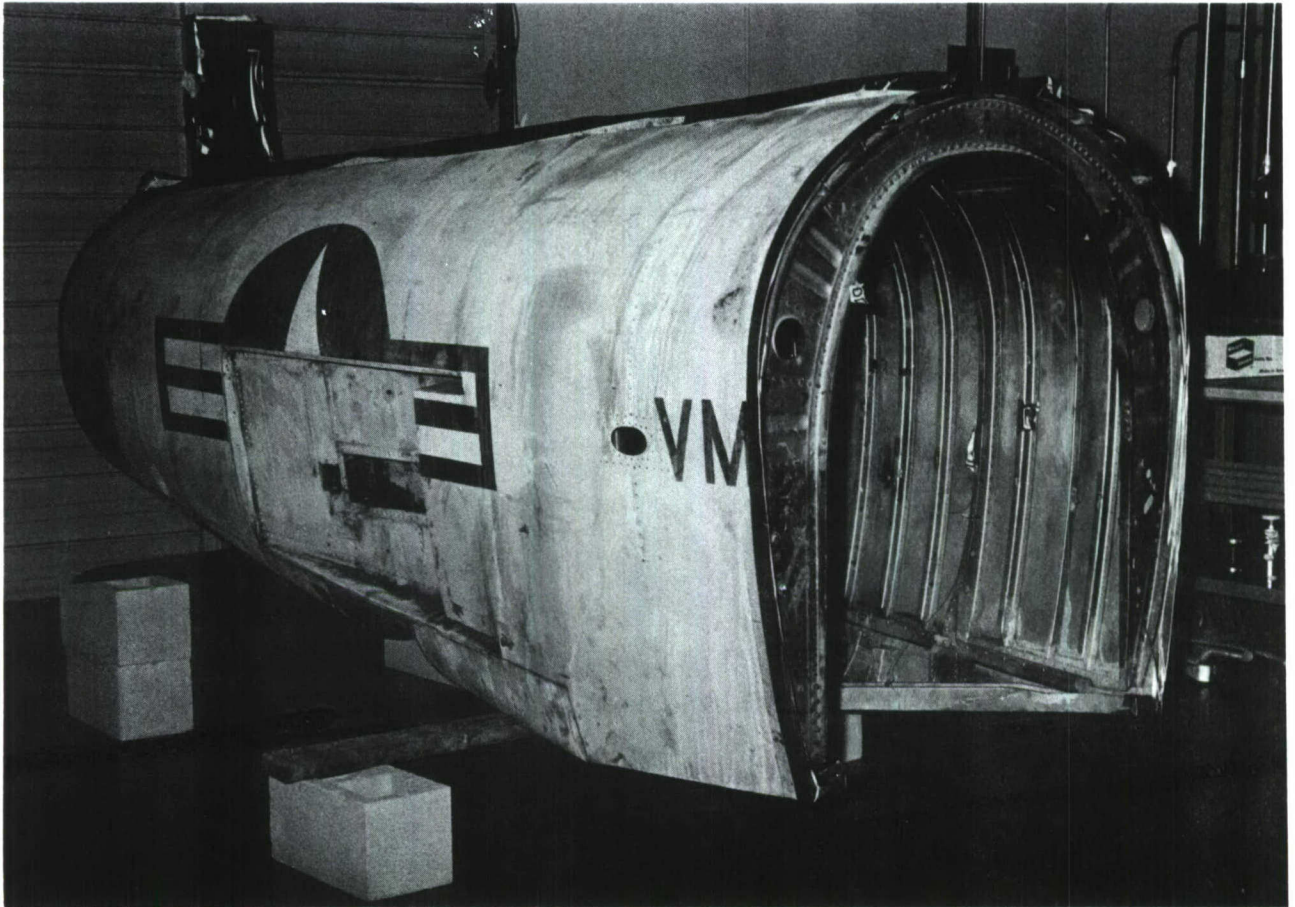


Figure 51 External View of Prepared Test Fuselage

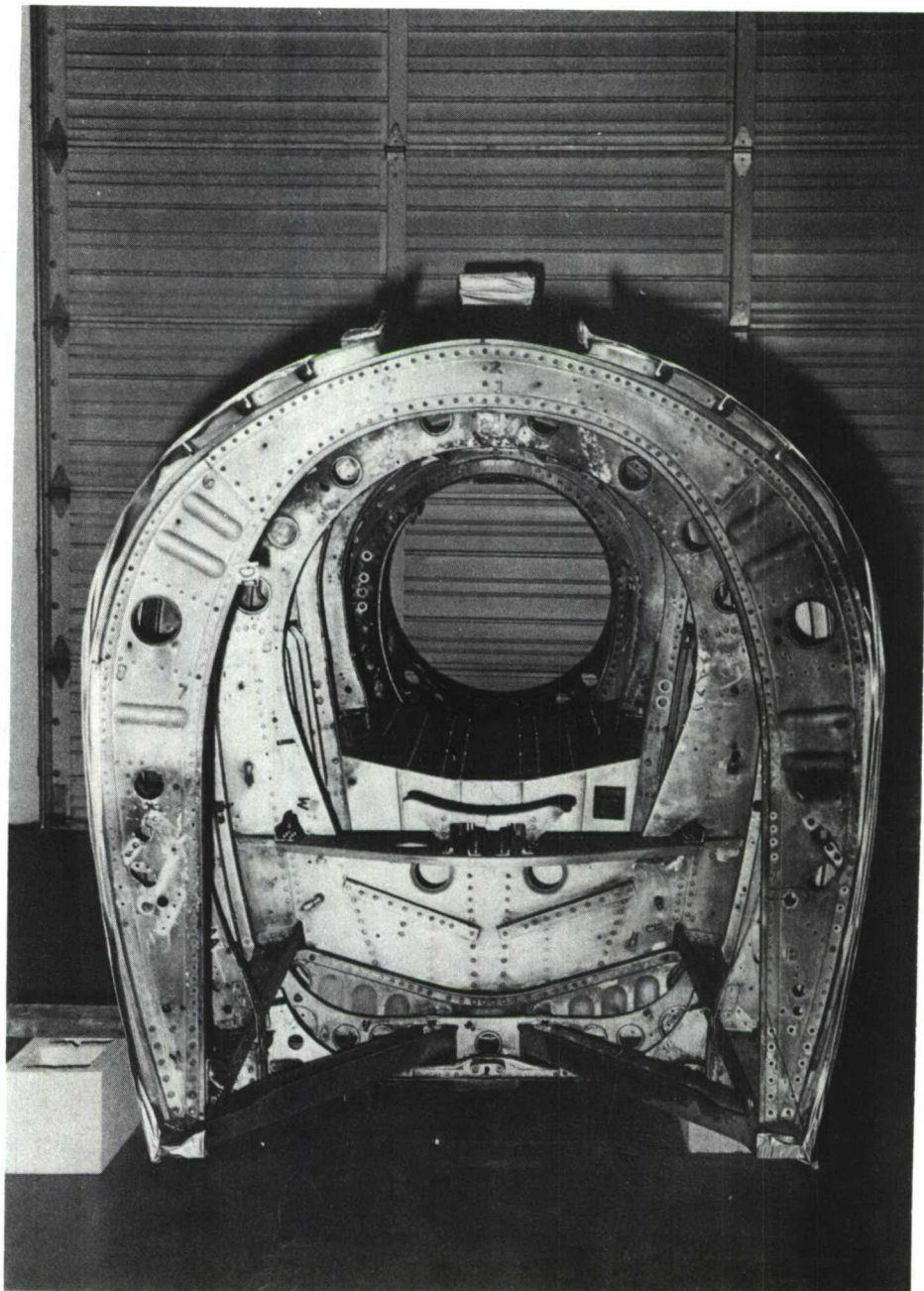


Figure 52 Internal View of Prepared Test Fuselage

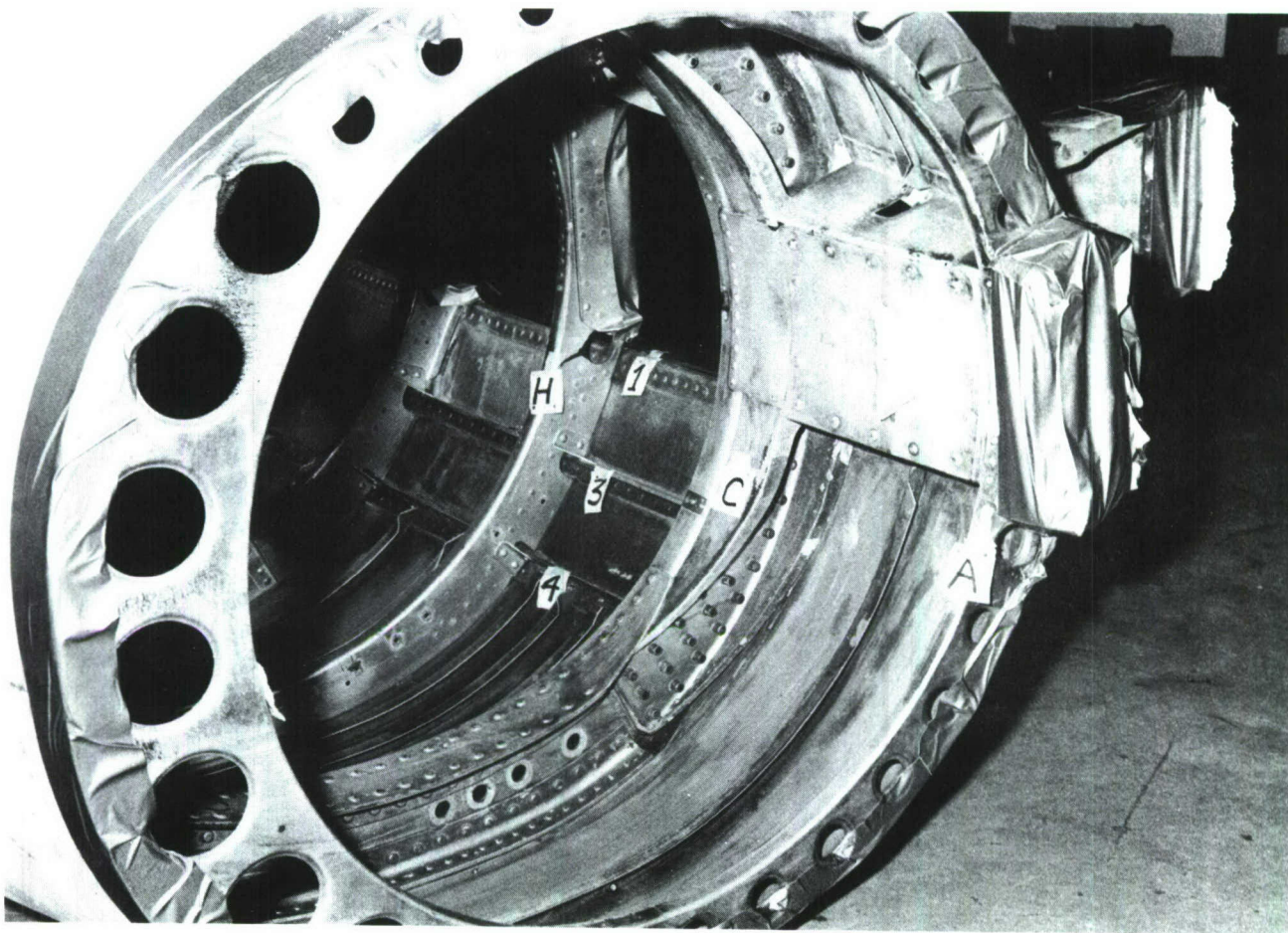


Figure 53 Rear View of Prepared Test Fuselage Showing Construction Details

7.2 FINITE ELEMENT MODEL

An outline of the Semi-Loof model of the fuselage is shown in Figure 54. The lines shown are the boundaries of the shell elements used. Beam elements are not shown. Symmetry was assumed about a vertical plane running along the axis of the fuselage, in spite of some minor differences between one side and the other. While the structure is geometrically symmetric, or nearly so, individual mode shapes are either symmetric or anti-symmetric. Two options were available for handling the symmetry plane. One was to compute two distinct sets of vibration nodes: one with symmetric boundary conditions [$f(y)=f(-y)$], and one with anti-symmetric boundary conditions [$f(y)=-f(-y)$]. These sets would then have to be merged prior to performing frequency response analysis. Although this approach was feasible, it was discarded because of the uncertain logistics involved in merging the two sets of nodes. Instead, the other side of the fuselage was generated by reflection about the symmetry plane, using a small Fortran program. A complete listing of the input data can be found in Appendix C.

Major rib stiffeners were modeled by offset curved L00F3 beam elements. The skin was modeled by a single row of L00F8 (and some L00F6) curved shell elements between major ribs. Most of the skin was stiffened by minor ribs and/or stringers, mostly with a "Z" cross-section. These minor stiffeners were "smeared" into the shell structure using the procedure outlined in Appendix A for calculating equivalent homogeneous orthotropic shell properties based on stiffener section properties, stiffener spacing, and skin properties. In fact, the PL00FX data card described in Appendix A was developed specifically for this application. In all, the model contains 585 grid points, 142 L00F8 quadrilaterals, 57 L00F6 triangles, and 194 L00F3 beam elements.

Additional L00F8 elements were used for the floor, mentioned previously, and miscellaneous beams and rigid masses were used as needed. Material density was increased 5% to

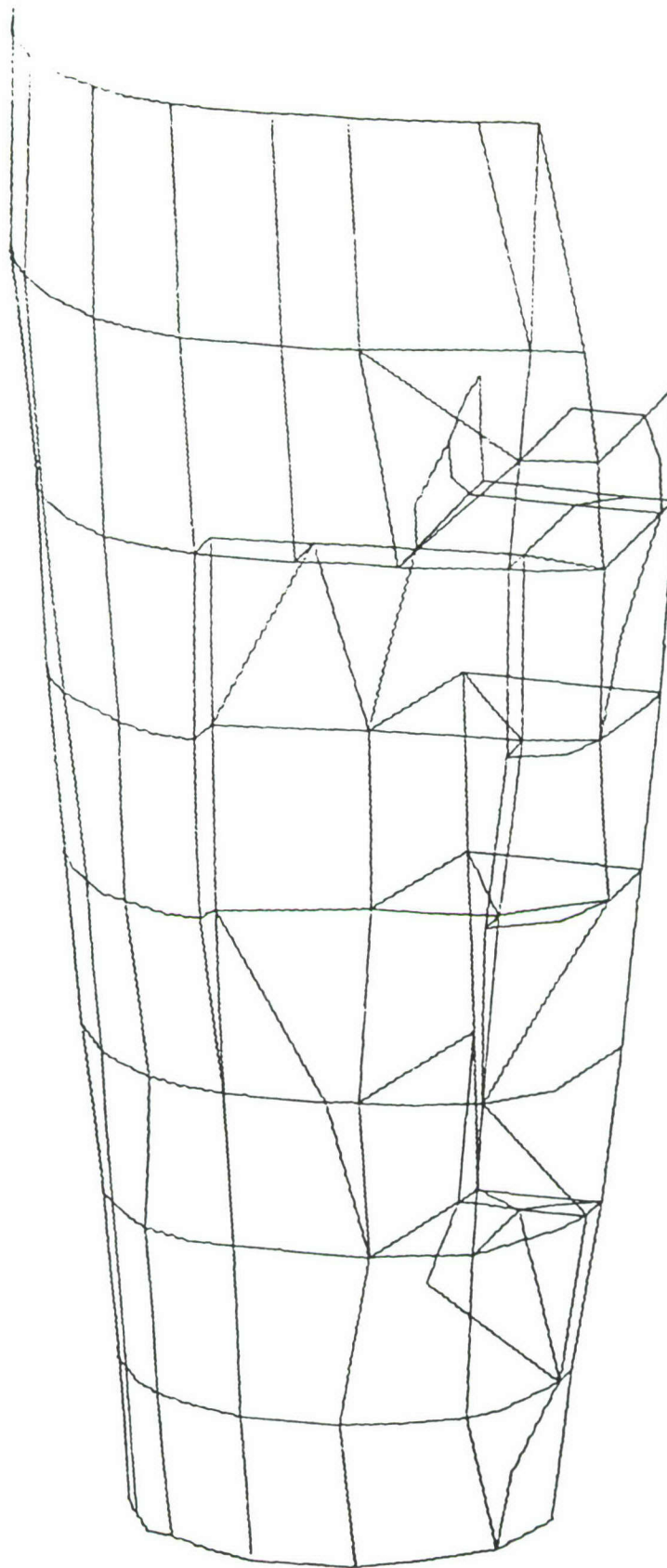


Figure 54 Semi-Loof Fuselage Model

account for rivets. With this adjustment, the total weight predicted by the model was 321 lbs. versus 334 lbs. measured.

The BANDIT node resequencing program was modified to recognize Semi-Loof elements and was used to generate SEQGP cards, which reorder the grid points for minimum bandwidth and thus reduced solution times.

Following usual analysis procedures, a static run with uniform external pressure was made to check out the connections and support conditions of the model. Following this a free vibration run was performed to extract natural frequencies and mode shapes. Finally, several short runs were made to produce frequency response and PSD curves. These latter runs used the NASTRAN Checkpoint/Restart feature to use information saved from the free vibration run. In each case, responses were computed over a range of 40 to 160 Hz. The 0-40 band was eliminated because the "free-free" modes actually had frequencies up to about 5 Hz. due to finite stiffness of the supports, and no attempt was made to model these supports with finite elements. Approximate modal damping values were used, based on experimental values. Damping was quite light and thus rather insignificant except perhaps for the heights of peak responses.

7.3 TEST DESCRIPTION

A series of forced vibration tests were carried out on the aircraft fuselage section described in Section 7.1. The tests were designed to furnish data for evaluating the Semi-Loof finite element model described in Section 7.2 and thereby, the Semi-Loof element itself. The experimental procedure and equipment are described in this section. Experimental results and comparisons with predictions are in Section 7.4.

7.3.1 Overview

Most linear dynamic response simulation under NASTRAN is done by normal mode methods. Thus, in developing a new element, one would like to make measured-vs.-predicted comparisons in terms of properties of individual modes. However, in the current case, this approach was not practical. In principle, it should be sufficient to compare predicted and measured values of natural frequencies and selected entries of the mode shape vectors (after normalizing to equal modal mass) which correspond to critical displacements and rotations. This implies that only a cursory modal survey covering the selected response points should be necessary. However, for a complex structure with numerous normal modes such as the fuselage, the task is not this easy. Unless natural frequency predictions are known in advance to be quite accurate, one is never sure which predicted normal mode should be compared with which measured mode. The usual solution is to measure mode vector entries for enough response locations to obtain a reliable physical picture of the overall mode shape. For example, one might identify a particular mode as the "first twisting" or "first bending" mode. Unfortunately, the time and cost allotted to testing of the fuselage under this contract were not adequate to allow a modal survey of this order, even with modern computer-based methods.

The approach actually taken was to perform a test which would measure the accuracy of the finite element model in an

end-to-end sense. What was given up in order to meet time and cost limits was detailed information about individual modes which one might use to improve the model.

The test consisted of mounting the fuselage on a soft suspension to simulate a free-free condition and measuring a number of acceleration/force frequency responses. Both linear and angular accelerations were taken as response variables. Input was a single point force in all cases. The frequency responses were measured in digital form and stored on disc for later processing and comparison with NASTRAN predictions. The functions were converted digitally to displacement/force frequency responses and these were used to predict displacement PSD's and r.m.s. values for a specified force input PSD. This computational procedure facilitated comparison with NASTRAN predictions in terms of output PSD and r.m.s. quantities. An important practical advantage is that NASTRAN and experimental output quantities could be compared on the basis of identical input force PSD without requiring precise closed loop control of the experimental input.

7.3.2 Test Procedure

Diagrams of the test object showing input and response point locations are presented as Figures 55 through 57. Table 5 gives the location, direction, and type (translation or rotation) of all degrees of freedom. Table 6 shows which of the possible entries of the frequency response matrix were actually measured. Each column represents a different input degree of freedom (shaker location) and each row represents a different response quantity. A total of 3 shaker locations were used in an attempt to acquire at least one frequency response function with a significant contribution from each mode in the 0-160 Hz measurement band. A total of 16 response degrees of freedom were used, of which 8 represented translations and 8 represented rotations. From Table 6 it may be observed that

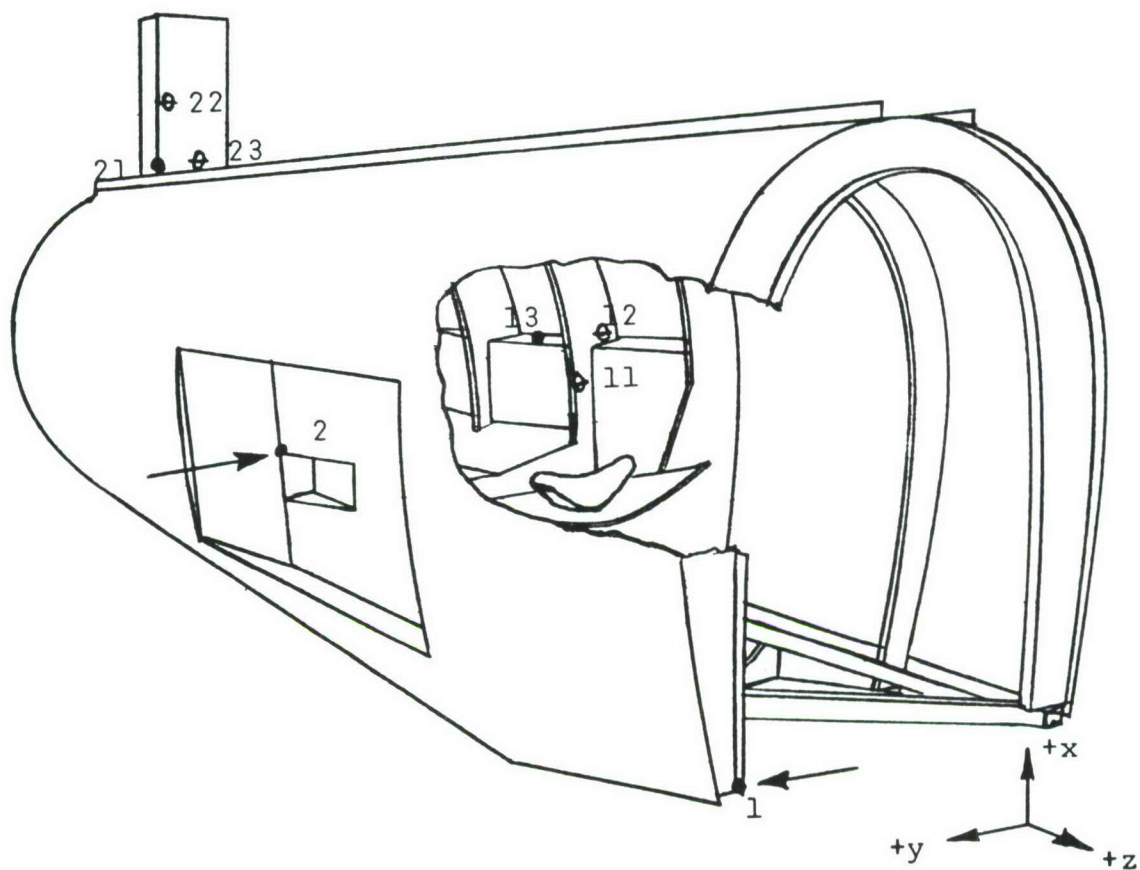


Figure 55 Sketch of Test Object Showing Input Force Location and Direction (Arrow), Translational Response Sensor •, and Rotational Response Sensor θ

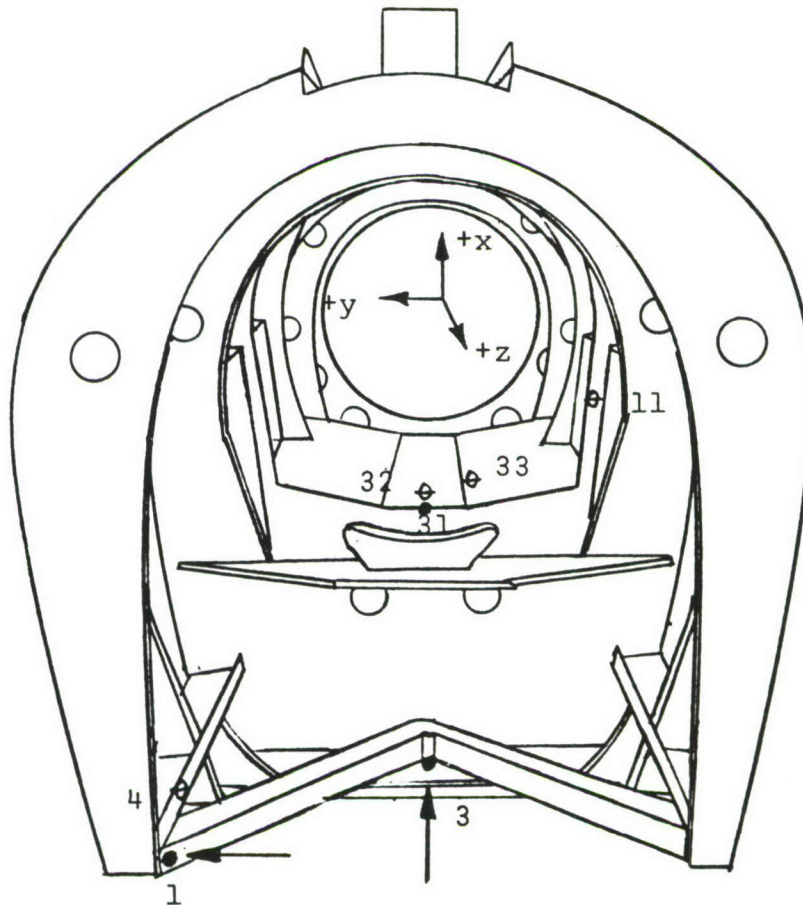


Figure 56 Front View Sketch of Test Object Showing Input Force Location and Direction (Arrow), Translational Response Sensor •, and Rotational Response Sensor θ

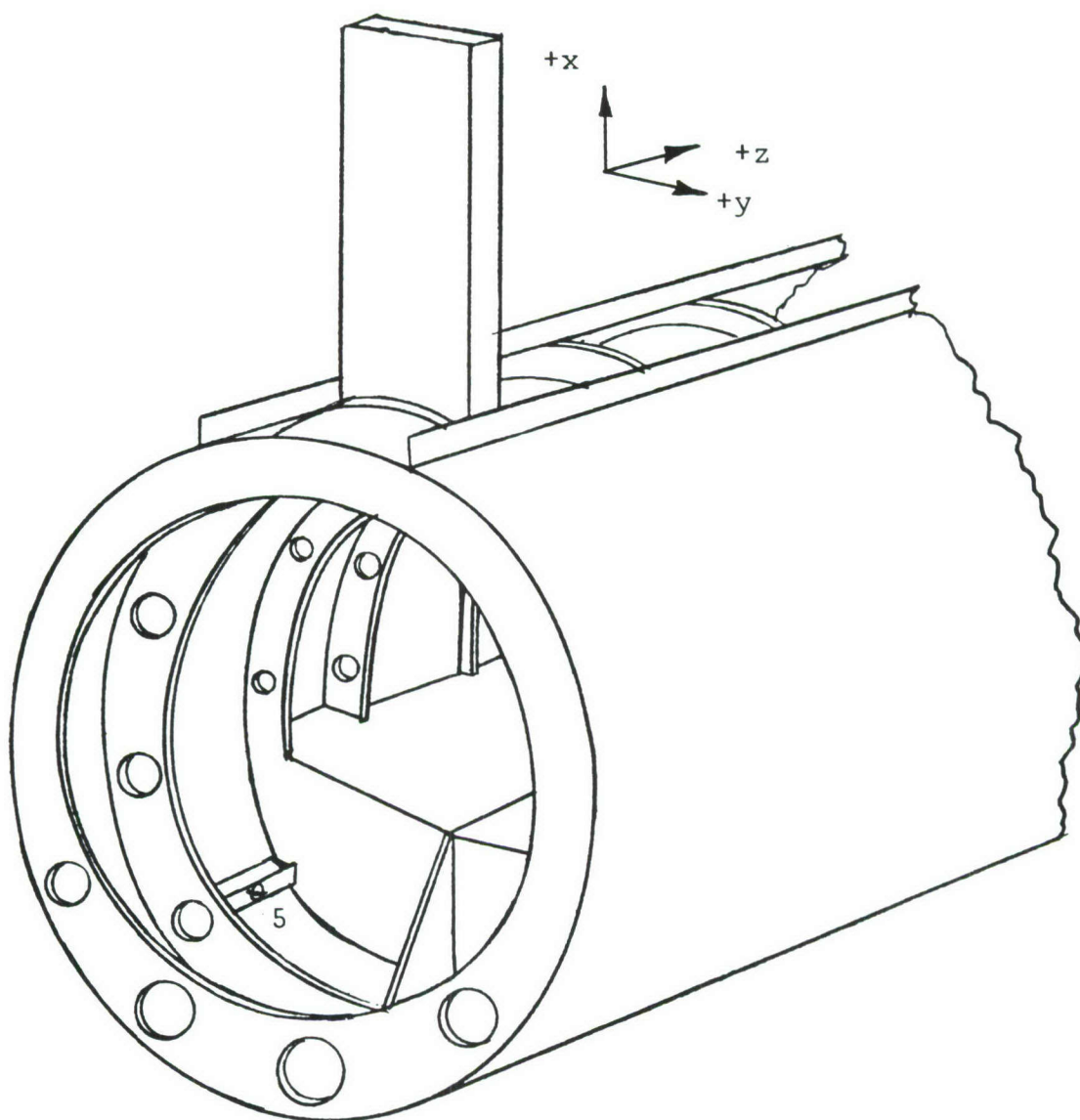


Figure 57 Partial Sketch of Test Object Showing Location of Rotational Sensor No. 5

TABLE 5
DESCRIPTION OF RESPONSE VARIABLES

<u>Number & Direction</u>	<u>Type</u>	<u>Location</u>
1y	trans.	Inside fuselage at extreme lower right front on end of heavy wishbone beam.
2y	trans.	Outside of fuselage, right side, near box section reinforcement for airbrake pivot bolt.
3x	trans.	Outside bottom of fuselage, near arresting hook bracket.
4x	rot.	Yaw at extreme lower right front on heavy beam outside wishbone.
5x	rot.	Yaw on major longitudinal stiffener at lower left, just behind electronics enclosure.
11z	rot.	Roll on major rib stiffener S, left inside of fuselage over electronics enclosure.
12y	rot.	Pitch at junction between major rib stiffener S and top of left airbrake box.
13x	trans.	Top of left airbrake box inside fuselage between ribs O and S.
21x	trans.	Outward normal direction (slightly to the right of vertical) on rib C near lower right corner of stabilizer stub.
21y	trans.	Lower right side of stabilizer stub.
21z	trans.	Lower right front of stabilizer stub.
22y	rot.	Pitch at base of stabilizer stub.
23x	rot.	Yaw at base of stabilizer stub.
31x	trans.	Top center front of electronics enclosure.
32z	rot.	Roll at top front center of electronics enclosure.
33y	rot.	Pitch at top front of electronics enclosure on stringer slightly left of center.

TABLE 6
MEASURED FREQUENCY RESPONSES

<u>Response</u>	<u>Input</u>		
	<u>1y</u>	<u>2y</u>	<u>3x</u>
1y	x		
2y		x	
3x			x
4x		x	
5x		x	
11z	x	x	x
12y	x	x	x
13x	x	x	x
21x	x	x	x
21y	x	x	x
21z	x	x	x
22y	x	x	x
23x	x	x	x
31x	x	x	x
32z	x	x	x
33y	x	x	x

a total of 38 different input/output pairs were used. Random excitation was provided by an electrodynamic shaker of 50 lbf (peak) capacity. Typical drive levels were 3-10 lbf (r.m.s.), depending on input and response locations. This drive level was found to give excellent coherence ($\gamma^2 > 0.97$) between force and acceleration over the frequency range of 50 to 160 Hz when the shaker drive signal was bandlimited to 5-180 Hz. However, a separate run with all drive power concentrated in the 5-60 Hz band was needed to get good coherence below 50 Hz.

Frequency responses for the 38 input/output pairs were measured over 0 to 160 Hz. The measurement was made using both a baseband FFT ($\Delta f = 0.312$ Hz) and a high resolution or zoom FFT ($\Delta f = 0.1$ Hz). The wideband measurements were intended as quick-look data, and the zoomed functions for estimation of modal parameters, such as damping. It was found that, in most cases, the wideband functions could be used to compute output PSD and r.m.s. quantities with negligible error.

A typical run was made by first setting the drive signal passband to 5-60 Hz and adjusting the force amplitude to about 10 lbf (r.m.s.). A baseband, 512 spectral line frequency response over 0-60 Hz was then measured and stored. The drive passband was next opened up to 5-180 Hz, and the amplitude adjusted as required. 400 seconds of the force and response signal, sampled at 640 Hz, were digitized and throughput to disc in real time. The shaker was then shut down and a stored program was run. It processed the stored data to obtain a 512 line baseband frequency response over 0-160 Hz, as well as five zoomed frequency responses of 200 spectral lines each, spaced over 60-80, 80-100, 100-120, 120-140, and 140-160 Hz. While this processing was taking place, the operator was free to move the sensors or shaker to the location desired for the next run.

A potential problem with the procedure as described above involves the rigid body modes of the test object. It is desirable that the actual test object and its NASTRAN idealization be identical with respect to support conditions. However, boundary conditions such as free or fixed which are easy to analyze may be quite difficult to implement. Likewise, a condition which is readily constructed may be impractical to model. The solution for this case was to use a soft, pneumatic suspension and model it as free-free. This implies that pure rigid body modes at zero frequency will be suppressed, but will reappear as nearly rigid modes at low frequency. Also, some slight modification of flexural modes will occur due to constraint forces. Some care is required if predicted-vs.-measured comparisons are made on the basis of displacement, which may be dominated by contributions from rigid body modes which were modeled only approximately. Two measures were taken to minimize this problem:

- (1) The hypothetical input force PSD used for measured-vs.-predicted comparison was taken to be zero below 40 Hz. Since the highest rigid body mode was at about 5 Hz. and the lowest elastic mode was at 45 Hz, this reduced the effect of rigid body modes considerably.
- (2) Frequency responses selected for comparison were those where, within the 40-160 Hz comparison band, contributions from elastic modes were large relative to those from rigid body modes.

It is possible to further reduce the effect of support conditions on the measured frequency response functions by careful data processing. A parametric representation of each function may be obtained by curve fitting a complex partial fraction expansion to the measured version. The analytic form represents the contribution of each mode, flexural or rigid, by a pair of terms. The function may then be resynthesized including only the terms corresponding to straining modes. While the data taken was sufficient to carry out this procedure, it was not done due to time and cost limitations.

7.3.3 Test Equipment

Figure 58 is a schematic of the test setup. A list of equipment used is supplied as Table 7.

Because of the low excitation force level available relative to the mass of the test object, it was found that very low-noise acceleration sensing was necessary. The combination of Endevco 2207-200 transducers and 2735 charge amplifiers was found to give an instrumentation noise level of about 28 μg (r.m.s.). High coherence measurements were possible for all input-output pairs as long as the available force was concentrated in the frequency band of interest and some effort was made to maintain a quiet test environment.

Angular acceleration was sensed by analog differencing of translational acceleration signals. Analog gain balancing was performed at 100 Hz prior to testing. Residual unbalance was estimated as less than 0.5% gain and 0.5° phase.

Typical r.m.s. response levels were on the order of 0.04 g. The lowest observed elastic mode of the fuselage was at 45 Hz. We may use Eqs. (6.27) and (6.37) of this report to estimate that at this frequency a differential narrowband signal/noise ratio of 30 dB could be maintained down to a dimensionless difference signal level of $1-\alpha = 0.05$ and 20 dB could be maintained down to $1-\alpha = 0.005$. This was considered adequate for the purpose at hand. Transducer separation varied for the different response locations, but was always in the 6.5 - 10.5 inch range. Flexural error was estimated using Figure 44 of this report and was found to be less than 10% for all cases. However, data from location 11z was eventually discarded because it was decided that the beam flange, on which the transducers were mounted, was too flexible to give a good approximation to a one-dimensional beam.

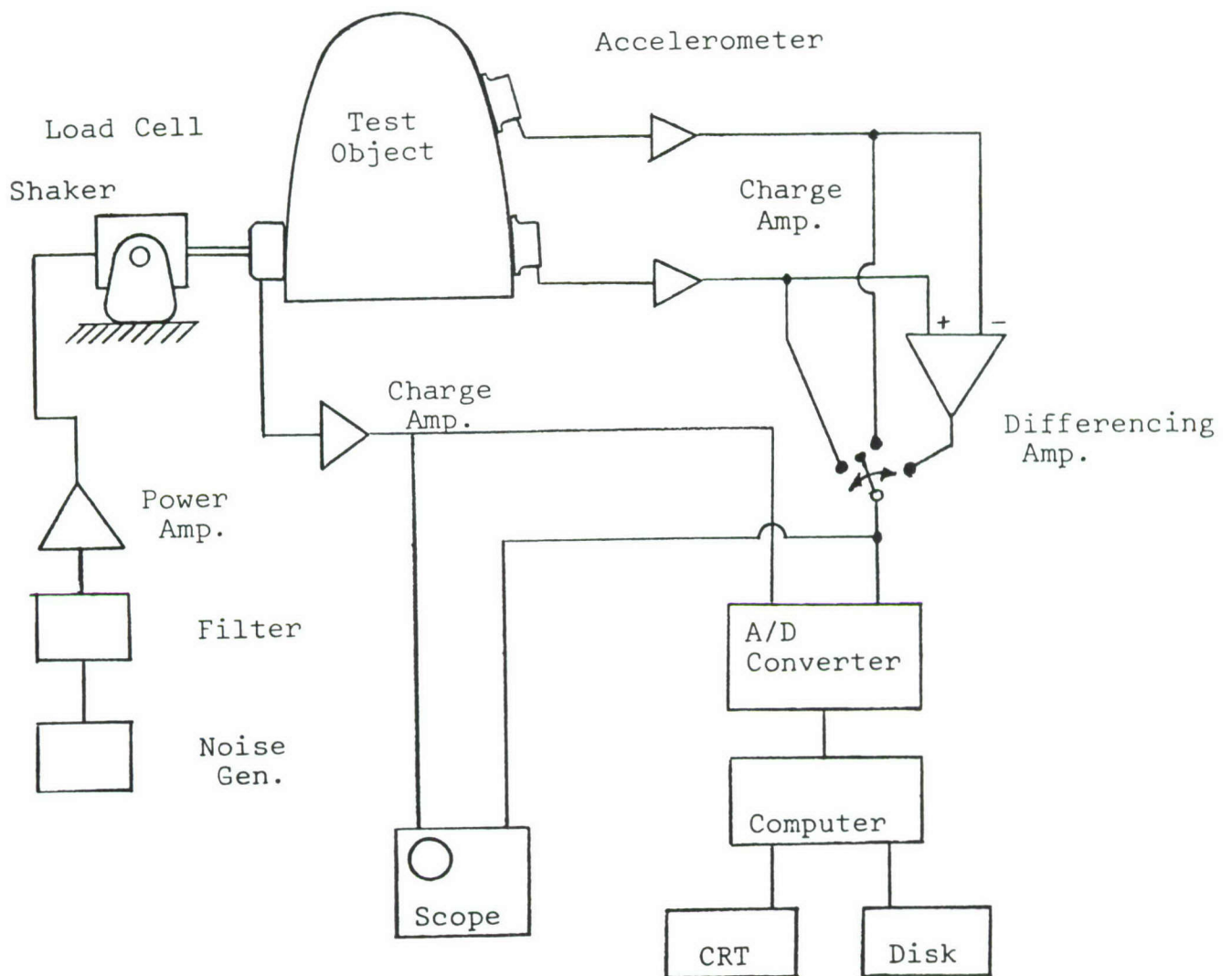


Figure 58 Schematic of Test Setup

TABLE 7
TEST EQUIPMENT

<u>Item</u>	<u>Manufacturer</u>	<u>Model</u>
Load Cell	Bruel & Kjaer	8200
Accelerometers	Endevco	2207-200
Charge Amp.	Endevco	2735
Dif. Amp.	Honeywell	122
Noise Generator	Wavetek	132
Variable Filter	Krohn-Hite	3323
Shaker	Unholtz-Dickie	Model 1 (50 lbf)
Power Amp.	Unholtz-Dickie	TA35I
A/D Converter	Time/Data	TDA-25
Computer	DEC	PDP 11/34
Scope	Tektronix	434

7.4 Results

Table 8 lists the natural frequencies measured in the test versus those predicted analytically. It should be emphasized that there is no way to know for certain which experimental mode corresponds to which analytical mode. For example, the eighth analytical mode (46.3 Hz) may correspond to either the seventh (45.5 Hz) or eighth (48.7 Hz) mode, and this could be determined only by a detailed investigation of the mode shape geometries. It should also be noted that the list of experimental modes is almost certainly incomplete due to the finite number of response points. The analytical model includes all modes up to a certain frequency, including any local modes that may have been missed in the test. "Generalized dynamic reduction" was the eigenvalue solution method that was chosen. If Guyan reduction had been used, with a judiciously selected analysis set, it probably would have been possible to eliminate these local modes.

PSD plots are shown in the following figures for three separate driving points as listed in Section 7.3. Figures 59 through 70 are experimental results and Figures 71 through 75 are analytical results. At best, there is some qualitative agreement between respective pairs of plots. Both exhibit peaks around 70 Hz. The analytical results show a peak around 90 Hz which is not seen in the experimental plots. Both also exhibit peaks in the area of 140 to 150 Hz. Table 9 lists r.m.s. values for several response PSD's. Correlation is fairly good for some of these values.

Power spectral density functions are plotted in linear amplitude-linear frequency format. This allows a direct interpretation of the area under the curve in terms of the mean square response. As a further aid to visualization of the frequency distribution of response, a quantity called the normalized cumulative r.m.s. response is plotted. It is a function of frequency defined as

TABLE 8
NATURAL FREQUENCIES OF THE FUSELAGE

<u>Experimental</u>	<u>Mode No.</u>	<u>Analytical</u>
45.5	7	22.7
48.7	8	46.3
70.1	9	71.9
79.7	10	74.2
83.0	11	79.3
109.8	12	85.1
111.8	13	88.4
113.9	14	88.9
121.9	15	90.9
126.6	16	95.5
130.0	17	97.1
133.3	18	112.0
137.0	19	112.2
140.0	20	116.4
148.7	21	118.0
155.9	22	120.7
	.	
	.	
	.	
	30	149.5

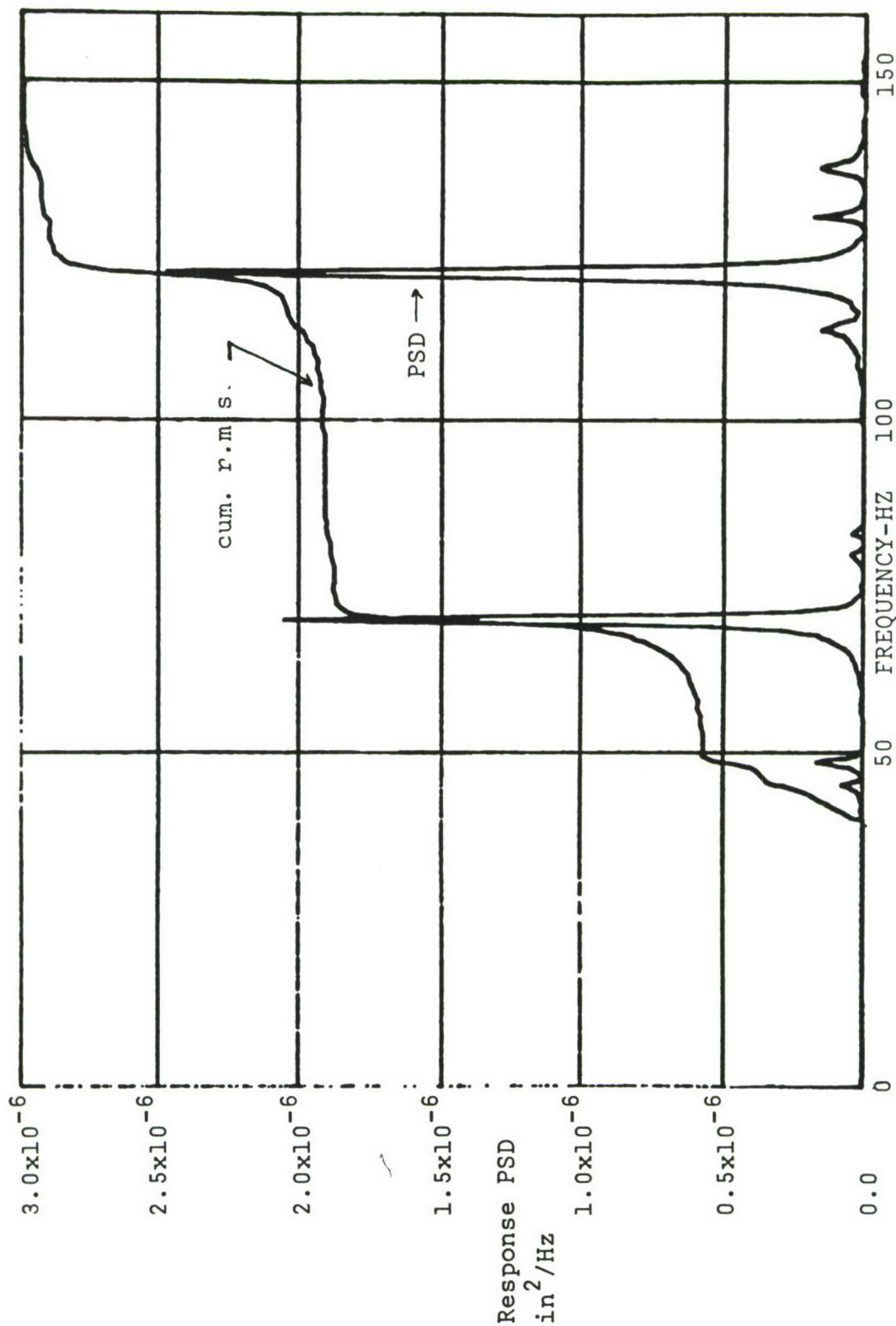
$$\text{cum. r.m.s.} = \sqrt{\frac{\int_0^f S(f) df}{\int_0^{f_{\max}} S(f) df}} \quad (7.1)$$

Since $S(f)$, the PSD, is always positive, the cumulative r.m.s. is also, and will always range from zero at $f=0$ to unity at $f=f_{\max}$. Thus, no scale is shown for it on the plots of Figures 59 to 70.

One would have to call this test inconclusive with respect to evaluation of the Semi-Loof elements. The performance of the elements is masked by a number of other considerations which dominated this exercise. These questions include the following: (1) What was the true damping in the system? (2) What were the true support conditions? (3) What was the effect of the angular accelerometers and their mounts, in terms of both stiffness and inertia (not considered in the model)? (4) What local structural deformations might have been missed by the finite element model.

TABLE 9
R.M.S. VALUES OF PSD RESPONSES OF THE FUSELAGE

<u>Drive Point</u>	<u>Response Point</u>	<u>Experimental</u>	<u>Analytical</u>
1Y+	22Y+	5.5×10^{-5}	7.6×10^{-6}
1Y+	23X+	7.3×10^{-5}	2.3×10^{-5}
2Y-	32Z+	2.1×10^{-5}	4.1×10^{-6}
3X+	22Y+	8.5×10^{-6}	7.6×10^{-6}
3X+	23X+	3.0×10^{-6}	2.4×10^{-6}



RESPONSE PSD AND NORMALIZED CUM. RMS FOR INPUT PSD -
1.00 LBFXX2/HZ OVER 40-160 HZ.

RESPONSE RMS = 3.27776E-03 IN.

FREQRESP-MODULUS
1Y+ 1Y-
020179-000000

Figure 59 Experimental Results
Response Point 1, Y Translation
Driving Point 1, Y Translation

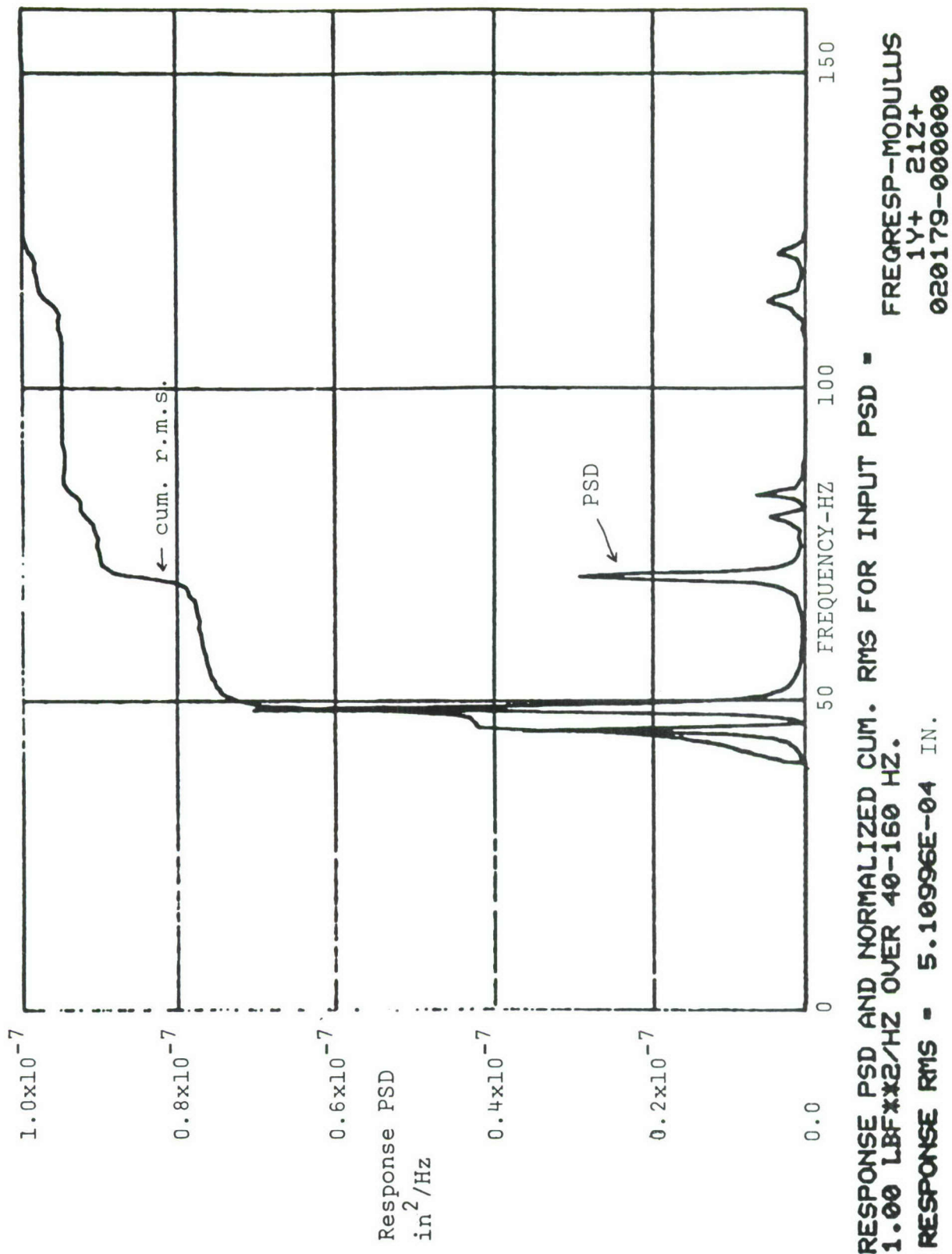
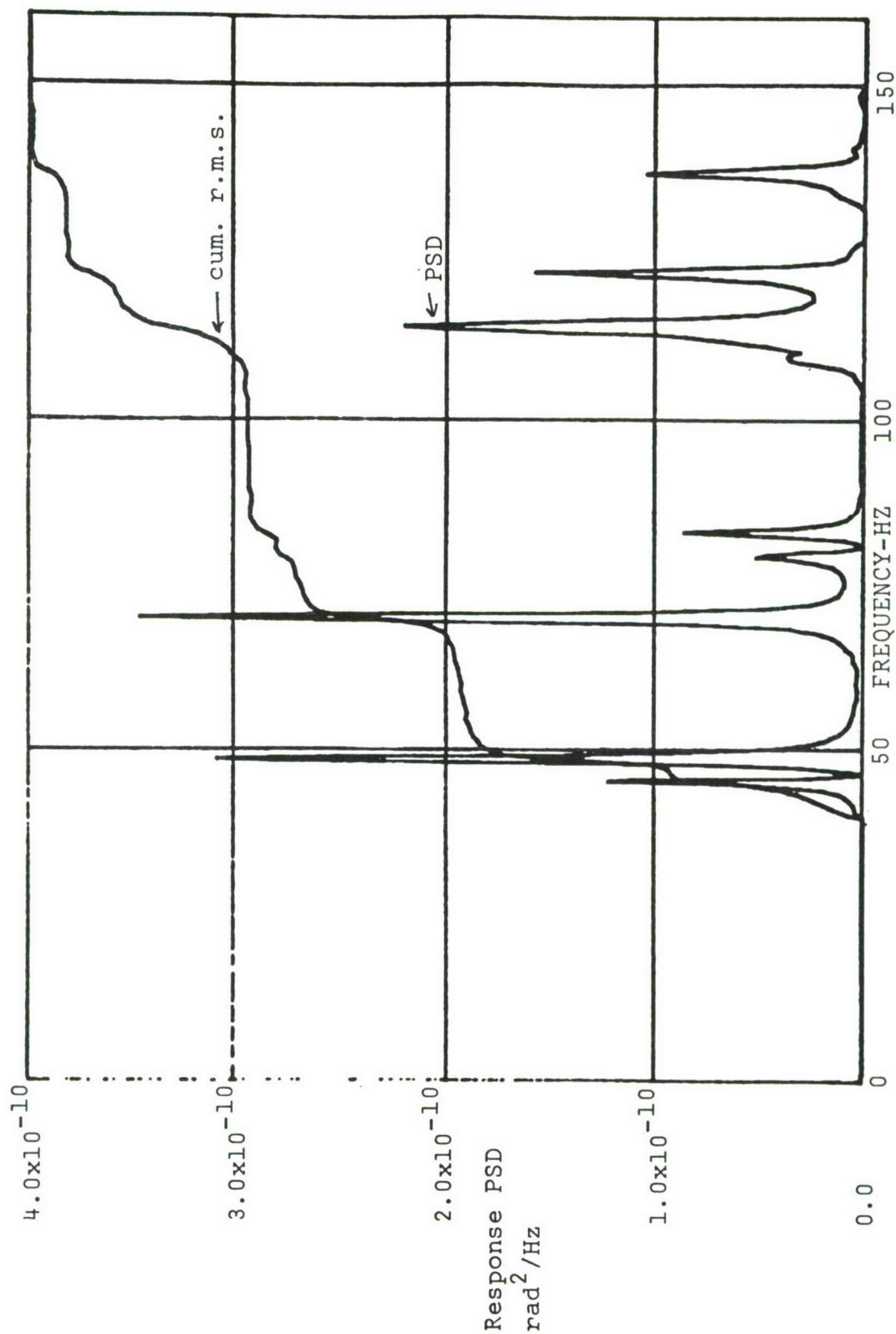


Figure 60 Experimental Results
Response Point 21, Z Translation
Driving Point 1, Y Translation



RESPONSE PSD AND NORMALIZED CUM. RMS FOR INPUT PSD -
 1.00 LBF*HZ OVER 40-160 HZ.

RESPONSE RMS = 5.46344E-05 RAD.

FREQRESP-MODULUS
 1Y+ 22Y+
 020179-000000

Figure 61 Experimental Results
 Response Point 22, Rotation about Y
 Driving Point 1, Y Translation

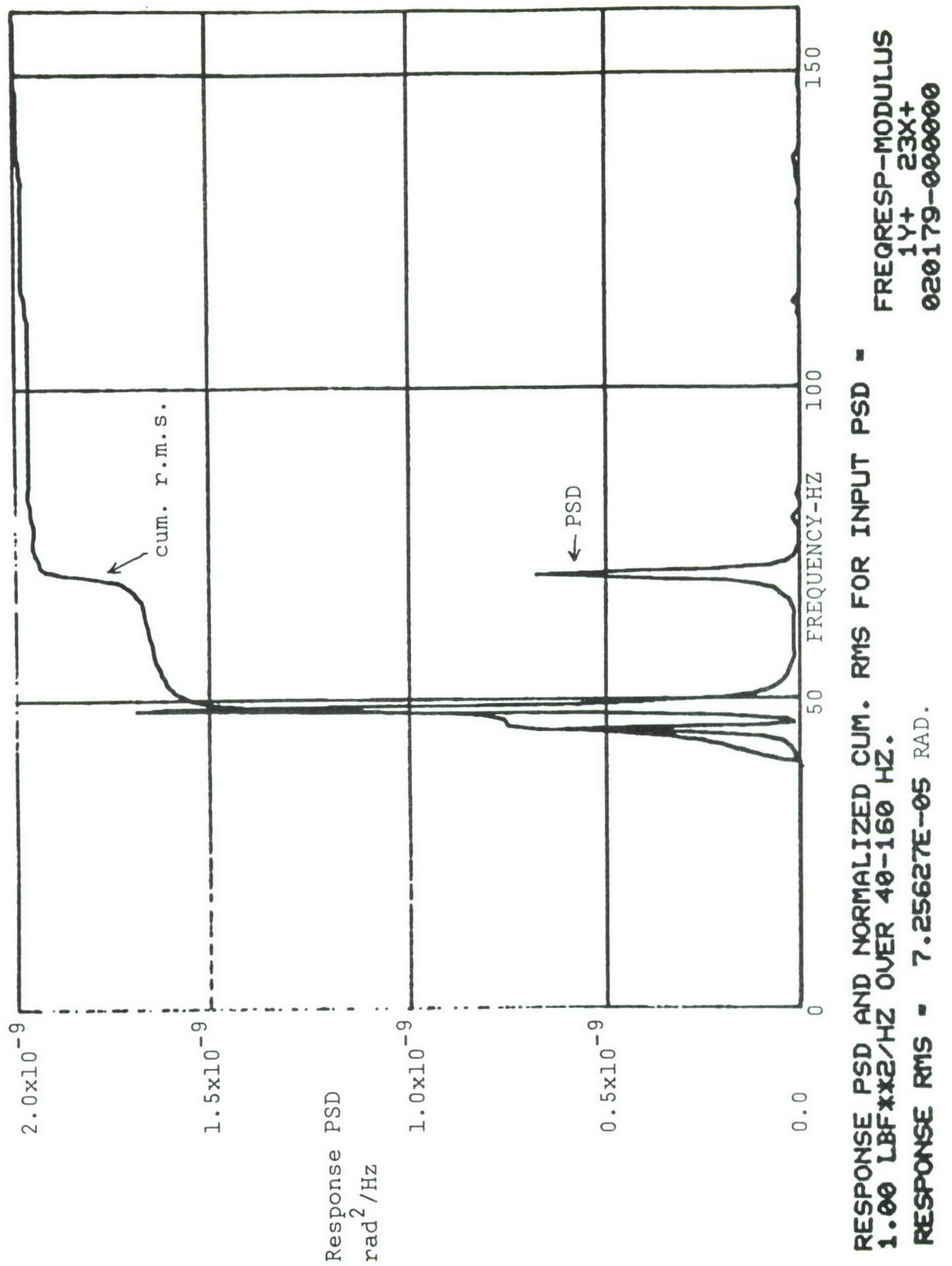
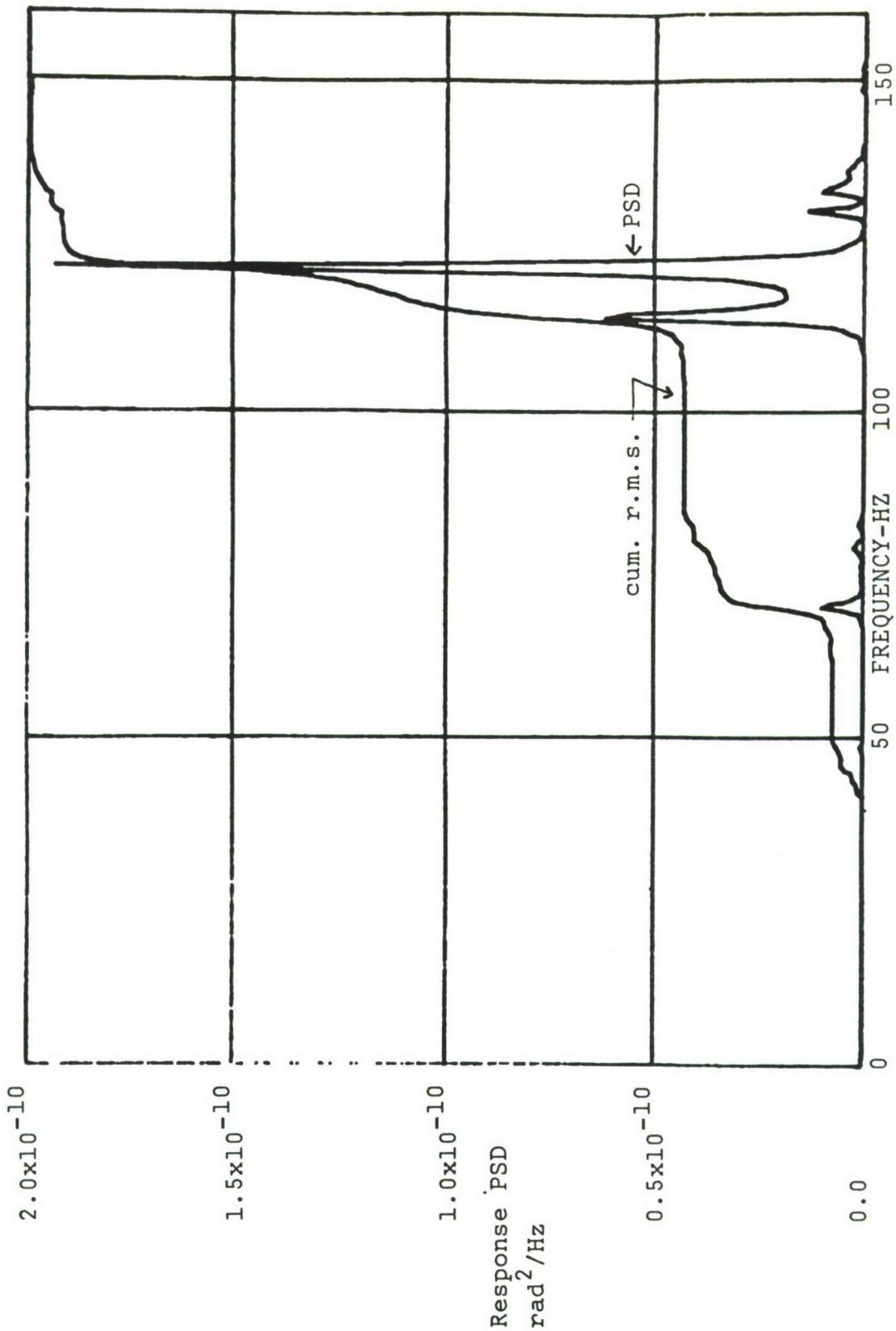


Figure 62

Experimental Results
Response Point 23, Rotation about X
Driving Point 1, Y Translation



RESPONSE PSD AND NORMALIZED CUM. RMS FOR INPUT PSD =
 1.00 LBFXX2/HZ OVER 40-160 HZ.
 RESPONSE RMS = 2.53336E-05 RAD.

FREQRESP-MODULUS
 2Y- 4X+

Figure 63

Experimental Results
 Response Point 4, Rotation about X
 Driving Point 2, Y Translation

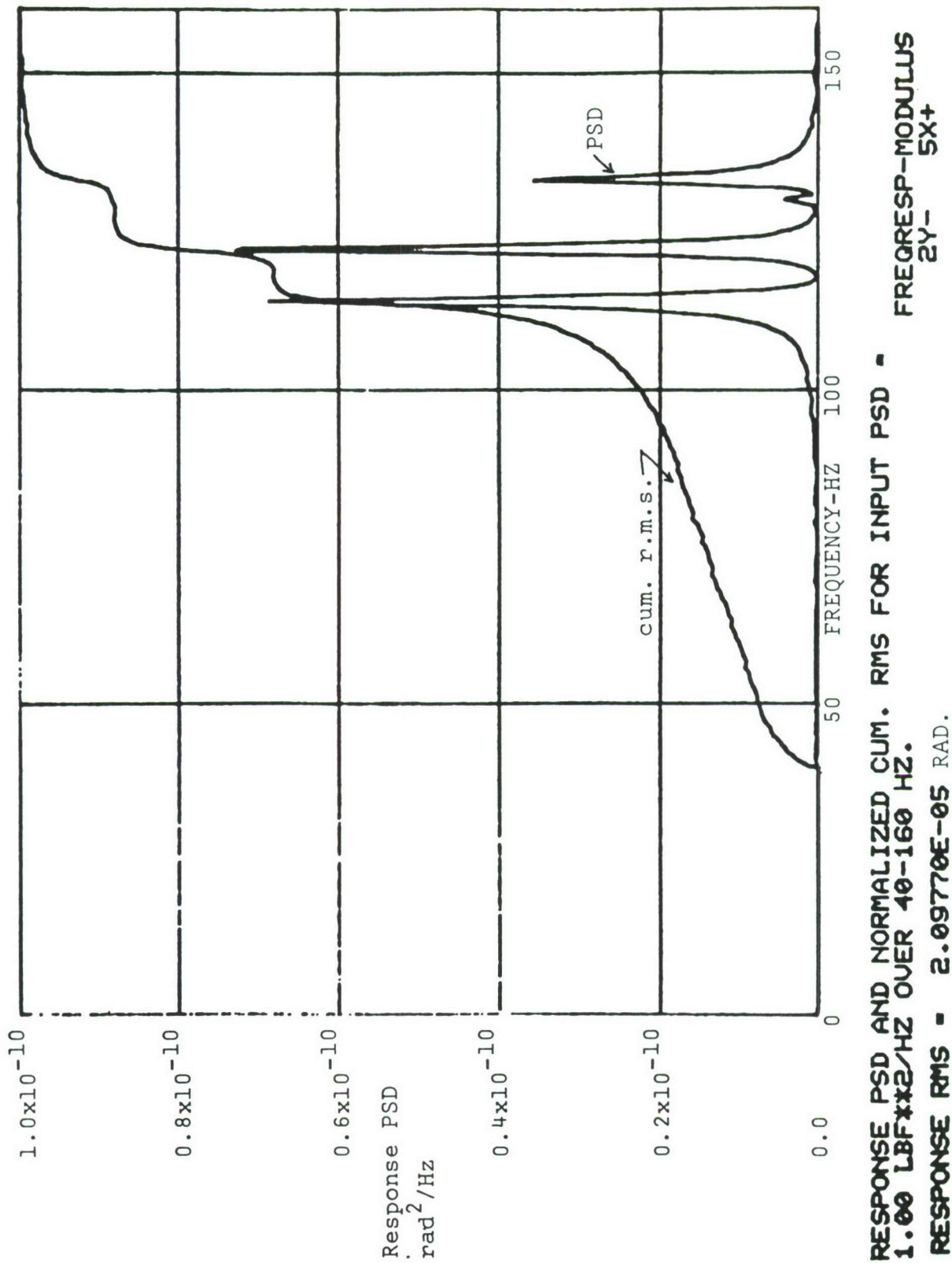


Figure 64

Experimental Results
Response Point 5, Rotation about X
Driving Point 2, Y Translation

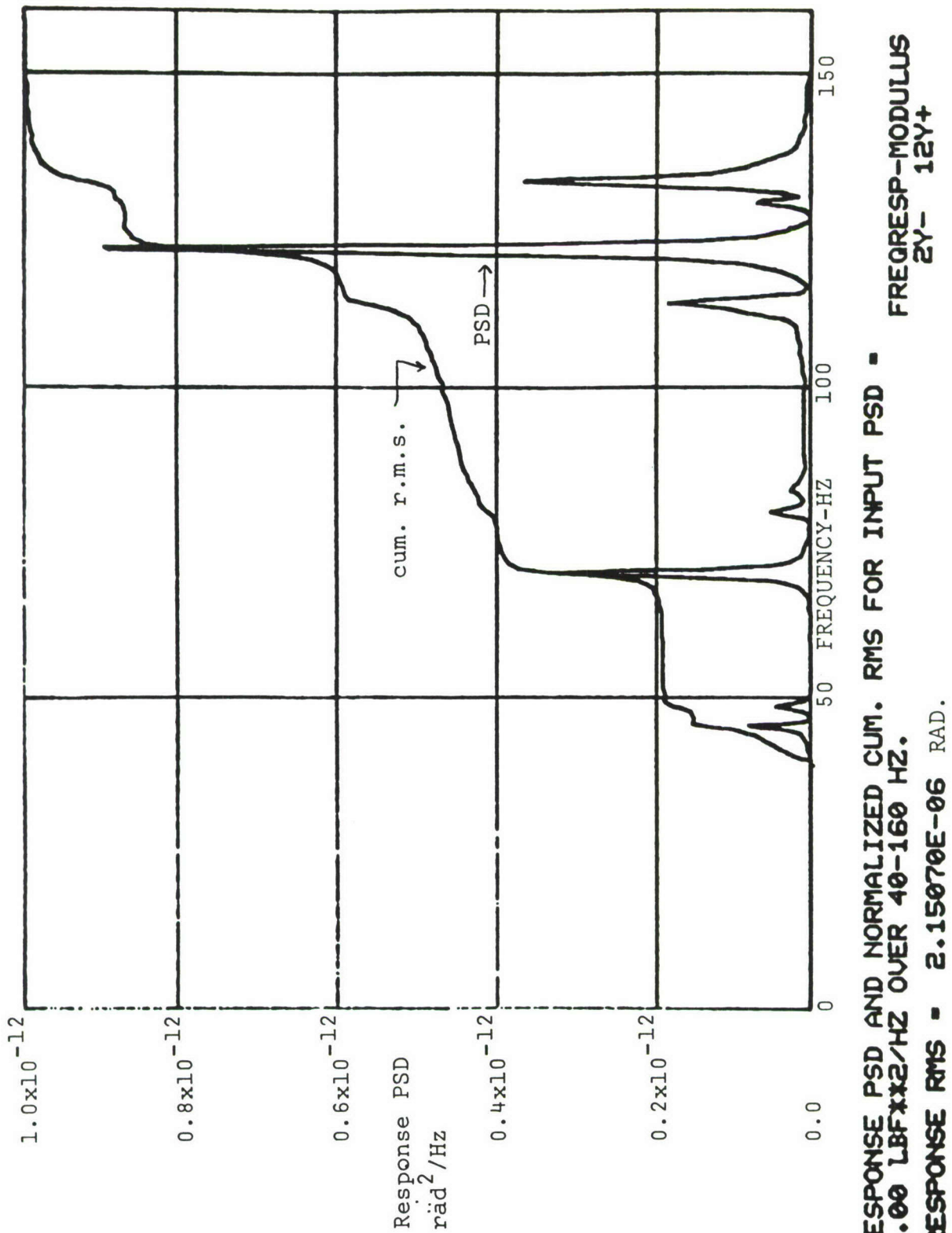


Figure 65

Experimental Results
Response Point 12, Rotation about Y
Driving Point 2, Y Translation

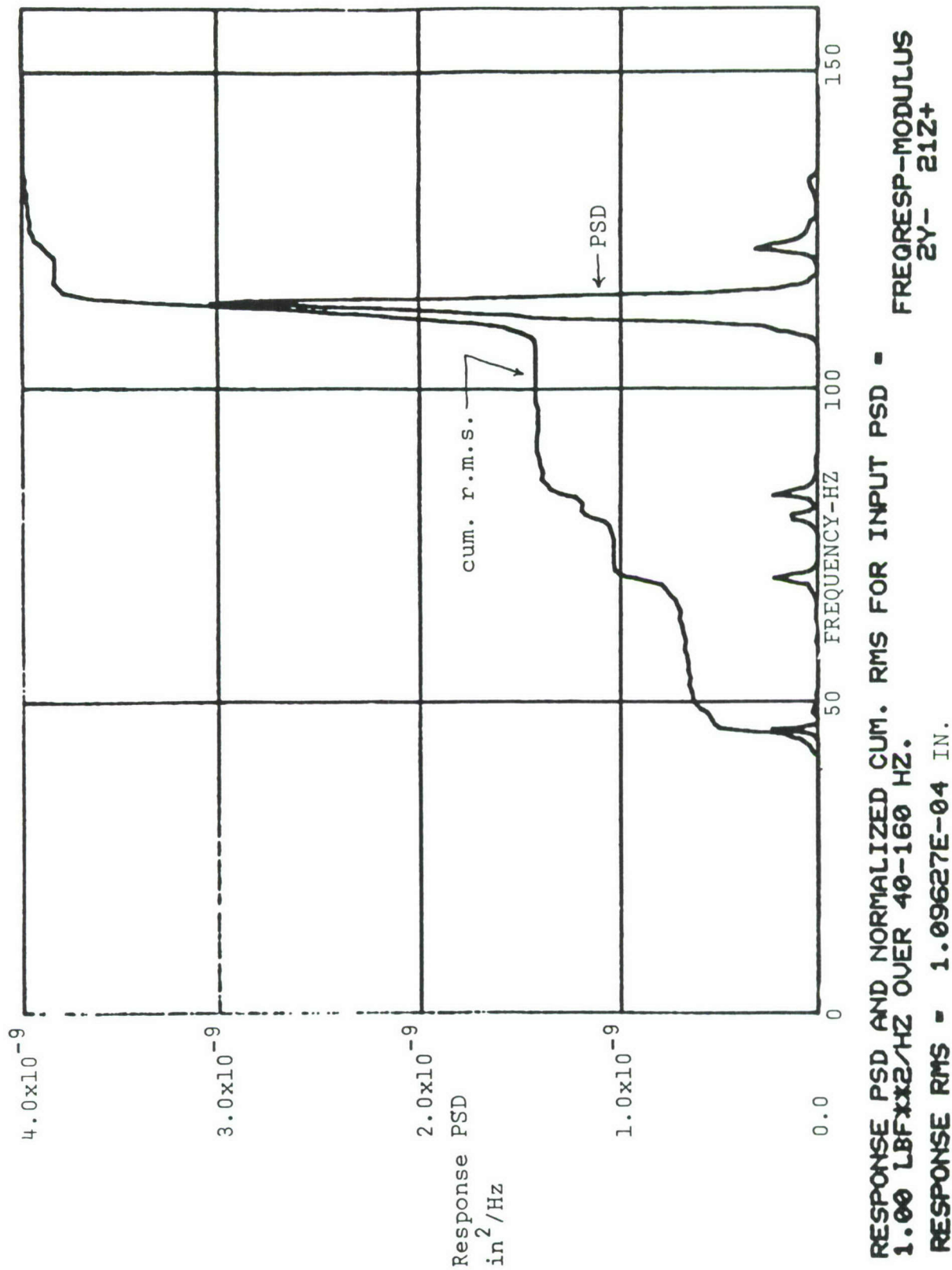


Figure 66 Experimental Results
Response Point 21, Z Translation
Driving Point 2, Y Translation

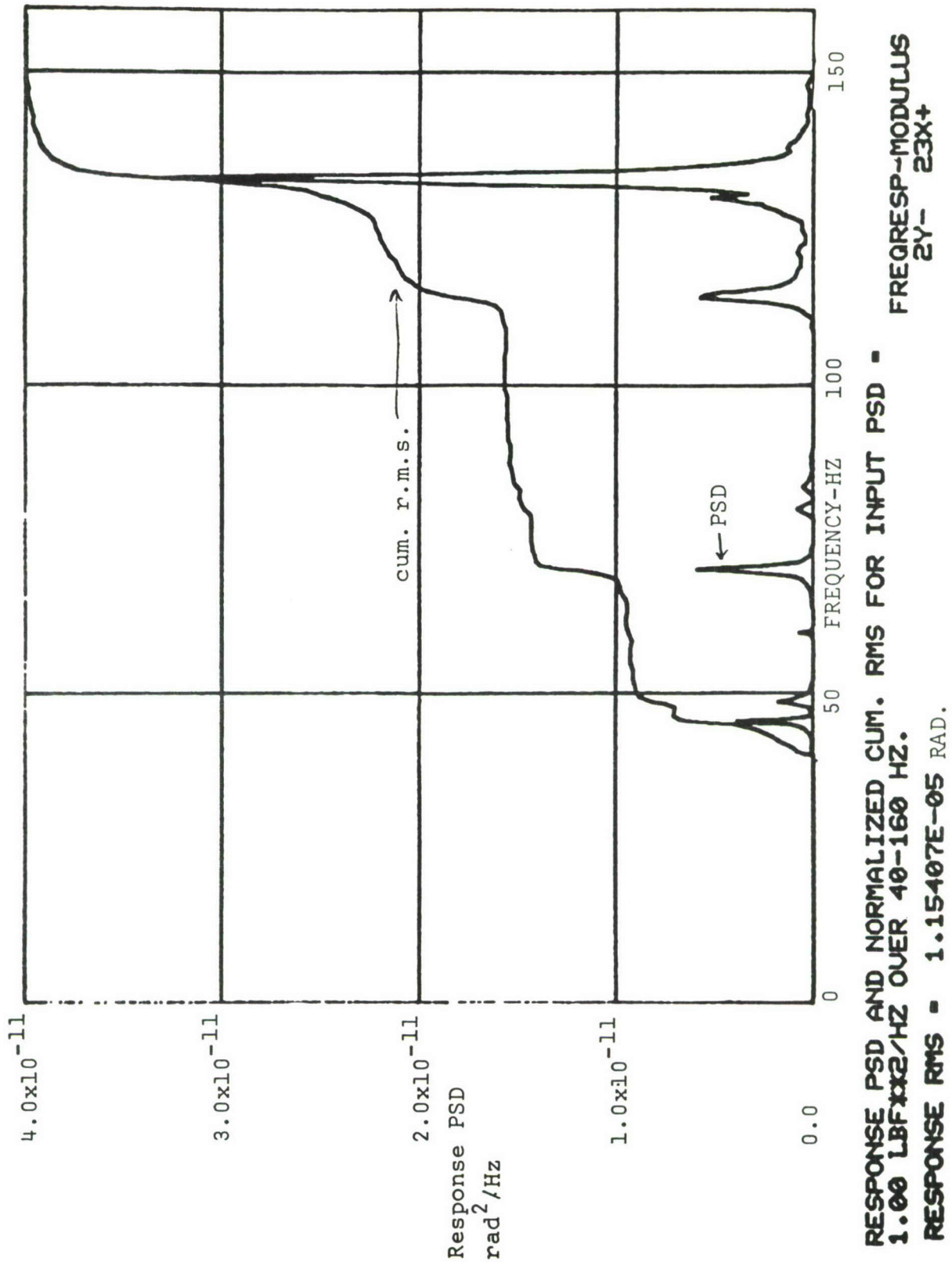


Figure 67 Experimental Results
Response Point 23, Rotation about X
Driving Point 2, Y Translation

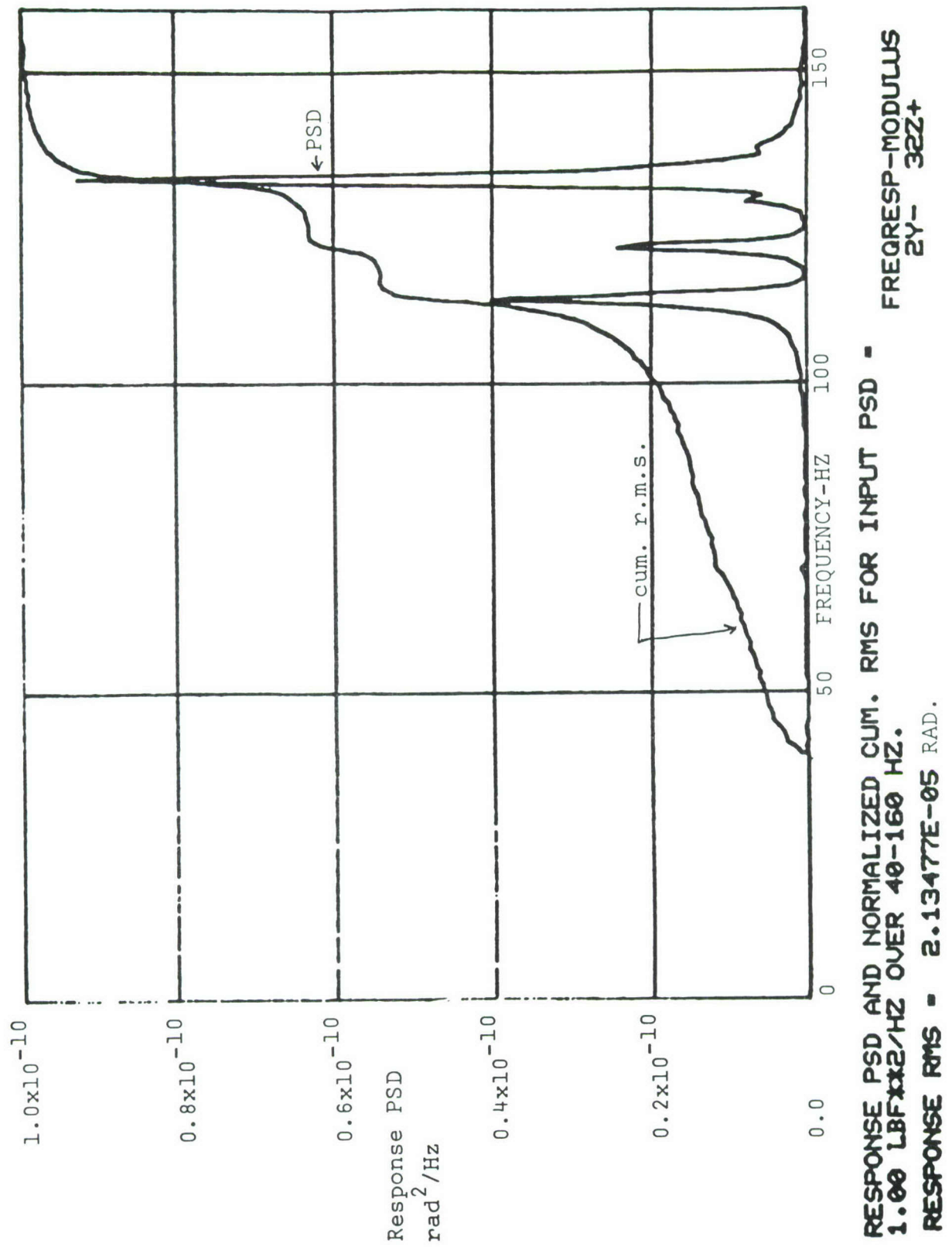


Figure 68 Experimental Results
Response Point 32, Rotation about Z
Driving Point 2, Y Translation

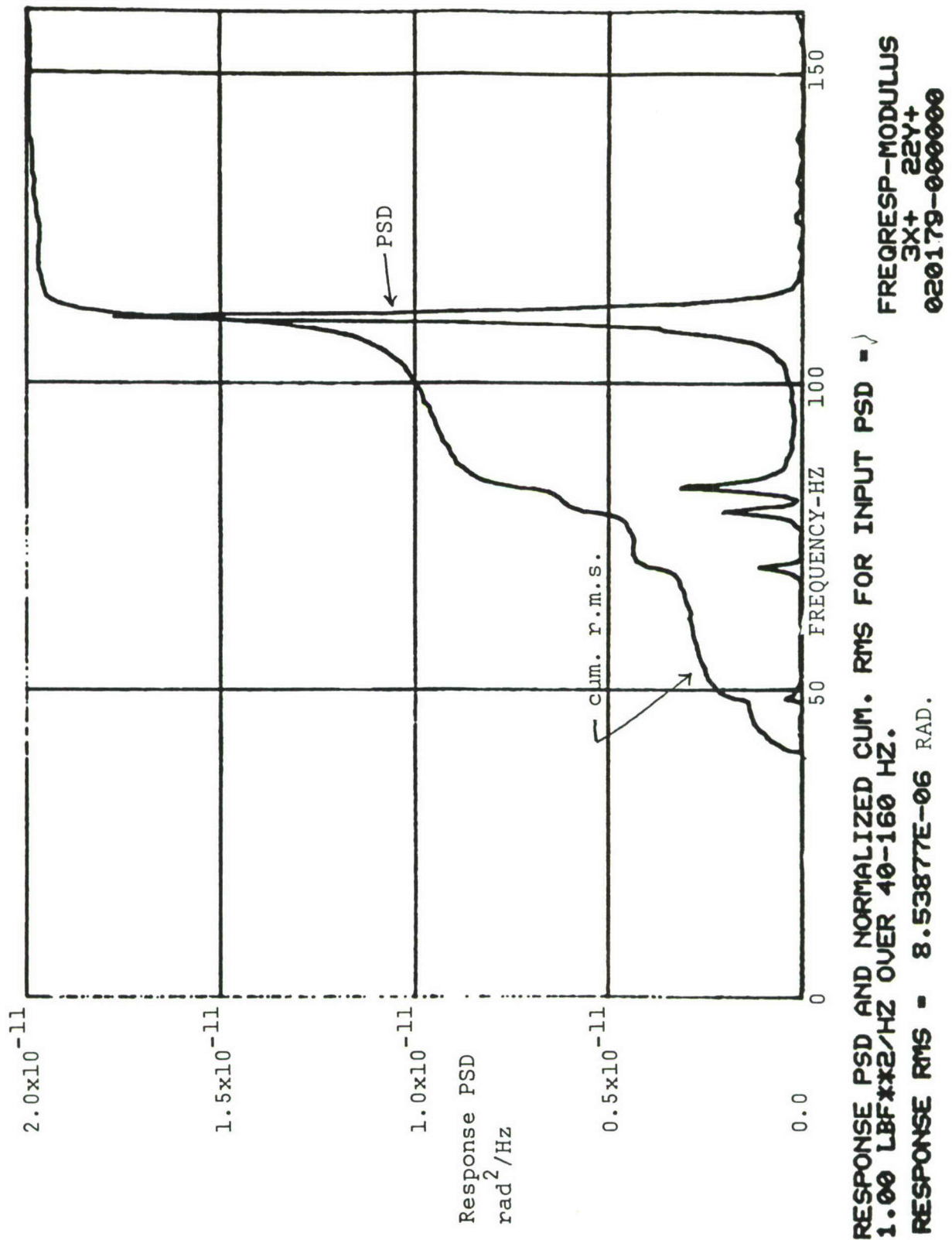


Figure 69 Experimental Results
Response Point 22, Rotation about X
Driving Point 3, X Translation

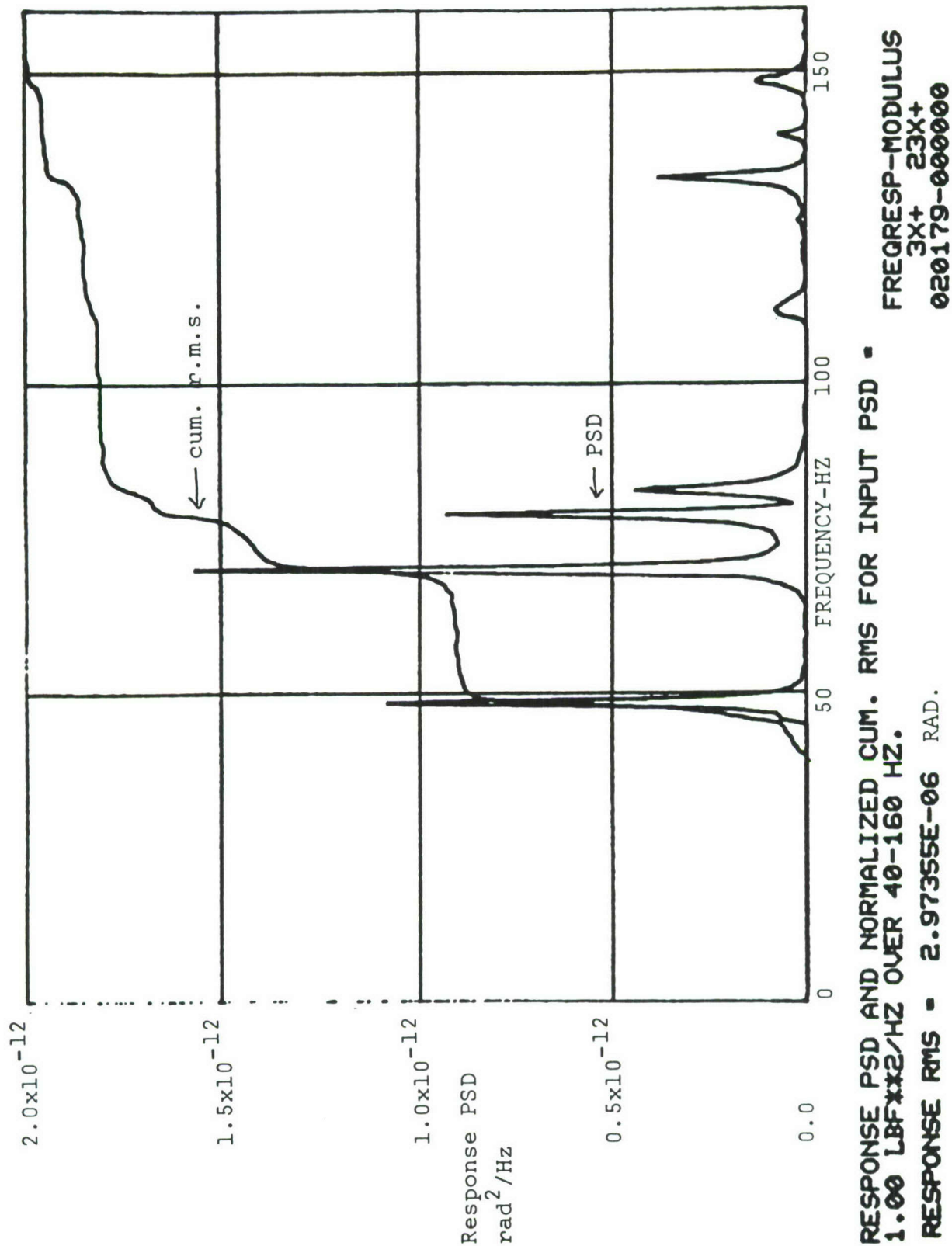


Figure 70

Experimental Results
Response Point 23, Rotation about Y
Driving Point 3, X Translation

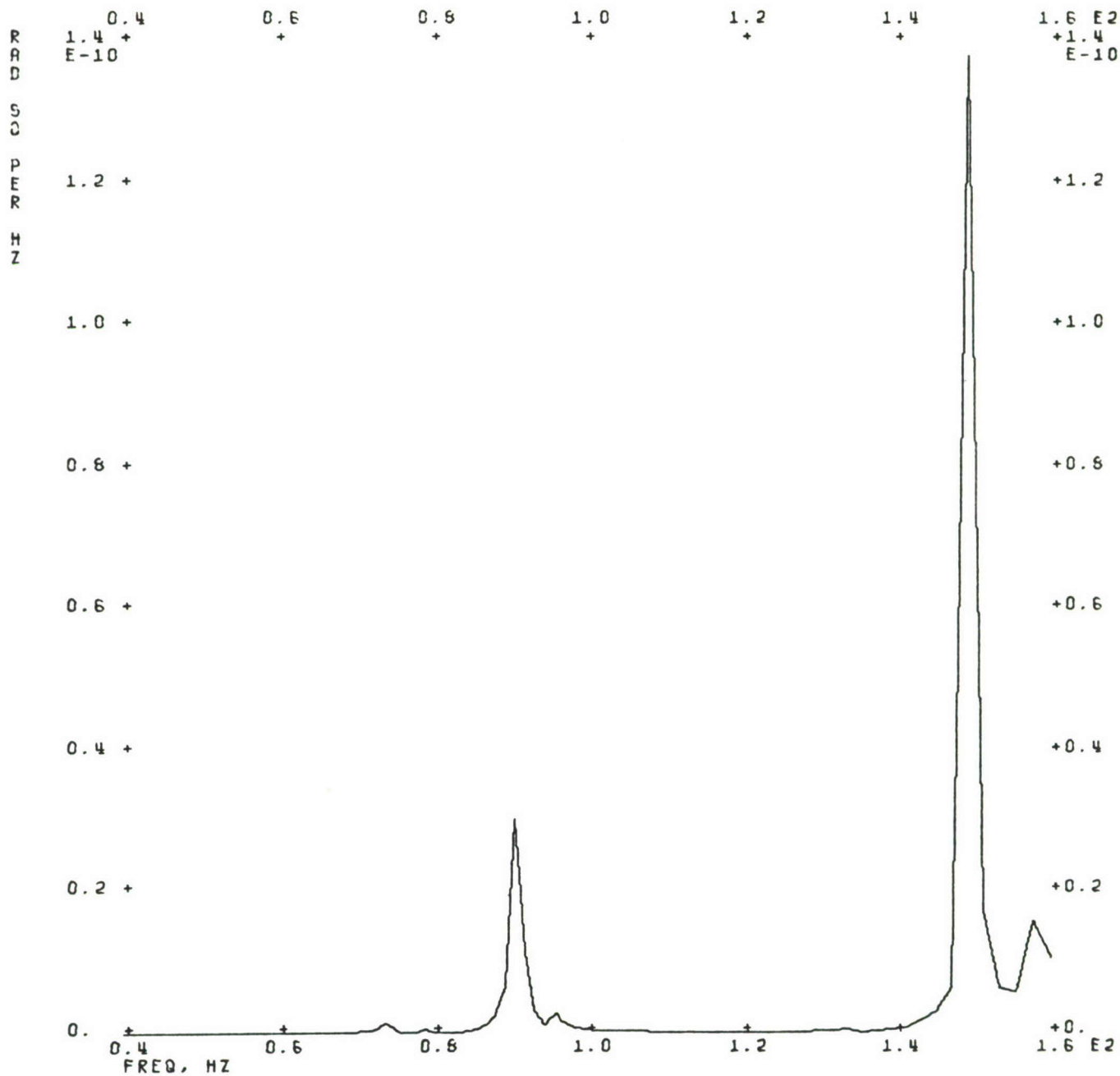


Figure 71 Analytical Results
Response Point 22, Rotation about Y
Driving Point 1, Y Translation

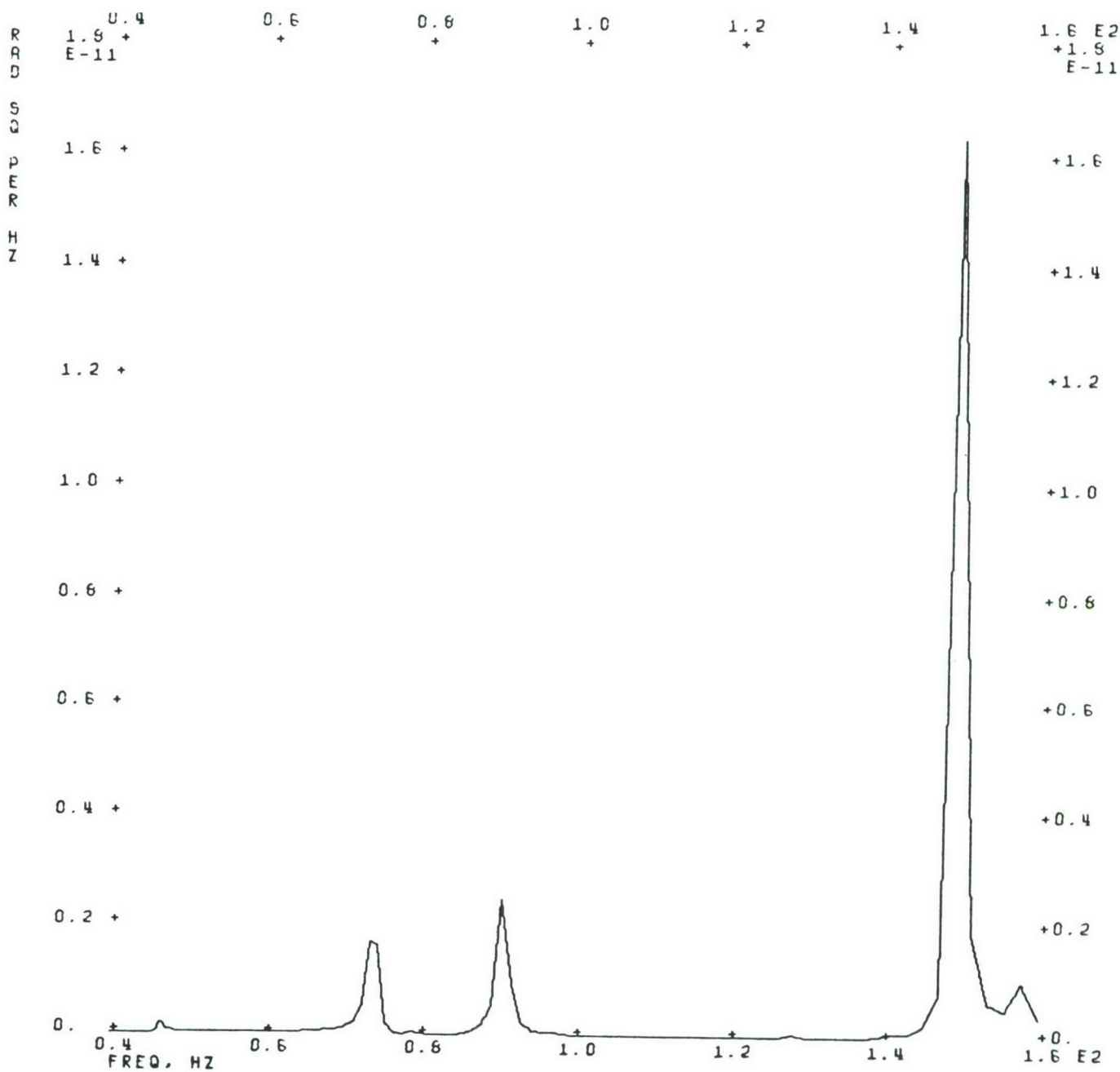


Figure 72 Analytical Results
Response Point 23, Rotation about X
Driving Point 1, Y Translation

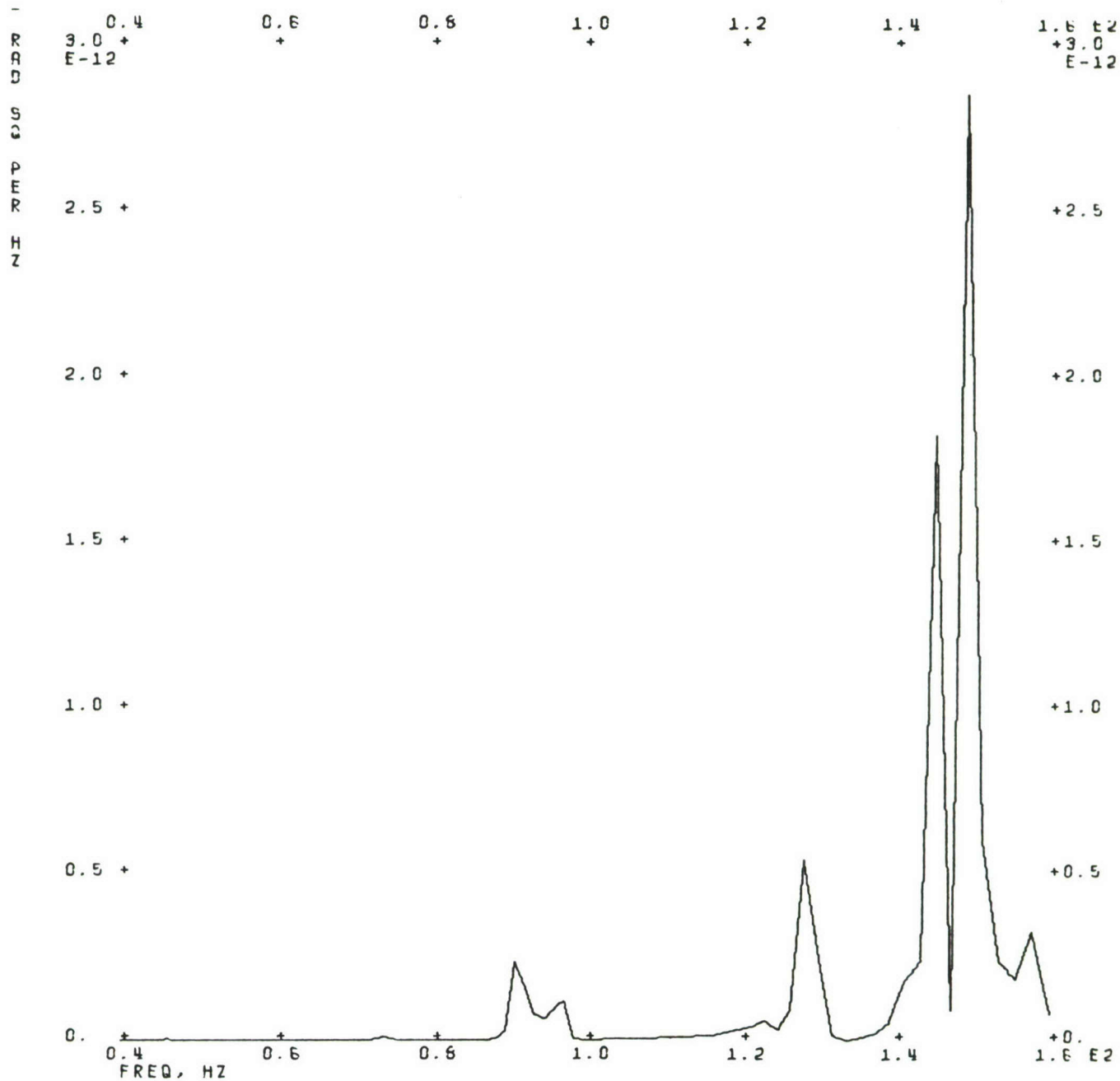


Figure 73 Analytical Results
Response Point 32, Rotation about Z
Driving Point 2, Y Translation

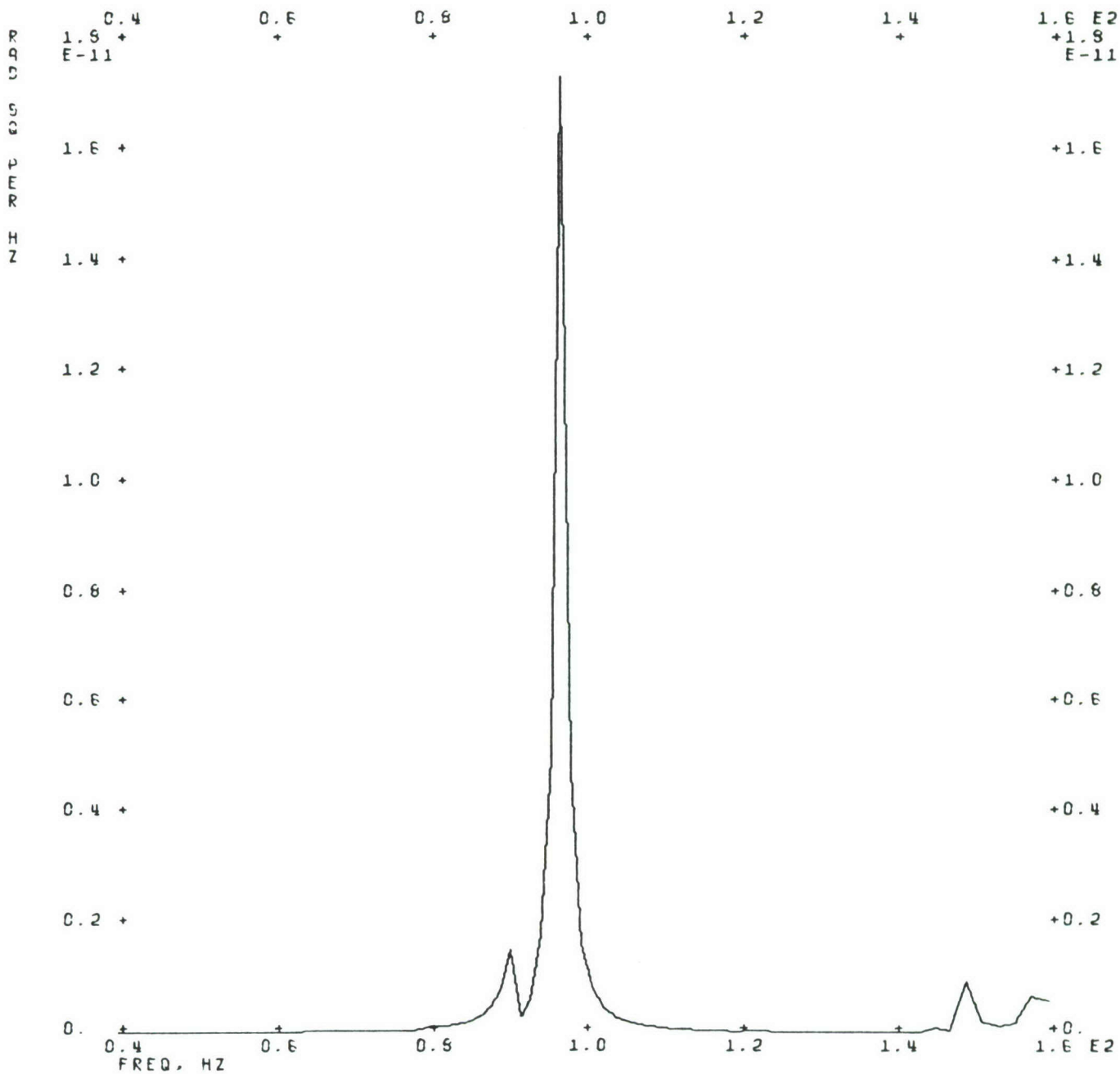


Figure 74 Analytical Results
 Response Point 22, Rotation about Y
 Driving Point 3, X Translation

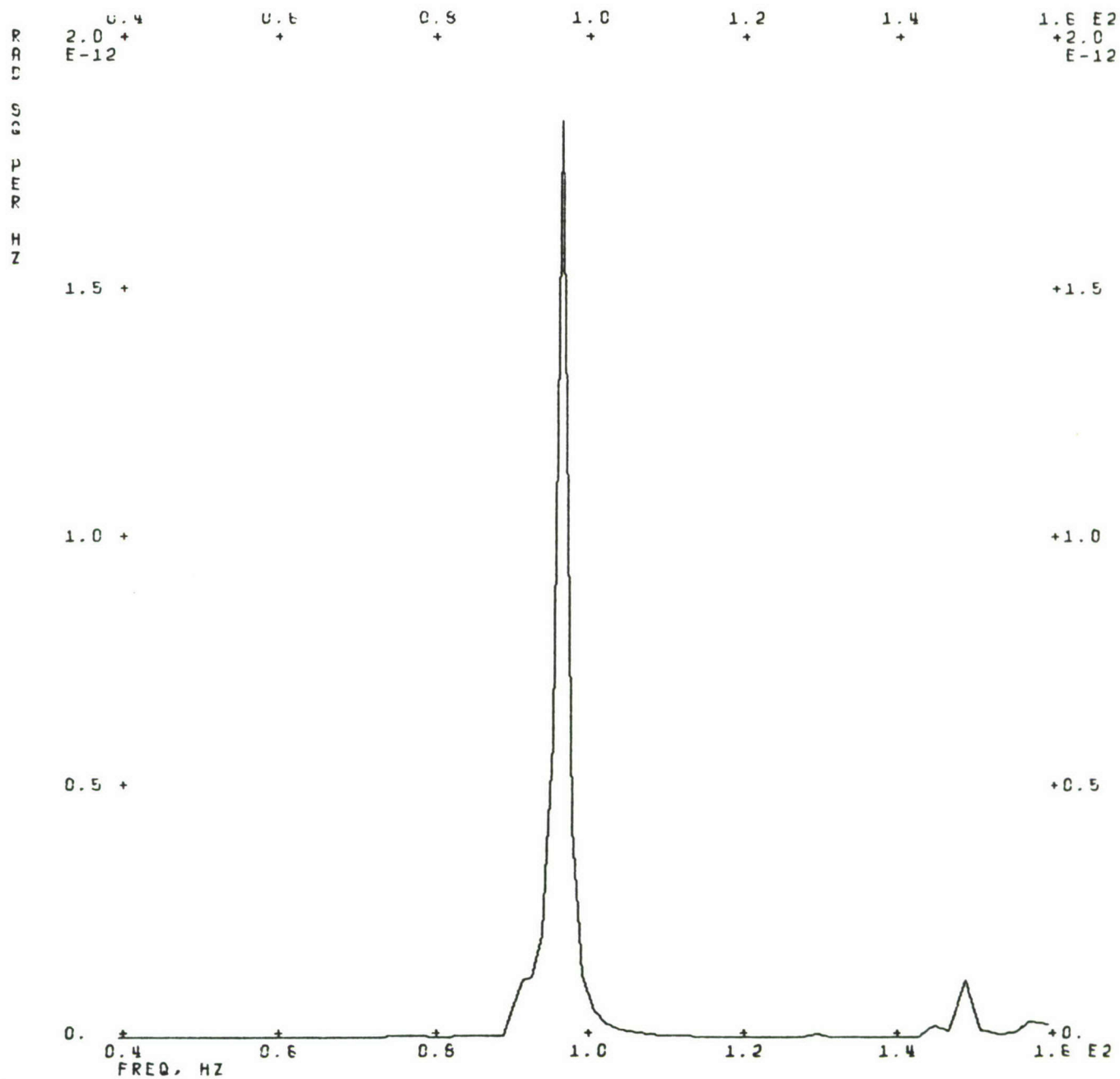


Figure 75 Analytical Results
 Response Point 23, Rotation about X
 Driving Point 3, X Translation

In retrospect, the fuselage analysis may not have been the ideal application for Semi-Loof given its present state of development. If the purpose of its development was to provide an advanced tool for analysis work, then it would have been informative to continue the policy followed in Phases I and II by modeling the same structure with both Semi-Loof and QUAD2 or other "standard" elements and then assessing the improvement in terms of accuracy per degree of freedom. Secondly, a smaller structure with equal complexity in terms of stiffeners, eccentricity, etc. would still have exercised the Semi-Loof elements just as well without the difficulties of getting large computer runs through a saturated central computer facility.

In spite of the disappointing results of the analysis, a number of valuable lessons were learned in the process. First, the participants emerged with a renewed appreciation of the amount of work involved in creating and debugging a large finite element model. Second, there was valuable feedback regarding some of the capabilities and input options needed in a large modeling job. In particular, these included convenient ways to specify beam orientations and offsets, and convenient specification of orthotropic (smeared stiffener) element orientation. Third, the exercise emphasized the number of variables or potential sources of discrepancy that can occur in an attempt to compare experimental to analytical results. Fourth, the Semi-Loof elements stood up well with respect to convenience of input and other operational details.

SECTION VIII

SUMMARY AND CONCLUSIONS

Angular vibration problems in aircraft have been studied in two distinct realms of vibration frequency. Low frequency problems have been studied from a deterministic viewpoint; that is, under the assumption that individual natural frequencies and mode shapes can be determined in some detail by analytical methods. Higher frequency problems have been approached in a statistical manner. The assumption here is that properties of individual modes cannot, and should not, be examined, but that aggregate response of mode groups can be estimated given only a limited structural description.

Considerable research and literature searching was carried out, but there was little to be found in the way of methods specifically addressed to angular vibration. Upon further reflection, it was decided that development work would have to be more general in nature, meaning that any new or improved methods would, for the most part, be relevant to vibration problems in general, and not just angular vibration.

The finite element method is the principal analytical tool currently in use for structural analysis. A number of advanced shell elements have been reported in the literature, but little of this progress has reached production users who mostly employ large, widely distributed codes such as NASTRAN. One set of curved shell elements known as "Semi-Loof" elements seemed especially promising. A skeleton computer program was obtained from its developer, Professor Bruce Irons. A preprocessing code was built around Irons' code to make it possible to use these elements with NASTRAN. This enables users to take full advantage of NASTRAN's capabilities. Small test problems showed encouraging results. A small stiffened panel was tested with results showing good agreement with Semi-Loof predictions. A large fuselage model was analyzed and tested, but with inconclusive results. A mathematical irregularity was encountered

in the quadrilateral element involving spurious mechanism modes. The problem was found to be avoidable, but still annoying.

The Semi-Loof pre- and post-processing programs are considered to be in production status, with careful attention having been given to input formats and error checking. Still, much more user experience will be needed before conclusive evaluations of these elements can be made. It is hoped that a real contribution has been made, not only to angular vibration methods, but to other shell analysis problems, as well.

For high frequency analysis, there is no widely accepted method with the flexibility and accuracy which finite elements allow in low frequency work. This is not surprising, considering the previously stated definition of the high frequency problem. During early literature searching, it was found that a rather diverse body of theory under the general name of Statistical Energy Analysis (SEA) had been developed during the 1960's in response to the difficulties of high frequency analysis. The method draws heavily on classical deterministic theories of acoustics, wave mechanics, and normal mode methods, and yet takes quite a unique view of vibration problems generally in that it uses component energies and intercomponent power flows as primary variables. SEA appeared to be well suited to the problem of predicting indirectly excited high frequency vibration of a piece of equipment mounted in an aircraft. The method is quite general and is not restricted to angular vibration.

It was decided that, due to the inexperience of the investigators with SEA, a small demonstration problem should be undertaken. The wave transmission method of SEA was used to predict equilibrium energy ratio and power transmission coefficient between two uniform plates coupled at a single degree of freedom. Comparison with experiment showed agreement which was quite encouraging in the light of the extensive assumptions involved. While physically simple, the experiment embodied the

typical factors which make high frequency predictions difficult. Mode counts were very high (over 600 total for the two plates) and individual mode shapes could not be determined. It was demonstrated how interactive digital FFT processing of measured data could be used to obtain the so-called coupling loss factors and internal loss factors of SEA. It was also demonstrated, and verified by experiment, that the basic assumptions of the wave transmission method can be utilized to estimate r.m.s. angular response given the component total energy for a uniform plate. While some software was developed in the course of the SEA experiment, it cannot yet be considered a production tool for the prediction of high frequency angular vibration.

Because of the extensive use of experimental methods for prediction of angular vibration, some effort was devoted to purely measurement problems. Differencing of output signals from translational accelerometers was identified as the most generally useful method of angular vibration measurement. Theoretical error analyses were performed to quantify the effects of intrachannel noise, interchannel gain and phase mismatch, and flexure of the mounting surface. It was concluded the measurement requirements involved in development of prediction methods could be satisfied by the differencing technique.

REFERENCES

1. Lyon, R. J., Statistical Energy Analysis for Designers, Part 1. Basic Theory and Part 2. The Engineering Application, AFFDL-TR-74-56. Part 1 - ADA 006 413, Part 2 - ADA 006 414.
2. Irons, B. M., "The Semi-Loof Shell Element," Finite Elements for Thin Shells and Curved Members, Chapter 11, ed. D. G. Ashwell and R. H. Gallagher, Wiley, London, 1976.
3. Zienkiewicz, The Finite Element Method in Engineering Service, McGraw-Hill, London, 1971.
4. Loof, H. W., "The Economical Computation of Large Structural Elements," Int. Symposium on Use of Computers in Structural Engineering, University of Newcastle-upon-Tyne, 1966.
5. Elaswaf, A., "An Application of the Semi-Loof Shell and Beam Elements," IASS World Conference on Space Enclosures.
6. Martins, R. A. F., and Owen, D. R. J., "Structural Instability and Natural Vibration Analysis of Thin Arbitrary Shells by Use of the Semi-Loof Element," J. Num. Methods in Engr., Volume 11, 1977, pp. 481-498.
7. Owen, D. R. J., and Martins, R. A. F., "Non-linear Analysis of Arbitrary Shell Structures by Use of the Semi-Loof Element."
8. NASTRAN Demonstration Problem Manual (Level 16.0), NASA SP-224(03), 1978.
9. Lyon, op cit, Section 3.2.
10. Remington, P. J., and Manning, J. E., "Comparison of Statistical Energy Analysis Power Flow Predictions with an 'Exact' Calculation," J. Acoustical Soc. Am., Vol. 57, No. 2, February 1975, pp. 374-379.
11. Manning, J. E., Cambridge Collaborative, Cambridge, Mass., personal communication.
12. Lyon, op cit, p. 38.
13. Lyon, op cit, Chapter 4.
14. Lyon, op cit, pp. 39-41.

REFERENCES (continued)

15. Lee, Jr., and Whaley, P. W., Prediction of the Angular Vibration of Aircraft Structures, AFFDL-TR-76-56, Air Force Flight Dynamics Laboratory, Wright-Patterson AFB, Ohio, June 1976; also J. Sound and Vibration, Vol. 49, December 1976, pp. 541-549.
16. Klosterman, A., and Zimmerman, R., "Modal Survey Activity via Frequency Response Functions," Soc. Automotive Engineers, Paper No. 751068, National Aerospace Engineering and Manufacturing Meeting, Culver City, CA, Nov. 17-20, 1975.
17. Richardson, M., and Potter, R., "Identification of the Modal Properties of an Elastic Structure from Measured Transfer Function Data," Instrument Society of America, Paper 74250, 20th International Instrumentation Symposium, May 21-23, 1974, Albuquerque, New Mexico.
18. Otnes, R. K., and Enochson, L., Applied Time Series Analysis, Section 5.8, Wiley-Interscience, 1975.
19. Johnson, C. D., Gibson, W. C., Kienholz, D. A., and Paxson, E. B., Jr., "Interim Technical Report Angular Vibration of Aircraft, Phase II," Anamet Report 877.501C, October, 1978, under Contract No. F33615-77-C-3050.
20. Lyon, op cit, pp. 36-38.
21. Lyon, op cit, pp. 46-48.
22. Anamet Laboratories R & D Status Report, May 1978, Air Force Contract No. F33615-77-C-3050, on Angular Vibration of Aircraft.
23. Lee, J., and Whaley, P., "Prediction of the Angular Vibration of Aircraft Structures," J. Sound and Vibration (1976), Vol. 49, No. 4, pp. 541-549.
24. Anamet Laboratories R & D Status Report, June 1978, Air Force Contract No. F33615-77-C-3050, on Angular Vibration of Aircraft.
25. Whaley, P. W., and Obal, M. W., "Measurement of Angular Vibration Using Conventional Accelerometers," Shock and Vibration Bulletin, September 1977, pp. 97-107.
26. ten Wolde, T., Verheij, J. W., and Steenhoek, H. F., "Reciprocity Method for the Measurement of Mechano-Acoustical Transfer Functions," J. Sound and Vibration, Vol. 42, No. 1, 1975, pp. 49-55

REFERENCES (continued)

27. Bendat, J. S., and Piersol, A. G., Random Data: Analysis and Measurement Procedures, Wiley-Interscience, 1971, pp. 153-156.
28. Papoulis, A., Probability, Random Variables, and Stochastic Processes, McGraw-Hill, 1965, pp. 343-344.
29. Bendat and Piersol, op cit, pp. 141-147.
30. Lyon, op cit, p. 43.
31. Flügge, William, Stresses in Shells, Springer Verlag, Berlin, 1962, pp. 221-222.
32. Irons, B. M., "The Semi-Loof Shell Element," Symposium on Thin Shells and Curved Members, University College, Cardiff, Wales, April 1974.

APPENDIX A

SEMI-LOOF USER'S MANUAL

A.1 INTRODUCTION

This appendix is intended to supply all the information needed to prepare finite element models using Semi-Loof shell and beam elements. The reader is assumed to have read Section III of this report, which presents the background and mathematical nature of Semi-Loof elements. In addition, an "intermediate" level of NASTRAN expertise is assumed. This implies fairly extensive experience with medium- to large-scale modeling problems. No facility with the DMAP language is assumed, however. DMAP Alter packages described in this appendix are designed to be used as "black boxes."

The remainder of this appendix is organized as follows:

Section A.2 is a general discussion of the Semi-Loof element, reviewing the capabilities, assumptions, and precautions that users must have in mind.

Section A.3 gives instructions for preparing input data for PRELOOF, the Semi-Loof preprocessing routine. For the most part, this data is prepared in the NASTRAN "BULK DATA" format.

Section A.4 discusses POSTLOOF, the post-processing routine for recovery of rotations and stresses.

Section A.5 presents deck setups for CDC computers. The Semi-Loof coding has not been implemented on any other machines at this time.

Finally, Section A.6 is a compilation of programmer's notes for those who are interested in programming details.

Listings of control card procedures, Fortran source code for PRELOOF and POSTLOOF, and DMAP Alters appear in Appendix B.

A.2 GENERAL DISCUSSION

The following is a summary of the salient features of the Semi-Loof shell and beam elements from the user's point of view.

First, the shell elements are strictly for thin shell analysis, as reflected in the assumption of zero transverse shear strain at selected points on the shell. The decision as to what constitutes a thin shell is an engineering judgement, based on the R/t ratio. R is a radius of curvature for a curved shell, or a typical span for a flat plate. R/t should normally exceed 10 for a thin shell.

The shell elements represent both stretching and bending action. For flat plates, one may suppress one or the other of these actions. Assuming the plate lies in the x - y plane, one would constrain degrees of freedom 1 and 2 to eliminate stretching, or 3, 4, 5, and 6 for bending. (Even though 4, 5, and 6 are already suppressed by permanent single point constraints at corner nodes, and 6 at Loof nodes, it is not an error to constrain these redundantly on SPC cards. See the discussion of permanent SPC's in the next section.)

It is important to keep in mind the nature of the degrees of freedom for these elements. As explained in Section 3.4, each node point uses degrees of freedom 1, 2, and 3 as translations, like other NASTRAN elements, but dof 4, 5, and 6 are different. For the shell element, no rotation degrees of freedom are defined at corner points. This does not mean that rotations are zero at such points, but only that these rotations are not independent degrees of freedom. At mid-side nodes, dof 4 and 5 represent rotations about the tangent to the element side. These rotations are not actually located at the mid-side node, but at the Loof nodes, located at about 58% ($1/\sqrt{3}$) of the distance from the mid-side to either corner. Thus the values of dof 4 and 5 are not usually of interest at mid-side nodes, but their nature must be understood when it comes to applying constraints, loads, or especially joining other elements. If another NASTRAN bending element were joined to the mid-side node

of a Semi-Loof shell, the NASTRAN element would interpret dof 4 and 5 as rotations about the x and y axes at that node point, while the Semi-Loof element would interpret them as explained above. This error would not be detected by the software. The Semi-Loof beam element defines all three translations and three rotations in the normal NASTRAN manner at the two end points. The middle node uses dof 4 and 5 as Loof rotations for compatibility with the shell elements.

From this discussion it may be seen that the mid-side nodes are in a sense subservient to the corner nodes. That is, each mid-side node serves to define the shape of the arc joining two corner nodes. The fact that L00F8 has eight nodes should not be construed to mean that L00F8 is an eight-sided element. In setting up a model, one should first locate the corner nodes and then place the mid-side nodes at the center of each element edge. Although tests have shown that a considerable deviation in location from the center point can be tolerated without losing much accuracy, still there is usually no reason not to pick the center point, so this practice is recommended.

At this time, no studies have been made with regard to aspect ratios, and particularly, what is a practical upper bound on aspect ratio. No problem is anticipated, due to the isoparametric formulation. It is known that no element side should not have too severe curvature. A 30° arc seems to be a good upper limit.

L00F8 elements are preferred over L00F6 triangle elements where the geometry permits. Two triangles would have five more degrees of freedom than a L00F8 element covering the same area, and probably less accuracy.

A special feature of L00F8 and L00F6 is their accommodation of variation of element thickness and/or surface pressure over the surface of any element. In these cases, thicknesses and/or pressures are assigned to grid points, and the grid point values are interpolated over the element surface. For constant thickness elements, thicknesses are assigned on element connection

cards, and these values override any grid point thickness values. Variation of pressure would be useful for representing hydrostatic or hydrodynamic fluid pressures, for example.

L00F8 and L00F6 can reference isotropic and orthotropic material properties. In addition, there is a provision that is useful in modeling minor stiffening ribs in airframe structures. The most economical and realistic strategy in these cases is to "smear" the minor ribs, that is, to distribute their stiffness to the attached panel. This is done by calculating properties for an anisotropic panel that has the same stretching and bending stiffnesses as the panel-rib assemblage. This approach will, of course, not provide accurate local representation of deformations but should be a good overall model. The following formulas taken from Reference [31] may be useful in such cases:

$$A_{11} = \frac{Et}{1-\nu^2} + \frac{EA_1}{b_1}$$

$$A_{22} = \frac{Et}{1-\nu^2} + \frac{EA_2}{b_2}$$

$$A_{12} = \frac{\nu Et}{1-\nu^2}$$

$$A_{33} = \frac{Et}{2(1+\nu)}$$

$$B_{11} = \frac{EA_1 C_1}{b_1}$$

$$B_{22} = \frac{EA_2 C_2}{b_2}$$

$$D_{11} = \frac{Et^3}{12(1-\nu^2)} + \frac{E(I_1 + A_1 C_1^2)}{b_1}$$

$$D_{22} = \frac{Et^3}{12(1-\nu^2)} + \frac{E(I_2 + A_2 C_2^2)}{b_2}$$

$$D_{12} = \frac{\nu Et^3}{12(1-\nu^2)}$$

$$D_{33} = \frac{Et^3}{12(1+\nu)} + \frac{G}{2} \left(\frac{J_1}{b_1} + \frac{J_2}{b_2} \right)$$

(A.1)

Here A_{ij} , B_{ij} , and D_{ij} are force-strain parameters coupling force and moment resultants to strains and curvatures:

$$\begin{Bmatrix} N_1 \\ N_2 \\ N_{12} \\ M_1 \\ M_2 \\ M_{12} \end{Bmatrix} = \begin{bmatrix} A_{11} & A_{12} & 0 & | & B_{11} & 0 & 0 \\ A_{12} & A_{22} & 0 & | & 0 & B_{22} & 0 \\ 0 & 0 & A_{33} & | & 0 & 0 & 0 \\ \hline B_{11} & 0 & 0 & | & D_{11} & D_{12} & 0 \\ 0 & B_{22} & 0 & | & D_{12} & D_{22} & 0 \\ 0 & 0 & 0 & | & 0 & 0 & D_{33} \end{bmatrix} \begin{Bmatrix} \epsilon_1 \\ \epsilon_2 \\ \epsilon_{12} \\ \chi_1 \\ \chi_2 \\ \chi_{12} \end{Bmatrix} \quad (A.2)$$

In these expressions, subscripts 1 and 2 represent two orthogonal directions. The beam properties referenced in Eq. (A.1) are

I_i = moment of inertia of a rib about its own center of gravity

A_i = area

b_i = spacing of ribs

C_i = eccentricity

J_i = torsion constant of a single rib

Also,

E, ν = material properties of both ribs and panel

t = panel thickness

Figure A-1 illustrates these properties

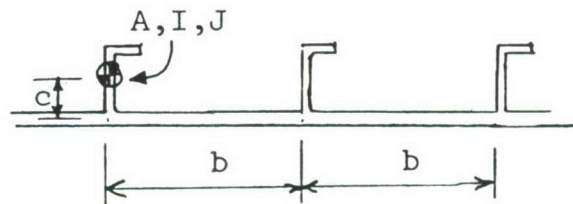


Figure A-1 Stiffened Panel Properties

It is important to get the signs of the B terms right. They are positive if the stiffeners are located on the positive side of the element; negative otherwise. The positive side is

determined by taking the cross product of unit vectors in the ξ and η directions (e.g. $\underline{V}_\zeta = \underline{V}_\xi \times \underline{V}_\eta$, see Figure A-2). The resulting vector points outward on the positive side of the element. An equivalent indicator is to proceed around the perimeter of the element in the sequence prescribed on the CLOOF card and use the right-hand rule.

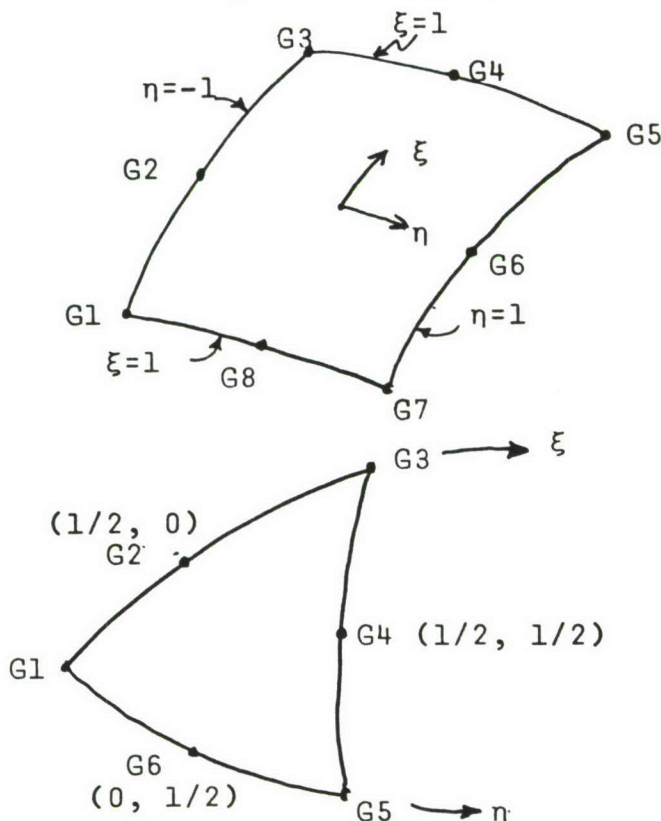


Figure A-2 Local Coordinates for L00F8 and L00F6 Elements

The L00F3 curved beam element can represent non-prismatic beams. In this case, cross-section properties are prescribed at each node point and are interpolated quadratically over the length of the element.

The following items are not incorporated in the Semi-Loof elements at this time:

- (1) Geometric stiffness
- (2) Thermal loads
- (3) Heat transfer matrices
- (4) Line loads

Two computer programs have been written to interface Semi-Loof to NASTRAN: PRELOOF and POSTLOOF. PRELOOF reads the bulk data deck and interprets certain bulk data cards that describe Semi-Loof elements. It also recognizes GRID cards and certain other bulk data cards. It then generates and assembles stiffness matrices, mass matrices, and distributed loads and writes them on a binary file in a form suitable for input to NASTRAN. The data cards, with Semi-Loof connection and property cards deleted, are then passed to NASTRAN. NASTRAN then saves output information in a form that can be read by POSTLOOF which recovers angular deformations or stresses as requested by the user.

Most of this process of transferring files into and out of NASTRAN has been automated by means of control card procedures and DMAP Alters, and can be set up by the user with a minimum knowledge of computer control details. As mentioned earlier, a moderate amount of experience with NASTRAN is assumed, however. The following sections describe the steps required to use PRELOOF and POSTLOOF.

A.3 PRELOOF INPUT

Input to PRELOOF consists of a complete NASTRAN input file: Executive Control, Case Control, and Bulk Data. This file is read by PRELOOF, modified, and then passed on to NASTRAN for execution. Some of the data cards are recognized by PRELOOF and then removed from the deck, some are recognized and retained, and others are passed directly to NASTRAN without any processing. These cards are described below and summarized in Table A-1.

The UPDATE utility is used to manipulate the input deck, so the first card must always be

*DECK X

which is immediately followed by a title card. Then begins the Executive Control Deck, which contains the cards ID, TIME, APP, SOL, DIAG, etc., as in a normal NASTRAN run. Next there is provision for a number of options which govern PRELOOF. These are listed below, with the default option underlined:

\$LOOFECHO	$\left\{ \begin{array}{c} \text{YES} \\ \underline{\text{NO}} \end{array} \right\}$	If yes, echo bulk data deck as read.
\$LOOFPLOT	$\left\{ \begin{array}{c} \text{YES} \\ \underline{\text{NO}} \end{array} \right\}$	If yes, generate PLOTTEL cards representing LOOF8 and LOOF6 elements. The PLOTTEL element numbers are $10*N+i$ where N is the LOOF element number and i ranges from 1 to 8 for LOOF8 or 1 to 6 for LOOF6.
\$LOOFMAT	$\left\{ \begin{array}{c} \text{YES} \\ \underline{\text{NO}} \end{array} \right\}$	If yes, generate matrices for NASTRAN. If no, process bulk data but do not generate matrices (e.g. for plots or debug).
\$LOOFDYN	$\left\{ \begin{array}{c} \text{YES} \\ \underline{\text{NO}} \end{array} \right\}$	If yes, generate mass matrix for dynamics problem. If no, generate stiffness but not mass matrix.
\$LOOFDEBUG	$\left\{ \begin{array}{c} \text{YES} \\ \underline{\text{NO}} \end{array} \right\}$	If yes, print out element matrices for diagnosis.
\$NASTRAN	$\left\{ \begin{array}{c} \text{MSC} \\ \text{COSMIC} \end{array} \right\}$	Selects either MSC/NASTRAN or COSMIC NASTRAN. Default set to correspond to each particular installation.
\$LOOFGAUSS	$\left\{ \begin{array}{c} 2 \\ \underline{3} \end{array} \right\}$	Integration rule: 2 means 2x2, 3 means 3x3 integration.

\$LOOFFACT $\left\{ \frac{0.0}{X} \right\}$ Sets a factor to be used for a fifth point in 2x2 integration. 0.2 is a recommended value for X when this option is used.

Following the \$ option cards, the following must be included to access rigid format alters:

*READ RFALTS

followed by

ENDALTER

CEND

the end of the executive control deck.

Next is the case control, which is the same as for any NASTRAN run with the following exception: when pressure loads are generated by PRELOOF there must not be any LOAD cards in the case control deck. This implies only a single load case.

The bulk data deck is somewhat restricted in format as compared with normal NASTRAN capabilities which the reader is assumed to know. The following rules apply:

(1) The first field of each card must be left-adjusted.

(2) Among the cards recognized by PRELOOF (listed below), continuation cards must immediately follow their parents. Also, among these cards, only GRID may use the double field option (i.e. GRID*).

(3) The bulk data cards must be sorted into groups, with groups in the order listed below. Within each group, cards may be in any order as long as continuation cards immediately follow their parents.

Group 1: CORD2R, CORD2C, and CORD2S cards. These cards must reference the basic coordinate system in field 3.

Group 2: GRID cards (the double field GRID* may also be used). Only coordinate systems defined in group 1 (or the basic system) may be used. Permanent single-point constraints must not be used. They are established by PRELOOF.

Group 3: SEQGP cards. The BANDIT bandwidth optimization code has been modified to recognize CLOOF8, CLOOF6, and CLOOF3 cards.

Group 4: CL00F8 and CL00F6 element connection cards for quadrilateral and triangular Semi-Loof elements, respectively. These cards are described in detail on pages 235 through 238.

Group 5: CL00F3 element connection cards for curved beam elements. See page 232 for format.

Group 6: PL00F cards giving properties of isotropic L00F8 and L00F6 elements. See page 241.

Group 7: PL00FX cards giving properties for anisotropic L00F8 and L00F6 elements. See page 243.

Group 8: PL00F3 cards giving properties for L00F3 elements. See page 242.

Group 9: MAT1 cards. As per the NASTRAN manual.

Group 10: MAT8 cards for homogeneous orthotropic shells (for smeared stiffeners see PLO0FX). See page 239.

Group 11: PLOAD4 cards defining pressures at grid points. See page 240.

Group 12: THL00F cards defining thicknesses at grid points. See page 245.

Group 13: PARAM cards. Only PARAM WTMASS, which has the same function as in standard NASTRAN runs, is used.

Group 14: SPOINT cards. Scalar point identification numbers must be greater than the highest grid point number.

Following these cards, all bulk data cards not processed by PREL00F are included. See Table A-1 for a summary of the foregoing.

TABLE A-1
SUMMARY OF PRELOOF INPUT DECK

*DECK X	
title card	

ID	Executive
TIME	Control
APP	Deck
SOL	
DIAG	
.	
.	
.	
\$LOOF options	
*READ RFALTS	
ENDALTER	
CEND	

TITLE=----	Case
.	Control
.	Deck
.	
BEGIN BULK	

CORD2R, CORD2C, CORD2S	Bulk
GRID	Data
SEQGP	Deck
CLOOF8, CLOOF6	
CLOOF3	
PLOOF	
PLOOFX	
PLOOF3	
MAT1	
MAT8	
PLOAD4	
THLOOF	
PARAM	
SPOINT	
other Bulk Data Cards	
ENDDATA	

BULK DATA DECK

Input Data Card CL00F3 Connections for L00F3 Curved Beam Element

Description: Defines a Semi-Loof beam element of the structural model.

Format and Example:

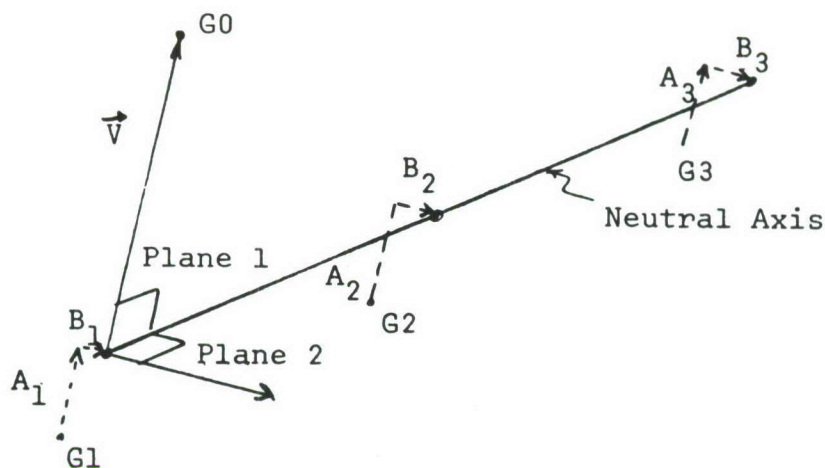
1	2	3	4	5	6	7	8	9	10
CL00F3	EID	MID	G1	G2	G3	P1	P2	P3	abc
CL00F3	1401	1	502	502	503	7			ABC
+bc	F1,G0	F2	F3	F					def
+BC	0.	1.	0.	1					DEF
+ef	A1	B1	A2	B2	A3	B3			
+EF	-0.2								

<u>Field</u>	<u>Contents</u>
EID	Element identification number (Integer > 0).
MID	Identification number of a MAT1 material card (Integer > 0).
G1, G2, G3	Grid point identification numbers of connection points (Integer > 0, G1 ≠ G2 ≠ G3)
P1, P2, P3	Identification numbers of PL00F3 cards defining section properties at G1, G2 and G3, respectively.
F1, F2, F3	Components of vector \vec{V} at the end of the beam corresponding to grid point G1. \vec{V} , together with the vector from G1 to G3, define the principal bending plane of the element. (Straight beams only).
G0	Grid point identification number to optionally define \vec{V} as the vector from G1 to G0.
F	Flag to specify the nature of fields 2, 3, 4 on the first continuation card as follows:

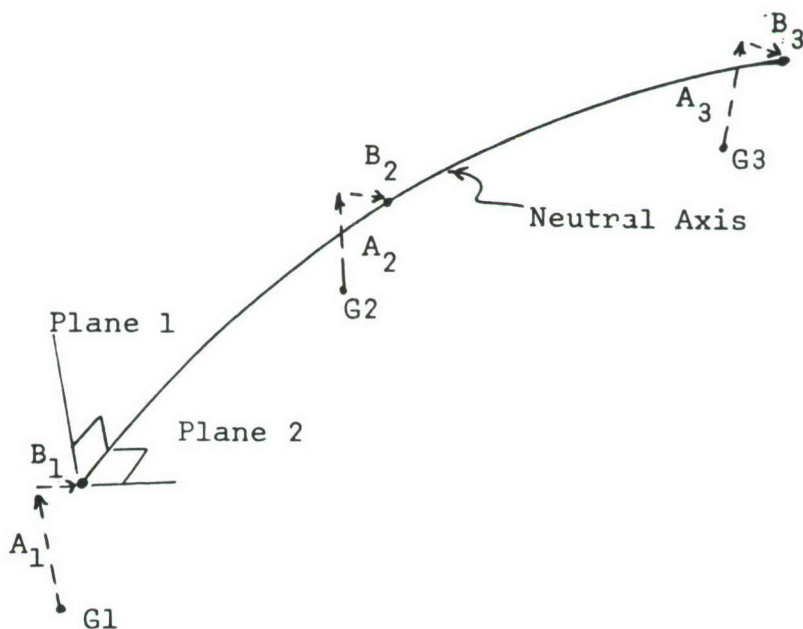
	2	3	4
F=1	F1	F2	F3
F=2	G0	blank	blank

BULK DATA DECK
CLOOF3 (Continued)

- A1, A2, A3 Offset distance of grid point G1, G2, or G3, respectively, in plane 1. Positive if the sense of A_i is the same as the sense of X₂ (see sketch below).
- B1, B2, B3 Offset distance of grid point G1, G2, or G3 respectively, in plane 2. Positive if the sense of B_i is the same as the sense of X₃ (see sketch below)



Straight Beam



Curved Beam

BULK DATA DECK
CL00F3 (Continued)

- Remarks:
1. Element identification numbers must be unique with respect to all other element identification numbers.
 2. The center node G2 should be located near the center of the arc connection G1 to G3.
 3. The continuation cards are optional.
 4. Degrees of freedom 4 and 5 at node G2 represent Loof rotations and are not the grid point rotations normally used by NASTRAN.
 5. The vector \vec{V} is used only for straight beams. For curved beams, the principal plane is always the plane in which the curved beam lies. The principal plane is plane 1, corresponding to I1, C1 and D1 on the PL00F3 card. Plane 2 corresponds to I2, C2 and D2.

BULK DATA DECK

Input Data Card CL00F6 Connections for L00F6 Element

Description: Defines a triangular Semi-Loof shell element of the structural model.

Format and Example:

1	2	3	4	5	6	7	8	9	10
CL00F6	EID	PID	G1	G2	G3	G4	G5	G6	abc
CL00F6	107	43	2	4	5	7	12	9	ABC
+bc	TH								
+BC	-14.0								

Field

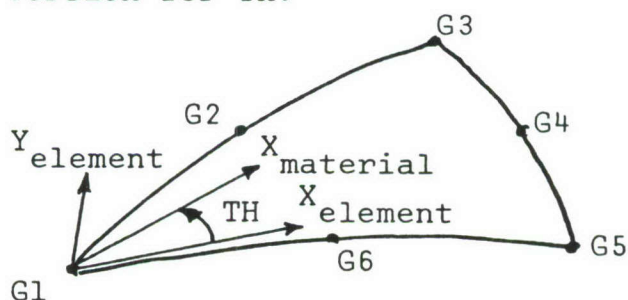
Contents

EID Element identification number (Integer > 0)

PID Identification number of a PL00F property card (Integer > 0)

G1,...,G6 Grid point numbers of connection points (Integer > 0, G1 ≠ G2 ≠ ... ≠ G6)

TH Material property orientation angle in degrees (Real). The sketch below gives the sign conversion for TH.



Remarks: 1. Element identification numbers must be unique with respect to all other element identification numbers.

2. Grid points G1 through G6 must be numbered consecutively around the perimeter of the element.

BULK DATA DECK
CL00F6 (Continued)

3. All interior angles must be less than 180° .
4. The Loof nodes G2, G4 and G6 should be located near the center of the arc connecting the respective corner points.
5. The continuation card is optional.
6. Degrees of freedom 4 and 5 at the Loof nodes represent Loof rotations and not the grid point rotations normally used by NASTRAN.

BULK DATA DECK

Input Data Card CL00F8 Connections for L00F8 Element

Description: Defines a quadrilateral Semi-Loof shell element of the structural model.

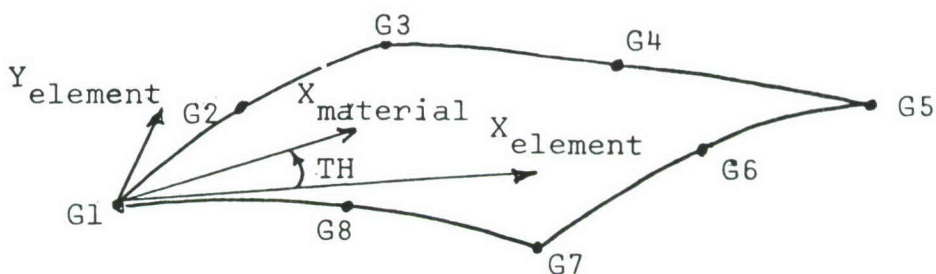
Format and Example:

1	2	3	4	5	6	7	8	9	10
CL00F8	EID	PID	G1	G2	G3	G4	G5	G6	abc
CL00F8	16	101	104	105	106	109	110	112	ABC
+bc	G7	G8	TH						
+BC	114	102							

Field

Contents

- EID Element identification number (Integer > 0)
- PID Identification number of a PL00F or PL00FX property card (Integer > 0)
- G1,...,G8 Grid point identification numbers of connection points (Integer > 0, $G1 \neq G2 \neq \dots \neq G8$)
- TH Material property orientation angle in degrees (Real).
The sketch below gives the sign conversion for TH.



- Remarks:
1. Element identification numbers must be unique with respect to all other element identification numbers.
 2. Grid points G1 through G8 must be numbered consecutively around the perimeter of the element.
 3. All interior angles must be less than 180° .

BULK DATA DECK
CLOOP8 (Continued)



4. Loof nodes G2, G4, G6 and G8 should lie near the mid point of the arc connecting the respective corner points.
5. Degrees of freedom 4 and 5 at the Loof nodes represent Loof rotations and not the grid point rotations normally used by NASTRAN.

BULK DATA DECK

Input Data Card MAT8 Material Property Definition, Form 8

Description: Defines the material properties for linear, temperature-independent, orthotropic materials.

Format and Example:

1	2	3	4	5	6	7	8	9	10
MAT8	MID	E11	E12	E22	G	RHO			
MAT8	1	10.2E6	3.2E6	8.2E6	4.1E6	0.2E-6			

Field

Contents

MID Material identification number (Integer > 0)

$\left. \begin{array}{l} E_{11} \\ E_{12} \\ E_{22} \\ G \end{array} \right\}$ Elements of stress-strain matrix

$$\left\{ \begin{array}{l} \sigma_1 \\ \sigma_2 \\ \tau_{12} \end{array} \right\} = \left[\begin{array}{ccc} E_{11} & E_{12} & 0 \\ E_{12} & E_{22} & 0 \\ 0 & 0 & G \end{array} \right] \left\{ \begin{array}{l} \epsilon_1 \\ \epsilon_2 \\ \gamma_{12} \end{array} \right\}$$

RHO Mass density

BULK DATA DECK

Input Data Card PLOAD4 Distributed Loads for L00F8 and L00F6 Elements

Description: Specifies distributed pressure loads by grid points. All L00F8 and L00F6 elements connected to these grids have these loads applied.

Format and Example:

1	2	3	4	5	6	7	8	9	10
PLOAD4	P	G1	G2	G3	G4	G5	G6	G7	
PLOAD4	-0.1	17	18	22					

Alternate Form:

PLOAD4	P	GA	"THRU"	GB					
PLOAD4	112.0	100	THRU	110					

<u>Field</u>	<u>Contents</u>
--------------	-----------------

P Pressure value (Real)

G1,G2,... } Grid point numbers (Integer > 0)
GA, GB }

- Remarks:
1. There is no set identification number. Distributed loads on L00F elements are not selected in the Case Control Deck.
 2. In the alternate form, all grid points from GA through GB are given the specified load.

BULK DATA DECK

Input Data Card PL00F Properties of L00F8 Element

Description: Defines the properties of a Semi-Loof shell element. Referenced by CL00F8 and CL00F6 cards.

Format and Example:

1	2	3	4	5	6	7	8	9	10
PL00F	PID	MID	THICK						
PL00F	2	1							

<u>Field</u>	<u>Contents</u>
PID	Property identification number (Integer > 0)
MID	Material identification number of a MAT1 or MAT8 card (Integer > 0)
THICK	Element thickness (Real)

Remarks: 1. If THICK is omitted, thicknesses will be taken from THLOOF cards.

BULK DATA DECK

Input Data Card PL00F3 Properties of a Semi-Loof Beam Element

Description: Defines the cross-sectional properties at one or more node points of a Semi-Loof curved beam element.

Format and Example:

1	2	3	4	5	6	7	8	9	10
PL00F3	PID	A	I1	I2	J	C1	C2	D1	abc
PL00F3	7	0.02	0.113	0.113	0.067	0.5	-0.5		+XY
+bc	D2								
+XY									

Field

Contents

PID Property identification number (Integer > 0)
A Cross-sectional area (Real > 0)
I1, I2 Moments of inertia (Real)
J Torsion constant (Real)
C1, C2, D1, D2 Stress recovery coefficients (Real)

Remarks:

1. The stress recovery coefficients are the local coordinates defining fiber distances for which stresses may be recovered.
2. The continuation card is optional.

BULK DATA DECK

Input Data Card PL00FX Orthotropic Membrane, Bending and Coupling Properties for L00F8 and L00F6 Elements

Description: Specifies a general force-strain law for a L00F8 or L00F6 element in matrix form.

Format and Example:

1	2	3	4	5	6	7	8	9	10
PL00FX	PID	A11	A12	A13	A22	A23	A33	X	abc
PL00FX		10.2E6	4.1E6		7.6E6		3.3E6		ABC
+bc	DENS	B11	B12	B13	B22	B23	B33	X	def
+BC	0.06								DEF
+ef	X	D11	D12	D13	D22	D23	D33	X	
DEF		5.2E6	0.9E6		3.2E6		1.8E6		

Field

Contents

PID Property identification number (Integer > 0)

$\left. \begin{matrix} A_{ij} \\ B_{ij} \\ D_{ij} \end{matrix} \right\}$ Matrices defining the force-strain law for the element as follows:

$$\begin{Bmatrix} N_1 \\ N_2 \\ N_{12} \\ \hline M_1 \\ M_2 \\ M_{12} \end{Bmatrix} = \begin{bmatrix} A_{11} & A_{12} & A_{13} & B_{11} & B_{12} & B_{13} \\ A_{12} & A_{22} & A_{23} & B_{12} & B_{22} & B_{23} \\ A_{13} & A_{23} & A_{33} & B_{13} & B_{23} & B_{33} \\ \hline B_{11} & B_{12} & B_{13} & D_{11} & D_{12} & D_{13} \\ B_{12} & B_{22} & B_{23} & D_{12} & D_{22} & D_{23} \\ B_{13} & B_{23} & B_{33} & D_{13} & D_{23} & D_{33} \end{bmatrix} \begin{Bmatrix} \epsilon_1 \\ \epsilon_2 \\ \gamma_{12} \\ \hline \chi_1 \\ \chi_2 \\ \chi_{12} \end{Bmatrix}$$

BULK DATA DECK

PL00FX (continued)

N_1 , N_2 , and N_{12} are in-plane force resultants and G_1 , G_2 , and γ_{12} are corresponding in-plane strains. M_1 , M_2 , and M_{12} are moment resultants, and χ_1 , χ_2 , and χ_{12} are corresponding curvatures.

DENS Effective mass per unit area for the shell

- Remarks:
1. No material property card is referenced since material properties are implicit in the matrices A, B, D, and DENS.
 2. The continuation cards may be omitted if no bending or coupling terms are desired.

BULK DATA DECK

Input Data Card THLOOF Grid-point Thicknesses

Description: Associates thicknesses with grid points for use by Semi-Loof shell elements.

Format and Example:

1	2	3	4	5	6	7	8	9	10
THLOOF	TH	GID	GID	GID	GID	GID	GID	GID	
THLOOF	0.07		4	7	6				

Alternate form:

THLOOF	TH	GID1	"THRU"	GID2					
THLOOF	0.2	16	THRU	48					

<u>Field</u>	<u>Contents</u>
TH	Thickness (Real > 0)
GID } GID1 } GID2 }	Grid-point identification number (Integer > 0, GID1 < GID2)

Remarks:

1. GID must be 0 or blank for omitted entries.
2. At least one positive GID must be present on each THLOOF card.

A.4 POSTLOOF INPUT

The POSTLOOF code serves to recover two kinds of information from solution vectors generated by NASTRAN: rotations and stresses. The solution vectors may be static displacements, eigenvectors, transient or frequency response vectors. If the NASTRAN output file contains more than one solution vector, information is recovered and printed successively for each vector.

The reason for recovering rotations in POSTLOOF is that node point rotations are not independent degrees of freedom with Semi-Loof shell elements, yet they may be of interest. Furthermore, rotations at arbitrary locations may be of interest, and these may be recovered as well. At node points, rotations are discontinuous between elements. The approach used for node points is to average the values indicated by all the elements that touch a given node point, except that when a LOOF3 beam element defines a rotation at a grid point, that value is used.

Stresses, if requested, are output at Gauss points only, which are optimal for stress calculations.

Input to POSTLOOF is as follows with default options underlined:

\$NASTRAN	$\left\{ \begin{array}{l} \text{MSC} \\ \text{COSMIC} \end{array} \right\}$	Default set for each installation.
\$LOOFROT	$\left\{ \begin{array}{l} \text{NONE} \\ \text{GRIDS} \\ n \end{array} \right\}$	Where "NONE" suppresses recovery of rotations, "GRIDS" causes recovery at all grid points, and n is an integer indicating that rotations are requested at n specific locations, listed on n data cards immediately following.

Rotation recovery locations: (LOOFROT n option)

Col.			
1-5	6-15	16-25	where m_i is an element number, and ξ_i, η_i
m_1	ξ_1	η_1	are local coordinates for that element
m_2	ξ_2	η_2	$(-1 \leq \xi_i, \eta_i \leq 1 \text{ for LOOF8, } 0 \leq \xi_i, \eta_i \leq 1$
.	.	.	for LOOF6). See Figure A-7 for local
.	.	.	coordinates.
m_n	ξ_n	η_n	

\$LOOFSTRESS $\left\{ \begin{array}{c} \text{NONE} \\ \text{ALL} \\ \text{R} \end{array} \right\}$ Where "NONE" suppresses recovery of stresses, "ALL" causes recovery for all elements, and n is an integer indicating that stresses are desired for n specific elements, listed on following data cards.

Stress recovery elements: (\$LOOFSTRESS n option)

Col.

1-5

Where m_i is a LOOF8 or LOOF6 element number.

m_1

m_2

.

.

.

m_n

POSTLOOF can easily be tailored by the user-programmer for specific requirements. A POSTLOOF listing appears in Appendix B.

A.5 DECK SETUP

A complete Semi-Loof analysis consists of three major steps: PRELOOF, NASTRAN, and POSTLOOF. In order to relieve the user of the burden of operating system details, control card procedures ("procs") have been set up so that each step can be called by a single control card:

```
BEGIN(PRE,LOOF)
NASTRAN(DATA)ATTACH
BEGIN(POST,LOOF)
```

Proc "PRE" uses the UPDATE utility to fetch the DMAP alter package required by the user and insert it into the user's deck. PRELOOF is executed, and the NASTRAN program file is attached. Proc "POST" reads input and executes POSTLOOF.

DMAP alter packages are maintained in an UPDATE program library, which is automatically accessed by PRE. There are two alter packages for each rigid format (currently RF 1, 3, 11, 24, and 25 are supported): one for preprocessing and one for postprocessing. Deck names for each DMAP sequence are formed by concatenating the letters RF, the rigid format number, and either PRE or POST, e.g. RF3PRE or RF25POST. Each run will use an RFxxPRE sequence, an RFxxPOST sequence, or both. They are listed on an UPDATE *C card followed by an end-of-record card (7/8/9).

Finally, PRE attaches the NASTRAN program file. Thus, a complete deck setup for a run including all three processing steps would have the following form (assume rigid format 3 is being run):

<u>Cards(s)</u>	<u>Remarks</u>
job card, account cards ATTACH(LOOF,ID=xxxx)	Follow installation rules. Attach proc file using ID specified for a particular installation.
BEGIN(PRE,LOOF) RFL(xxxxxx)	Assemble input and execute PRELOOF. Request appropriate field length for NASTRAN.
NASTRAN(DATA)ATTACH BEGIN(POST,LOOF) 7/8/9	Execute NASTRAN. Execute POSTLOOF.
*C RF3PRE,RF3POST	Fetch DMAP alter sequences.

<u>Card(s)</u>	<u>Remarks</u>
7/8/9 *DECK X . . .	PRELOOF input as explained in Section A.3.
7/8/9 . . .	POSTLOOF input as explained in Section A.4.
6/7/8/9	

Many users like to maintain large data decks in UPDATE program libraries. To accommodate this feature, PRE checks for the presence of a file named OLDPL and if one is present, it makes corrections to decks on that file using correction cards supplied by the user. In this case, a card such as

```
ATTACH(OLDPL,MYSEMILOOFDECK,ID=xxxx)
```

would precede BEGIN(PRE,LOOF) and the input record for PRELOOF would consist of UPDATE corrections such as

```
*ID FIXMYDECK
*I EXEC.10
*READ RFALTS
*D BULK.200
PARAM WTMASS .00259
7/8/9
```

For large problems, the user may desire to rerun NASTRAN without rerunning PRELOOF, or to rerun POSTLOOF without rerunning NASTRAN. To do this, the user must know what files are passed from one step to the next. Four files are involved:

<u>File</u>	<u>Remarks</u>
UT1	Contains stiffness matrices, mass matrices, and load vectors. Written by PRELOOF and read by NASTRAN. Binary format.
DATA	NASTRAN data deck after modification by PRELOOF. Written by PRELOOF and read by NASTRAN. Card image format.
LSTRES	Stress matrices and element connection tables. Written by PRELOOF and read by POSTLOOF. Binary format.

<u>File</u>	<u>Remarks</u>
UT2	Solution vector(s). Written by NASTRAN and read by POSTLOOF. Binary format.

Figure A-3 shows how these files pass information from PRELOOF to NASTRAN to POSTLOOF.

In running restart cases one must know what information must be regenerated and what information can be salvaged from previous runs. If any structural data changes, PRELOOF must be rerun and UT1 and DATA regenerated. Proc PRE checks for the presence of UT1 and if present, skips PRELOOF execution and just does UPDATE processing and attaches NASTRAN. If the NASTRAN restart facility is to be used, involving problem tapes and checkpoint dictionaries, one needs to know whether the rigid format will be restarted before or after the point at which UT1 is read in. If before, UT1 must either be generated or recovered from a previous run before executing NASTRAN. If after, UT1 is omitted, and the DMAP sequence used to read UT1 (i.e. RFxxPRE) is also omitted.

Now consider a sequence of runs for a large problem. Assume that input data is maintained on an UPDATE program library with deck names EXEC, CASE, and BULK, and that the following sequence of runs is executed:

1. Check out data and plot the undeformed structure.
2. Run static load check case.
3. Run dynamic case using NASTRAN checkpoint facility.
4. Restart NASTRAN without rerunning PRELOOF, save output from NASTRAN, then run POSTLOOF.
5. Rerun POSTLOOF.

The deck setups would be as follows:

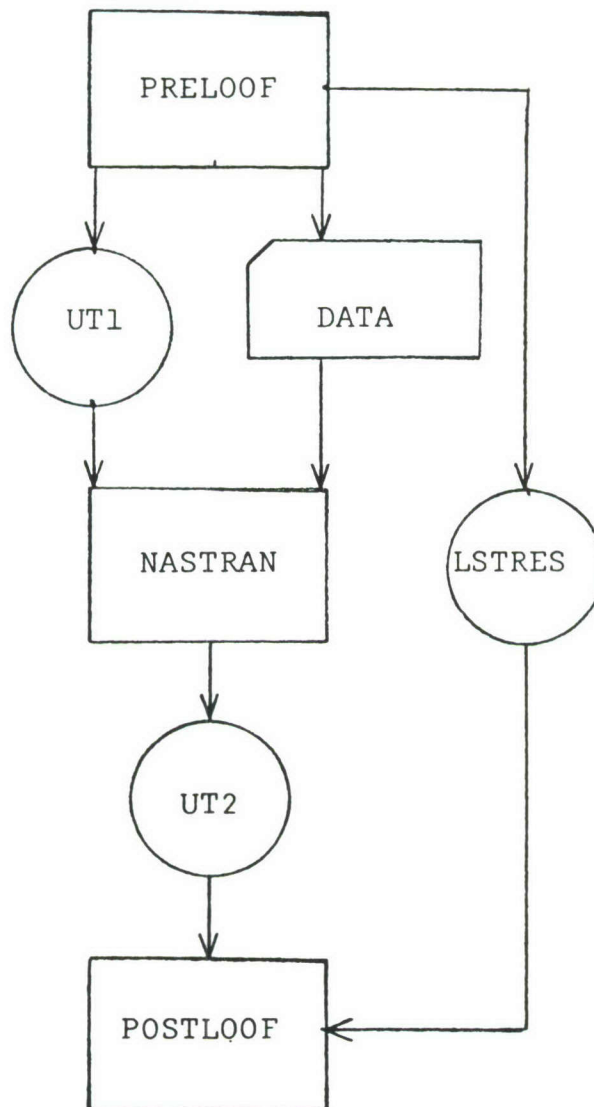


Figure A-3 Flow of Control from PRELOOF to NASTRAN to POSTLOOF

Run 1: Data Check and Plot

Card(s)

Remarks

```

job card, account card
ATTACH(LOOF,ID=xxxx)
ATTACH(OLDPL,BIGPROBLEM,ID=...)
BEGIN(PRE,LOOF)
RFL(160000)
NASTRAN(DATA)ATTACH
(plot control cards)
7/8/9

```

Null input record for UPDATE.
No DMAP sequences used since
execution will halt right after
plot generation.

7/8/9

```
*ID FIXES
*I EXEC.10
$LOOFMAT NO
$LOOFPLOT YES
ALTER 23
EXIT$
*I CASE.10
OUTPUT(PLOT)
```

```
Suppress matrix generation.
Generate PLOTEL cards.
Stop rigid format after plots.
```

Plot commands.

*C EXEC,CASE,BULK
6/7/8/9

Run 2: Static Test Case (Rigid Format 1)

```

job card, account card
ATTACH(LOOF,ID=xxxx)
ATTACH(OLDPL,BIGPROBLEM,ID=...)
BEGIN(PRE,LOOF)
RFL(160000)
NASTRAN(DATA)ATTACH
7/8/9
*C RFLPRE

```

Get DMAP sequence for rigid format
1 preprocessing.

7/8/9

```
*ID FIXES
*I EXEC.10
*READ RFALTS
```

Get DMAP sequence.
Other changes as needed.

```

.
.
.
*C EXEC,CASE,BULK
6/7/8/9

```

Run 3: Dynamic Run

Save PRELOOF output on tape, NASTRAN NPTP on tape, and checkpoint dictionary on permanent file.

<u>Card(s)</u>	<u>Remarks</u>
job card, account card	
ATTACH(LOOF,ID=xxxx)	
ATTACH(OLDPL,BIGPROBLEM,ID=...)	
BEGIN(PRE,LOOF)	
LABEL(SAVE,L=SAVEPRELOOF,VSN=xxx)	Tape to save UT1.
COPY(UT1,SAVE)	Save UT1 and LSTRES.
COPY(LSTRES,SAVE)	
REWIND(UT1)	
RETURN(LSTRES,SAVE)	
REQUEST(CHKPNT,*PF)	Permanent file for checkpoint dictionary.
LABEL(NPTP,W,L=NPTP,VSN=xxx)	New problem tape.
RFL(240000)	
NASTRAN(DATA,,CHKPNT)ATTACH	Punch output (checkpoint dictionary) routed to file CHKPNT.
CATALOG(CHKPNT,MYDICT,ID=xxxx)	
7/8/9	
*C RF3PRE,RF3POST	
7/8/9	
*ID FIXES	
*D EXEC.6	
SOL 3.0	Switch to rigid format 3.
CHKPNT YES	
*I EXEC.10	
*READ RFALTS	
.	Other changes as needed.
.	
.	
*C EXEC,CASE,BULK	
6/7/8/9	

Run 4: Restart NASTRAN, Execute POSTLOOF

<u>Card(s)</u>	<u>Remarks</u>
job card, account card	
ATTACH(RESTART,MYDICT,ID=xxxx)	Get checkpoint dictionary and old problem tape.
LABEL(OTPT,R,L,VSN=xxx)	
ATTACH(OLDPL,BIGPROBLEM,ID=xxx)	
UPDATE(Q,D,8,C=DATA,L=0)	Use UPDATE to set up DATA file.
LABEL(SAVE,W,L=SAVEUT2,VSN=xxx)	Tape to save file UT2.
COPY(UT2,SAVE)	

<u>Card(s)</u>	<u>Remarks</u>
REWIND(UT2)	
RETURN(SAVE)	
BEGIN(POST,LOOF)	
7/8/9	
*C RF3POST	
7/8/9	
*IC FIXES	
*D EXEC.6	
SOL 3,0	
*READ RESTART	Insert checkpoint dictionary.
*I EXEC.10	
*READ RFALTS	
.	Other changes as needed.
.	
.	
*I CASE.15	
BEGIN BULK	Do not retrieve bulk data from OLDPL. Insert bulk data additions and corrections after Case Control Deck.
.	
.	
.	
ENDDATA	
7/8/9	
\$LOOFROT ALL	POSTLOOF input.
.	
.	
.	
6/7/8/9	
<u>Run 5: Rerun POSTLOOF</u>	
job card, account card	
LABEL(SAVE1,R,L=SAVEPRELOOF,VSN=xxx)	
SKIPF(SAVE1,1,17)	Skip UT1.
COPY(SAVE1,LSTRES)	Recover LSTRES.
RETURN(SAVE)	
REWIND(LSTRES)	
LABEL(SAVE2,R,L=SAVEUT2,VSN=xxx)	Recover UT2.
COPY(SAVE2,UT2)	
RETURN(SAVE2)	
REWIND(UT2)	
BEGIN(POST,LOOF)	
7/8/9	
\$LOOFROT ALL	POSTLOOF input.
.	
.	
.	
6/7/8/9	

A.6 PROGRAMMER'S NOTES

This section is a brief description of the programming techniques used in PRELOOF, a list of subroutines, and a list of important variables. A listing of the source codes for both PRELOOF and POSTLOOF can be found in Appendix B.

Because the amount of main storage required varies so much from one application to another, a dynamic core allocation procedure is used. Blank common space is added as needed by calls to subroutine MORCOR. The address of the newly acquired core relative to the beginning of blank common is available in variable NCORE, and this address is passed to a subroutine. For example, to allocate a new array C of size N by M, one would code the following:

```
COMMON A(1)
.
.
.
CALL MORCOR(N*M)
CALL SUBR(A(NCORE),N,M)
.
.
.
SUBROUTINE SUBR(C,N,M)
DIMENSION C(N,M)
```

Examples of this technique can be seen throughout PRELOOF.

PRELOOF checks input for errors. When a data card format error is detected, subroutine BUM notes the card which was erroneous. When any error is detected, a flag is set so that the job can be terminated after all input has been read.

Since grid point numbers, element numbers, etc., are arbitrary, a search procedure is needed so that the array entry corresponding to a given grid point or element can be located when the grid point number or element number is given. Function INDEX does this search, and returns a -1 if a match is not found, so that an error can be noted.

Since element matrices are accessed in random sequence during the assembly process, it is convenient to store these

matrices on random disk files. Element stiffness matrices are stored on unit 98, mass matrices on 99, and load vectors on 97. Utility subroutines READMS and WRITMS are used to access these files. This approach also allows two elements to share the same matrices if they are congruent.

The sequence of operations of PRELOOF is as follows:

The main program first calls INIT, which sets default values and reads \$LOOF option cards. It then calls INPUT to read the bulk data deck.

INPUT calls FETCH once for each bulk data card. FETCH puts the alphanumeric card image in common block CARD, along with indicators to show whether that card is to be processed by PRELOOF. If so INPUT branches to the appropriate subroutine to decode information from the card images. These routines are:

- CORD for CORD2R, CORD2C and CORD2S cards
- GRID for GRID cards
- SEQGP for SEQGP cards
- ISHELL for CLOOF8 and CLOOF6 cards
- IBEAM for CLOOF3 cards
- CONGR for CONGR cards
- IPSHEL for PLOOF cards
- IPSHELX for PLOOFX cards
- IPBEAM for PLOOF3 cards
- IMAT for MAT1 and MAT8 cards
- IPLOAD for PLOAD4 cards
- ITHICK for THLOOF cards
- PARAM for PARAM cards
- ISPNT for SPOINT cards

All these routines use subroutine DCODE to decode 8-character fields in either integer, floating point, or alphabetic format, left or right adjusted. DCODE2 is available for 16-character fields.

PRELOOF then calls PROCESS to do some additional processing not handled by the input routines. First grid points are sequenced (unless SEQGP cards were present), using subroutine SORT. Then all shell elements are scanned, corner nodes are distinguished from mid-side nodes, and any attempts to use a node as a mid-side node of one element and a corner node of

another are noted. Appropriate permanent single-point constraints are generated. PLOTTEL cards are generated if requested. Beam elements are then scanned in the same manner. Modified GRID cards are written out to file DATA.

PREL00F then calls XSHELL, which controls shell element generation. XSHELL loops through all elements, gathering connection and property information, then calling QSHELL once for each element. After QSHELL has generated matrices, XSHELL stores them on disk and prints them out, if requested.

QSHELL loops through the Gauss points required for integration, calling ZHELL once for each Gauss point.

ZHELL calls HAL00F to generate shape functions and their derivatives, assembles the required functions into matrix B, assembles stiffnesses into matrix D, and computes the matrix product B^TDB , which is then accumulated onto the stiffness matrix. The mass matrix and load vector are handled in a similar fashion.

HAL00F is really the heart of the code. It generates values of all the shape functions and their required derivatives in array WSHEL, direction cosines in matrix FRAM, an area integrating factor in AREA, and a boundary integrating factor in SIDE. HAL00F calls SFR to obtain actual shape function values, then manipulates these to reduce out unwanted degrees of freedom. This subroutine was supplied by Prof. Irons, and more information about it may be found in his publications [2] and [32].

For beam elements, PREL00F calls XBEAM, which controls beam element matrix generation. It calls ZBEAM once for each element. ZBEAM calls LOFBEM once for each integrating point, and LOFBEM in turn calls SFR1 to obtain shape function values. LOFBEM and SFR1 perform the same functions for beams as HAL00F and SFR do for shells.

Having generated element matrices, PREL00F assembles them into global matrices. This is done by calls to ASSY, once for the stiffness matrix and once for the mass matrix, if requested.

ASSY generates six columns at a time and writes them out in a format compatible with NASTRAN (INPUTT2 module for COSMIC NASTRAN, INPUTT4 for MSC/NASTRAN).

If element load vectors have been computed, they are assembled into a master load vector by ASSL which writes out a single column.

The following is a list of the major variables used in PRELOOF. All array values are floating point.

<u>Variable</u>	<u>Significance</u>
NCORD	Number of coordinate systems
NUMNP	Number of node points
NSHELL	Number of L00F8 and L00F6 elements
NPSHEL	Number of PL00F shell property cards
NPSHELX	Number of PL00FX shell property cards
NSHLTY	Number of distinct shell element matrices
NBEAM	Number of beam elements
NPBEAM	Number of PL00F3 beam property cards
NBMTY	Number of distinct beam element matrices
NMAT	Number of material properties
NSPOINT	Number of SPOINT's
CORD(3,5,NCORD)	Coordinate system data (1,1) External id. (2,1) Type = 1 = rectangular 2 = cylindrical 3 = spherical (i,2) Origin (i,j+2) Transformation matrix
GRID(13,NUMNP)	Grid point data (1) External id. (2) Location coordinate system id. (3 to 5) Location (location coordinate system) (6 to 8) Location (basic coordinate system) (9) Displacement coordinate system id. (10) Permanent single-point constraints (11) Sequence number (12) Pressure (13) Thickness
SEQ(NUMNP)	Same as GRID (11)

<u>Variable</u>	<u>Significance</u>
SHELL(12,NSHELL)	Shell connections (1) External id. (2) Property id. (3-10) Grid point numbers (external id. initially, sequence numbers later) (11) Element matrix number (12) Material orientation angle for orthotropic materials
PSHELL(3,NPSHEL)	Isotropic shell properties (1) External id. (2) Material id. (3) Thickness
PSHELX(20,NPSHELX)	Anisotropic shell properties (1) External id. (2-7) A matrix (8-13) B matrix (14-19) D matrix (20) Area density
BEAM(21,NBEAM)	Beam connections (1) External id. (2) Material id. (3-5) Grid points (external id. initially, sequence numbers later) (6-8) Property id. for cross-sections at the three grid points (9-14) Offsets (15) Element matrix number (16-18) Not used (19-21) Orientation vector
PBEAM(9)	Beam cross-section properties (1) External id. (2) Area (3,4) Moments of inertia (5) Torsion constant (6-9) Stress recovery coefficients
XMAT(6)	Material properties (1) External id. (2) E_{11} (3) E_{12} (4) E_{22} (5) G^{22} (6) ρ

APPENDIX B

SEMI-LOOF PROGRAM LISTINGS

B.1 CONTROL CARD PROCEDURES

```
.PROC,UPDATE.
ATTACH(LUOFPL,ID=D740292,MR=1)
UPDATE(Q,P=LUOFPL)
FTN(I=COMPILE,B=NEW)
ATTACH(OLD,LUOFGU,ID=D740292,MR=1)
REWIND(OLD,NEW)
COPYL(OLD,NEW,LUOFGU)
RETURN(LUOFPL,COMPILE,OLD,NEW)
REVERT.
X
.PROC,PRE.
COMMENT. PRELOUF,WPAFB
IFE,FILE(UT1,.NOT.(LO.OR.PF)),SKIPB.
ATTACH(RFALTER,LOUFALTER,ID=D740292,MR=1)
UPDATE(P=RFALTER,L=1,C=RFALTS,D,8)
RETURN(RFALTER)
UPDATE(D,8,C=PRE,L=1)
RETURN(RFALTS)
IFE,FILE(LUOFGU,.NOT.(LO.OR.PF)),SKIPB.
ATTACH(LUOFGU,ID=D740292,MR=1)
ENDIF(SKIPB)
LDSET(MAP=B/MAPFILE)
LUOFGU(PRE)
RETURN(LUOFGU,PRE,TAPE97,TAPE98,TAPE99)
ENDIF(SKIPB)
ATTACH(NASTRAN,ID=NASTRAN,MR=1)
LIBRARY(FORTRAN,SYSIU)
MAP(UFF)
RETURN(MAPFILE)
REVERT.
EXIT.
REWIND(MAPFILE)
COPY(MAPFILE,OUTPUT)
REVERT(ABORT)
X
.PROC,PLOT.
REDUCE.
ATTACH(M,MANDMCALCOMP,ID=D740292)
ATTACH(NASPLT,ID=D740292)
MAP(UFF)
REWIND(PLT2)
ROUTE(ANAPLT,DEF,DC=PR,ST=CSA,TID=TB)
LDSET(LIB=M)
NASPLT.
RETURN(ANAPLT,NASPLT,M)
REVERT.
X
```

```

.PROC,RESTORE,PROB=L740292.
LABEL(SAVE,R,L=LOOF,VSIN=L04454)    REMOTE REQUEST, OWNER=D730368
REQUEST(LOOF,*PF)
COPYBF(SAVE,LOOF)
CATALOG(LOOF,ID=PROB,RP=999)
RETURN(LOOF)
REQUEST(LOOFGU,*PF)
COPYBF(SAVE,LOOFGU)
CATALOG(LOOFGU,ID=PROB,RP=999)
RETURN(LOOFGU)
COPYBF(SAVE,OLDPL)
REQUEST(NEWPL,*PF)
UPDATE(N,F,C=0)
CATALOG(NEWPL,LOOFALTER,ID=PROB,RP=999)
RETURN(OLDPL,NEWPL)
COPYBF(SAVE,OLDPL)
RETURN(SAVE)
UPDATE(N,C=0)
REQUEST(NEWPL,*PF)
CATALOG(NEWPL,LOOFPL,ID=PROB,RP=999)
RETURN(OLDPL,NEWPL)
X
W

```


B.2

DMAP ALTER LIBRARY
(Partial Listing)

```

$ *** ALTER MSC/NASTRAN, RIGID FORMAT 24 ***
$ *** TO ACCEPT FILES FROM PRELOOF FOR SEMI-LOOF ELEMENTS ***
$
$     PARAMS USED BY THIS ALTER ...
$
$     NAME          DEFAULT    REMARKS
$     ----          -
$
$     NSPOINT      0          NUMBER OF SPOINT-S IN THE MODEL
$
$     STIFFMAT     -1          POSITIVE VALUE INDICATES THAT PRELOOF
$                               GENERATED A STIFFNESS MATRIX
$
$     MASSMAT      -1          POSITIVE VALUE INDICATES THAT PRELOOF
$                               HAS GENERATED A MASS MATRIX.
$
$     LOADVEC      -1          POSITIVE VALUE INDICATES THAT PRELOOF
$                               GENERATED A LOAD VECTOR CORRESPONDING
$                               TO PRESSURE LOADING FOR THIS MODEL.
$
$     *NOTE* PRELOOF GENERATES ALL THESE PARAMS AUTOMATICALLY
$
$ ALTER 93
$ PARAM //SUB/V,N,NOEXTRA/V,Y,NEXTRA=0/1 $
$ PARAM //NOP/V,Y,STIFFMAT=-1 $
$ COND LBLNOK,STIFFMAT $
$ INPUTT4 /KLGG,,,,/1/11 $
$ LABEL LBLNOK $
$ EQUIV KLGG,KGG/NOEXTRA $
$ COND LBLNOX,NOEXTRA $
$ PARAM //MPY/V,N,NDOF/V,Y,NEXTRA/6 $
$ PARAM //ADD/V,N,MDOF/V,N,NDOF/V,Y,NSPOINT=0 $
$ PARAM //SUB/V,N,LUSET2/V,N,LUSET/V,N,MDOF $
$ MATGEN ,/RP/6/V,N,LUSET/V,N,LUSET2/V,N,MDOF $
$ MERGE KLGG,,,,RP,/KLGG1/ $
$ ADD KLGG1,KGG/KGG1 $
$ EQUIV KGG1,KGG/ALWAYS $
$ LABEL LBLNOX $
$ PARAM //NOP/V,Y,MASSMAT=-1 $
$ COND LBLNUM,MASSMAT $
$ INPUTT4 /MLGG,,,,/1/11/0 $
$ EQUIV MLGG,MGG/NOEXTRA $
$ COND LBLNUM,NOEXTRA $
$ MERGE MLGG,,,,RP,/MLGG1/ $
$ ADD MLGG1,MGG/MGG1 $
$ EQUIV MLGG1,MGG/ALWAYS $
$ LABEL LBLNUM $

```

```

ALTER 145
PARAM      //NOP/V,Y,LOADVEC=-1 $
COND       LBLNOPGL,LOADVEC $
INPUTT4    /PGL,,,,/1/11/0 $
EQUIV      PGL,PGLL/NOEXTRA $
COND       LBLNOP,NOEXTRA $
MERGE      PGL,,,,,RP/PGLL/0 $
LABEL      LBLNOP $
ADD        PGLL,PG/PGW $
EQUIV      PGW,PG/ALWAYS $
LABEL      LBLNOPGL $

```

```

$ *** ALTER MSC/NASTRAN, RIGID FORMAT 24 ***
$ *** TO PRODUCE FILES FOR POSTLOOF FOR SEMI-LOOF ELEMENTS ***
ALTER 160
OUTPUT4    UGV,,,,//0/12 $

```

```

$ *** ALTER MSC/HASIRAN, RIGID FORMAT 25 ***
$ *** TO ACCEPT FILES FROM PRELOOF FOR SEMI-LOOF ELEMENTS ***
$
$      PARAMS USED BY THIS ALTER ...
$
$      NAME          DEFAULT    REMARKS
$      ----          -
$
$      NSPOINT      0          NUMBER OF SPOINT-S IN THE MODEL
$
$      STIFFMAT     -1          POSITIVE VALUE INDICATES THAT PRELOOF
$                               GENERATED A STIFFNESS MATRIX
$
$      MASSMAT      -1          POSITIVE VALUE INDICATES THAT PRELOOF
$
$      *NOTE* PRELOOF GENERATES ALL THESE PARAMS AUTOMATICALLY
$
ALTER 76,80
PARAM      //AND/V,N,NOM/NOMGG/V,Y,MASSMAT=-1 $
COND       RFERR,NOM $
ALTER 83,83
PARAM      //AND/V,N,NOK/NOKGG/V,Y,STIFFMAT=-1 $
ALTER 92
PARAM      //SUB/V,N,NOBLOW/V,Y,NSPOINT=0/1 $
PARAM      //SUB/V,N,LUSET2/LUSET/NSPOINT $
COND       LNOBLOW,NOBLOW $
MATGEN     ,/BLOW/6/LUSET/LUSET2/NSPOINT $
LABEL      LNOBLOW $
COND       NOLOOFK,STIFFMAT $
INPUTT4    /KGG1,,,/1/11 $
EQUIV      KGG1,KGG2/NOBLOW $
COND       NOLOOFK,NOBLOW $
MERGE      KGG1,,,BLOW,/KGG2/ $
LABEL      NOLOOFK $
ADD        KGG2,KGG/KGGL $
EQUIV      KGGL,KGG/ALWAYS $
COND       NOLOUFM,MASSMAT $
INPUTT4    /MGG1,,,/1/11/0 $
EQUIV      MGG1,MGG2/NOBLOW $
COND       NOLOUFM,NOBLOW $
MERGE      MGG1,,,BLOW,/MGG2/ $
LABEL      NOLOUFM $
ADD        MGG2,MGG/MGGL $
EQUIV      MGGL,MGG/ALWAYS $
COND       LGPWG,GRDPNT $
GPWG       BGPDT,CSTM,EGEXIN,MGG/UGPWG/V,Y,GRDPNT=-1/C,Y,WTMASS $
OFF        UGPWG// $
LABEL      LGPWG $

```


B.3 PRELOOF PROGRAM LISTING

```

PROGRAM PRELOOF (INPUT=201B,DATA=201B,OUTPUT=201B,UT1,LSTRES,
.      TAPE3=DATA,TAPE4=UT1,TAPE7=LSTRES,
.      TAPE1=INPUT,TAPE5=INPUT,TAPE6=OUTPUT,
.      TAPE97,TAPE98,TAPE99)

C
C   MATRIX GENERATION PREPROCESSOR FOR SEMI-LOOP
C   SHELL AND BEAM ELEMENTS FOR USE IN NASTRAN
C   W.G. 12-78
C
COMMON/CONTRL/ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD,NCORE,
.      WTMASS,NASTY
LOGICAL ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD
COMMON/LOC/LCORD,LGRID,LSEQ,LSHELL,LBEAM,LPSHEL,LPSHELX,LPBEAM,
.      LMAT,LPSPC
COMMON/N/NCORD,NUMNP,NSHELL,NPSHEL,NPSHELX,NSHLTY,NBEAM,NPBEAM,
.      NBMTY,NMAT,NSPOINT,MAXGRD,NDOF
COMMON A(1)
CALL INIT

C
C   READ STUFF IN; DO PARTIAL ERROR CHECK
C
CALL INPUT
IF (ERR) CALL SYSTEM(52,32H JOB ABORTED DUE TO ABOVE ERRORS)

C
C   DO SOME MORE PROCESSING
C
CALL PROCESS(A(LGRID),A(LSHELL),A(LBEAM),A(LSEQ),A(LPSPC))
IF (ERR) CALL SYSTEM(52,32H JOB ABORTED DUE TO ABOVE ERRORS)
REWIND 3

C
C   GET MATRICES, IF WANTED
C
IF (,.NOT.MATGEN) STOP

C
C   RANDOM FILE INDICES
C
NL=NSHLTY+NBMTY+1
LSINDEX=NCORE
CALL MORCOR(NL)
LMINDEX=NCORE
CALL MORCOR(NL)
LLINDEX=NCORE
CALL MORCOR(NL)
CALL OPENMS(98,A(LSINDEX),NL,0)
IF (DYN) CALL OPENMS(99,A(LMINDEX),NL,0)
IF (APPLD) CALL OPENMS(97,A(LLINDEX),NL,0)

C
C   SHELL ELEMENTS
C
LELSTIF=NCORE

```

```

CALL MURCOR(32*32)
LELMASS=NCORE
CALL MURCOR(32*32)
LELOAD=NCORE
CALL MURCOR(32)
LELSTRS=NCORE
CALL MURCOR(6*32)
LSHAPE=NCORE
CALL MURCOR(13*45)
CALL XSHELL(A(LGRID),A(LSHELL),A(LPSHEL),A(LPSHELX),A(LMAT),
.      A(LELSTIF),A(LELMASS),A(LELOAD),A(LELSTRS),
.      A(LSINDEX),A(LSHAPE),A(LSEQ))
NCORE=LELSTIF
C
C BEAM ELEMENTS
C
LELSTIF=NCORE
CALL MURCOR(17*17)
LELMASS=NCORE
CALL MURCOR(17*17)
LELOAD=NCORE
CALL MURCOR(17)
LELSTFS=NCORE
CALL MURCOR(17*6)
LSHAPE=NCORE
CALL MURCOR(13*45)
CALL XBEAM(A(LGRID),A(LBEAM),A(LPBEAM),A(LMAT),
.      A(LELSTIF),A(LELMASS),A(LELOAD),A(LELSTRS),
.      A(LSINDEX),A(LSHAPE),A(LSEQ))
IF (ERR) CALL SYSTEM(52,32H JOB ABORTED DUE TO ABOVE ERRORS)
NCORE=LELSTIF
NDOF=6*NUMNP+NSPOINT
CALL MURCOR(32*32)
LELMAT=NCORE
CALL MURCOR(32*32)
LBIG=NCORE
CALL MURCOR(6*NDOF)
CALL ASSY(4HSTIF,98,
.      A(LBIG),A(LSHELL),A(LBEAM),A(LCORD),A(LGRID),A(LSEQ),
.      A(LPSPC),A(LELMAT),A(LELMAT),NDOF)
IF (ERR) CALL SYSTEM(52,32H JOB ABORTED DUE TO ABOVE ERRORS)
IF (DYN)
. CALL ASSY(4HMASS,99,
.      A(LBIG),A(LSHELL),A(LBEAM),A(LCORD),A(LGRID),A(LSEQ),
.      A(LPSPC),A(LELMAT),A(LELMAT),NDOF)
IF (ERR) CALL SYSTEM(52,32H JOB ABORTED DUE TO ABOVE ERRORS)
IF (NSPOINT.GT.0.AND.NASTY.EQ.1) CALL PARVEC
IF (APPLD) CALL ASSL(4HLOAD,97,A(LBIG),
.      A(LSHELL),A(LBEAM),A(LCORD),A(LGRID),A(LSEQ),
.      A(LPSPC),A(LELMAT),A(LELMAT))

```

IF (ERR) CALL SYSTEM(52,32H JOB ABORTED DUE TO ABOVE ERRORS)
REWIND 4
REWIND 7
END


```

SUBROUTINE DUMP(NAME,A,M,N,MN)
DIMENSION A(M,N)
CALL SYSTEM(51,5H DUMP)
GO TO (10,20,30),MN
10 DO 11 I=1,N
11 PRINT 12,NAME,I,LUCF(A(1,I)),(A(J,I),J=1,M)
12 FORMAT (1X,A10,15,3X,06/(5(1X,020)))
RETURN
20 DO 21 I=1,N
21 PRINT 22,NAME,I,LUCF(A(1,I)),(A(J,I),J=1,M)
22 FORMAT (1X,A10,15,3X,06/(10E13,5))
RETURN
30 DO 31 I=1,N
31 PRINT 32,NAME,I,LUCF(A(1,I)),(A(J,I),J=1,M)
32 FORMAT (1X,A10,15,3X,06/(10I13))
RETURN
END

```

```

SUBROUTINE MURCOR(N)
COMMON/CUNTRL/ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD,NCORE,
      WTMASS,NASTY
LOGICAL ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD
DIMENSION X(3)
COMMON A(1)
DATA LA,IFL/0,0/
DATA KFL/40008/
IF (LA.GT.0) GO TO 5
LA=LOCFA(A)
CALL MEM(IFL)
5 NCORE=NCORE+N
10 IF (NCORE+LA.LT.IFL-4) RETURN
INCR=NCORE+LA-IFL+5
INCR=(INCR+KFL-1)/KFL
INCR=KFL*INCR
IFL=IFL+INCR
ENCODE(20,20,X) IFL
20 FORMAT (4HRFL(,06,1H))
X(3)=0
CALL REMARK(X)
CALL MEM(IFL)
GO TO 10
END

```

	IDENT	MEM
	ENTRY	MEM
MEM	BSSZ	1
	SA1	X1
	Sx7	A1
	LX1	30D
	SA7	LSTAT
	BX6	X1
	SA6	STAT
	MEMURY	SCM,STAT,RECALL
	SA1	STAT
	SA2	LSTAT
	AX1	30D
	BX6	X1
	SA6	X2
	EQ	MEM
STAT	BSSZ	1
LSTAT	BSSZ	1
	END	


```

C      FUNCTION INDEX(LIST,N,L,ITEM)
C
C      FINDS A NUMBER IN A LIST (E.G. GRID POINT NUMBERS)
C
      DIMENSION LIST(L,1)
      REAL LIST,ITEM
      DO 10 INDEX=1,N
      IF (LIST(1,INDEX).EQ.ITEM) RETURN
10 CONTINUE
      INDEX=-1
      RETURN
      END

```

```

SUBROUTINE BUM
COMMON/CONTRL/ECHO,ERR,MAIGEN,DYN,APPLD,PLOT,DBUG,BAD,NCORE,
      •      WTMASS,NASTY
LOGICAL ECHO,ERR,MAIGEN,DYN,APPLD,PLOT,DBUG,BAD
COMMON/CARD/NAME,CARD(400),MXCARD,NCARD,NTYPE,NSUB,NO
C
C      ERROR MESSAGE FOR BAD BULK DATA CARD
C
      IF (.NOT.ECHO) PRINT 10,NO,NAME,(CARD(I),I=1,NCARD)
10  FORMAT (110,1H.,5X,9A8/(24X,8A8))
      PRINT 20,NO
20  FORMAT (29H *** FORMAT ERROR ON CARD NO.,15)
      ERR=.TRUE.
      RETURN
      END

```

```

IDENT      MOVER      (A,IC,B,JC,N)
ENTRY      MOVER
VFD        30/0HMOVER,30/5

```

```

*
* COMPASS PROGRAM FOR TRANSFER OF DATA (FIN COMPILER)
*

```

```

* EQUIVALENT FORTRAN SUBROUTINE
*

```

```

* SUBROUTINE MOVER (A,IC,B,JC,N)
*

```

```

* DIMENSION A(1),B(1)
*

```

```

* IF(N.EQ.0) RETURN
*

```

```

* I=1
*

```

```

* J=1
*

```

```

* DO 100 K=1,N
*

```

```

* B(J)=A(1)
*

```

```

* I=I+IC
*

```

```

* 100 J=J+JC
*

```

```

* RETURN
*

```

```

* END
*

```

```

MUVER      BSSZ      1
SA5        A1+4      ADS(N) TO X5
SB1        1
SB3        2        2 TO B3
SA2        X5        N TO X2
SA3        A1+B1
SX6        B1        1 TO X6
SA4        A3+B3      ADS(JC) TO X4
SA3        X3        IC TO X3
SA5        A1+B3      ADS(B) TO X5
SB4        X2        N TO B4
SA4        X4        JC TO X4
SB7        X3        IC TO B7
BX7        X6*X2      MOD(N,2) TO X7
SX6        B3        2 TO X6
SB6        X4        JC TO B6
LE         B4,B0,MUVER RETURN IF N=0 OR N IS NEGATIVE
ZR         X7,EVEN
SA3        X1        ODDD, MOVE FIRST ELEMENT
SB4        B4-B1
SX1        X1+B7
BX7        X3
SX2        B4        N TO X2
SA7        X5
SX5        X5+B6
EVEN      BX7        X6*X2      MOD(N,4) TO X7
SB5        B3+B3      4 TO B5
ZR         X7,FOUR
SA3        X1
SA4        X1+B7
SX1        X1+B7

```


	SB4	B4-B3	
	HX0	X3	
	SX1	X1+B7	
	LX7	X4	
	SA0	X5	
	SA7	X5+B0	
	SX5	X5+B6	
	SX5	X5+B6	
FOUR	GE	B0,B4,MOVER	
	SX0	X5	
	SB2	B0	
	SA3	X1+B7	
	SB7	B7+B7	
	SA4	X1+B7	
	SA2	X1	
	SA5	A3+B7	
	SX1	X0+B6	
	SB7	B7+B7	
	SB0	B6+B6	
	EQ	B7,B0,ZERO	TEST FOR ZERO INCREMENT
	EQ	B0,B0,MIDD	
LOOP	SA2	A2+B7	
	SA3	A3+B7	
	SA4	A4+B7	
	SA5	A5+B7	
MIDD	BX0	X2	
	SB4	B4-B5	
	LX7	X3	
	SA6	X0+B2	
	SA7	X1+B2	
	SB2	B2+B6	
	BX0	X4	
	LX7	X5	
	SA6	X0+B2	
	SA7	X1+B2	
	SB2	B2+B6	
	GT	B4,LOOP	
	EQ	B0,B0,MOVER	
ZERO	BX6	X2	
	LX7	X2	
LOOP2	SA6	X0+B2	
	SA7	X1+B2	
	SB2	B2+B6	
	SB4	B4-B5	
	SA6	X0+B2	
	SA7	X1+B2	
	SB2	B2+B6	
	GT	B4,LOOP2	
	EQ	B0,B0,MOVER	
	END		

LOOP	RX0	X1*X2	X0=A(1)*B(1)
	SA1	A1+B2	X1=A(1+2)
	NX6	X6	NORMALIZE SUM1
	SA2	A2+B3	X2=B(1+2)
	FX7	X7+X5	SUM2=SUM2+A(1-1)*B(1-1)
	RX5	X3*X4	X5=A(1+1)*B(1+1)
	SB1	B1-B4	N=N-4
	SA3	A3+B2	X3=A(1+3)
	NX7	X7	NORMALIZE SUM2
	SA4	A4+B3	X4=B(1+3)
	FX6	X6+X0	SUM1=SUM1+A(1)*B(1)
	RX0	X1*X2	X0=A(1+2)*B(1+2)
TWO	SA1	A1+B2	X1=A(1+4)
	SA2	A2+B3	X2=B(1+4)
	FX7	X7+X5	SUM2=SUM2+A(1+1)*B(1+1)
	NX6	X6	NORMALIZE SUM1
	RX5	X3*X4	X5=A(1+3)*B(1+3)
	SA3	A3+B2	X3=A(1+5)
	SA4	A4+B3	X4=B(1+5)
	NX7	X7	NORMALIZE SUM2
TEST	FX6	X6+X0	SUM1=SUM1+A(1+2)*B(1+2)
	LT	B4,B1,LOOP	RECYCLE IF 4 LT N
	RX0	X1*X2	CLEANUP BEGINS
	NX1	X6	
	FX7	X7+X5	
	RX5	X3*X4	
	NX3	X7	
	FX2	X0+X1	
	NX6	X2	
	FX4	X3+X5	
	NX7	X4	
	FX6	X6+X7	SUM=SUM1+SUM2
	NX6	X6	
LAST	SA6	B6	
	EQ	SCPR0D	
	END		


```

IDENT          REDCLP          (N,IC,JC,FAC,A,B)
VFD            36/0HREDCLP,24/0
ENTRY          REDCLP
REDCLP         BSSZ            1
*
*
* GAUSSIAN REDUCTION LOOP.      FTN COMPILER
*
* FORTRAN EQUIVALENT
*
SUBROUTINE REDCLP (N,IC,JC,FAC,A,B)
DIMENSION A(1),B(1)
I=1
J=1
DO 100 K=1,N
B(J)=B(J)-FAC*A(I)
I=I+IC
* 100 J=J+JC
*
*
SA2           X1              N TO X2
SB3           2              2 TO B3
SA3           A1+1          ADS(IC) TO X3
SX6           1              1 TO X6
SA5           A1+B3          ADS(JC) TO X5
SB2           X2              N TO B2
SA4           A3+B3          ADS(FAC) TO X4
SA3           X3              IC TO X3
BX7           X6*X2          MOD(N,2) TO X7
SA2           A5+B3          ADS(A) TO X2
SA1           X5              JC TO X1
SA5           A4+B3          ADS(B) TO X5
SA4           X4              FAC TO X4
SB7           X3              IC TO B7
SB4           X1              JC TO B4
BX0           X4              FAC TO X0
SA4           X5              B TO X4
SB1           X1+B4          2*JC TO B1
SA1           X2              A TO X1
SA2           X2+B7          A(1+IC) TO X2
SB7           X3+B7          2*IC TO B7
RX1           X0*X1          FIRST MULTIPLY
ZR            X7,EVEN        SKIP IF N IS EVEN
FX4           X4-X1          FINISH FIRST OPERATION IF N ODD
SB2           B2-1          DECREMENT B2
SA1           A1+B7          A(1+2*IC) TO X1
NX6           X4              NORMALIZE(X4) TO X6
SA6           A4              X6 TO B(1)
ZR            B2,REDCLP      RETURN IF N=1
RX1           X0*X1          X0*X1 TP X1

```

	SA4	A4+B1	B(1+2*JC) TO X4
EVEN	RX3	X0*X2	X0*X2 TO X3
	SA5	X5+B4	B(1+JC) TO X5
	GE	B3,B2,DUNE	NO LOOP IF N=2 OR N=3
LOOP	FX4	X4-X1	
	SA2	A2+B7	
	SA1	A1+B7	
	NX6	X4	
	SA4	A4+B1	
	FX3	X5-X3	
	SA5	A5+B1	
	NX7	X3	
	SB2	B2-B3	
	SA6	A4-B1	
	RX1	X0*X1	
	KX3	X0*X2	
	SA7	A5-B1	
	LT	B3,B2,LOOP	
DUNE	FX4	X4-X1	
	NX6	X4	
	FX5	X5-X3	
	SA6	A4	
	NX7	X5	
	SA7	A5	
	EQ	B0,B0,REDCLP	
	END		

C
C
C
C

SUBROUTINE INIT

WADES THRU EXEC + CASE DECKS, SETS DEFAULT VALUES, READS
\$LOOF OPTIONS, INITIALIZES STUFF

```

COMMON/CONTRL/ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD,NCORE,
    WTMASS,NASTY
LOGICAL ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD
COMMON/LUC/LCORD,LGRID,LSEQ,LSHELL,LBEAM,LPSHEL,LPSHELX,LPBEAM,
    LMAT,LSPSC
COMMON/CARD/NAME,CARD(400),MXCARD,NCARD,NTYPE,NSUB,NO
COMMON/TITLE/TITLE(8)
COMMON/GAUSS/NGAUS,GFACT
COMMON/N/NCORD,NUMNP,NSHELL,NPSHEL,NPSHELX,NSHLTY,NBEAM,NPBEAM,
    NBMTY,NMAT,NSPOINT,MAXGRD,NDOF
DIMENSION STATUS(6)
DATA YES,NEIN/3HYES,3H NO/
READ(5,1) TITLE
1 FORMAT (8A10)
IF (EOF(5).NE.0) STOP
CALL MOVER(0,0,LCORD,1,9)
CALL MOVER(0,0,NCORD,1,12)
NCORE=1
PLOT=.FALSE.
ECHO=.FALSE.
MATGEN=.TRUE.
APPLD=.FALSE.
ERR=.FALSE.
DYN=.FALSE.
DEBUG=.FALSE.
NGAUS=2
GFACT=0
WTMASS=1.
NASTY=2
MXCARD=400
NO=0
5 READ(1,10) (CARD(I),I=1,8)
10 FORMAT (8A10)
IF (EOF(1).NE.0) GO TO 40
WRITE(3,10) (CARD(I),I=1,8)
IF (CARD(1).EQ.10HBEGIN BULK) GO TO 20
PLOT=PLOT.OR.(CARD(1).EQ.10H$LOOF PLOT .AND.CARD(2).EQ.3HYES)
ECHO=ECHO.OR.(CARD(1).EQ.10H$LOOF ECHO .AND.CARD(2).EQ.3HYES)
DEBUG=DEBUG.OR.(CARD(1).EQ.10H$LOOF DEBUG.AND.CARD(2).EQ.4H YES)
MATGEN=MATGEN.AND.CARD(1).NE.10H$LOOF MAT N
IF (CARD(1).EQ.10H$LOOF GAUSS.AND.CARD(2).EQ.2H 3) NGAUS=3
IF (CARD(1).EQ.10H$LOOF GFACT ) GFACT=DCODE(CARD(2),2,BAD)
DYN=DYN.OR.CARD(1).EQ.10H$LOOF DYN Y
IF (CARD(1).EQ.10H$NASTRAN C) NASTY=1
IF (CARD(1).EQ.10H$NASTRAN M) NASTY=2

```



```

        GO TO 5
20 CALL MUVER(NEIN,0,STATUS,1,6)
   IF (MATGEN) STATUS(1)=YES
   IF (DYN) STATUS(2)=YES
   IF (APPLD) STATUS(3)=YES
   IF (ECHO) STATUS(4)=YES
   IF (PLOT) STATUS(5)=YES
   IF (DEBUG) STATUS(6)=YES
   WRITE(6,30) TITLE,STATUS,NGAUS,NGAUS,GFACT
30 FORMAT (47H1 PRELOOF -- PREPROCESSOR FOR SEMILOOF ELEMENTS//
   .      1X,8A10//21H OPTIONS IN EFFECT --//
   .      20H MATRIX GENERATION           ,T30,A3/
   .      20H DYNAMICS                     ,T30,A3/
   .      20H DISTRIBUTED LOAD             ,T30,A3/
   .      20H ECHO                         ,T30,A3/
   .      20H PLOT                         ,T30,A3/
   .      20H DEBUG                        ,T30,A3/
   .      20H INTEGRATION RULE             ,T30,I1,IHX,I1/
   .      20H INTEGRATION FACTOR          T27,F6,4)
   IF (ECHO) WRITE(6,35)
35 FORMAT (///15X,27HB U L K   D A T A   E C H O)
   RETURN
40 CALL SYSTEM(52,20H BULK DATA MISSING  )
   END

```

```

SUBROUTINE INPUT
C
C   READS BULK DATA DECK, INTERPRETS CARDS PERTINENT TO SEMILOOF,
C   ECHOES OTHERS, DOES SOME ERROR CHECKS AND TABLE SETUPS.
C
COMMON/CARD/NAME,CARD(400),MXCARD,NCARD,NTYPE,NSUB,NO
COMMON/CONTRL/ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD,NCORE,
  WTMASS,NASTY
  LOGICAL ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD
COMMON/LOC/LCORD,LGRID,LSEQ,LSHELL,LBEAM,LPSHEL,LPSHELX,LPBEAM,
  LMAT,LPSPC
COMMON/N/NCORD,NUMNP,NSHELL,NPSHEL,NPSHELX,NSHLTY,NBEAM,NPBEAM,
  NBMTY,NMAT,NSPUIN1,MAXGRD,NDOF
COMMON A(1)
LTYPE=1
C
C   FETCH UP A BULK DATA CARD
C
10 CALL FETCH
  IF (NTYPE) 9000,10,15
15 IF (NTYPE.GE.LTYPE) GO TO 30
  IF (.NOT.ECHO) PRINT 17,NO,NAME,(CARD(I),I=1,NCARD)
17 FORMAT (110,1H.,5X,9A8/(24X,8A8))
  PRINT 20,NO
20 FORMAT (42H *** BULK DATA OUT OF ORDER, IGNORING CARD,15)
  ERR=.TRUE.
  GO TO 10
30 LTYPE=NTYPE
C
C   SEMILOOF CARD RECOGNIZED.  BRANCH TO APPROPRIATE ROUTINE
C
  GO TO
  .(100,200,300,400,500,600,700,800,900,1000,1100,1200,1300,1400),
  NTYPE
C
C   CURD2X CARDS
C
100 IF (LCORD.EQ.0) LCORD=NCORE
  CALL MORCOR(15)
  CALL CORD(A(LCORD))
  GO TO 10
C
C   GRID CARDS
C
200 IF (LGRID.EQ.0) LGRID=NCORE
  CALL MORCOR(13)
  CALL GRID(A(LCORD),A(LGRID))
  GO TO 10
C
C   SEQUENCE CARDS

```

```

C
300 IF (LSEQ.EQ.0) LSEQ=NCORE
    IF (LSEQ.EQ.NCURE) CALL MURCOR(NUMINP)
    CALL SEQGP(A(LSEQ),A(LGRID))
    GO TO 10

C
C    CLOUF8 AND CLOUF6 CARDS
C
400 IF (LSHELL.EQ.0) LSHELL=NCORE
    CALL MURCOR(12)
    CALL ISHELL(A(LSHELL),A(LGRID))
    GO TO 10

C
C    CLOUF3
C
500 IF (LBEAM.EQ.0) LBEAM=NCORE
    CALL MURCOR(21)
    CALL IBEAM(A(LBEAM),A(LGRID))
    GO TO 10

C
C    CONGRUENCE CARDS
C
600 CALL CONGR(A(LGRID),A(LSHELL),A(LBEAM))
    GO TO 10

C
C    PLOOF
C
700 IF (LPSHEL.EQ.0) LPSHEL=NCORE
    CALL MURCOR(3)
    CALL IPSHEL(A(LPSHEL))
    GO TO 10

C
C    PLOOFX
C
800 IF (LPSHELX.EQ.0) LPSHELX=NCORE
    CALL MURCOR(20)
    CALL IPSHELX(A(LPSHELX))
    GO TO 10

C
C    PLOUF3
C
900 IF (LPBEAM.EQ.0) LPBEAM=NCORE
    CALL MURCOR(9)
    CALL IPBEAM(A(LPBEAM))
    GO TO 10

C
C    MAT1 OR MAT8
C
1000 IF (LMAT.EQ.0) LMAT=NCORE
    CALL MURCOR(6)

```



```

        CALL IMAT(A(LMAT))
        GO TO 10
C
C      PLOAD4
C
1100 CALL IPLOAD(A(LGRID))
      APPLD=.TRUE.
      GO TO 10
C
C      THLOOF
C
1200 CALL ITHICK(A(LGRID))
      GO TO 10
C
C      PARAM
C
1300 CALL PARAM
      GO TO 10
C
C      SPOINI
C
1400 CALL ISPN1
      GO TO 10
C
9000 LPSPC=NCORE
      CALL MORCUR(NUMNP)
      IF (LSEQ.NE.0) RETURN
      LSEQ=NCORE
      CALL MORCOR(NUMNP)
      CALL MOVER(0,0,A(LSEQ),1,NUMNP)
      RETURN
      END

```

```

SUBROUTINE FETCH
C
C READS A BULK DATA CARD, PLUS ANY CONTINUATIONS, AND CHECKS
C AGAINST LIST OF KNOWN CARDS
C
COMMON/CUNTRL/ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD,NCORE,
.      WTMASS,NASTY
LOGICAL ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD
COMMON/CARD/NAME,CARD(400),MXCARD,NCARD,NTYPE,NSUB,NO
DIMENSION BUFFER(8)
DIMENSION LL(20),NN(20),MM(20),MECH(20)
LOGICAL DUP
LOGICAL MECH,FULL
DATA PLUS,STAR/1H+,1H*/
DATA LL/8HCORD2R ,
2      8HCORD2C ,
3      8HCORD2S ,
4      8HGRID ,
5      8HGRID* ,
6      8HSEQGP ,
7      8HCLOUF8 ,
8      8HCLOUF6 ,
9      8HCLOUF3 ,
.      8HCONGR ,
.      8HPLOUF ,
.      8HPLOUFx ,
.      8HPLOUF3 ,
.      8HMAT1 ,
.      8HMAT8 ,
.      8HPLOAD4 ,
.      8HTHLOUF ,
.      8HPARAM ,
.      8HSPPOINT /
DATA NN/1,
.      1,
.      1,
.      2,
.      2,
.      3,
.      4,
.      4,
.      5,
.      6,
.      7,
.      8,
.      9,
.      10,
.      10,
.      11,
.      12,

```

```

      13,
      14/
DATA MM/1,
      2,
      3,
      1,
      2,
      1,
      1,
      2,
      1,
      1,
      1,
      1,
      1,
      1,
      1,
      2,
      1,
      1,
      1,
      1/
DATA MECH/
9      .TRUE.,
      .TRUE.,
      .TRUE.,
      .FALSE.,
      .FALSE.,
      .TRUE.,
      .FALSE.,
      .FALSE.,
      .FALSE.,
      .FALSE.,
      .FALSE.,
      .FALSE.,
      .FALSE.,
      .FALSE.,
      .FALSE.,
      .FALSE.,
      .TRUE.,
      .FALSE.,
      .FALSE.,
      .FALSE.,
      .TRUE.,
      .TRUE./
DATA KNOWN/19/
DATA FULL/.FALSE./
IF (.NOT.FULL) READ(1,10) BUFFER
10 FORMAT (8A10)
IF (EOF(1).NE.0.OR.BUFFER(1).EQ.7HENDDATA) GO TO 200
FULL=.TRUE.
DECODE(80,20,BUFFER) NAME,(CARD(I),I=1,8),C,CONT
20 FORMAT (9A8,A1,A7)
IF (NAME.EQ.6HPLOOF2) NAME=6HPLOOFX

```



```

      CALL LSHIFT(CARD,8)
      NU=NU+1
      FULL=.FALSE.
      DO 30 K=1,KNOWN
      IF (NAME.EQ.LL(K)) GO TO 40
30  CONTINUE
      NTYPE=0
      NSUB=0
      DUP=.TRUE.
      GO TO 50
40  NTYPE=NN(K)
      NSUB=MM(K)
      DUP=MECH(K)
C
C      PASS THIS CARD TO NASTRAN
C
50  IF (DUP) WRITE(3,20) NAME,(CARD(I),I=1,8),C,CONT
C
C      AND ECHO, IF DESIRED
C
      IF (ECHO) PRINT 60,NU,NAME,(CARD(I),I=1,8),C,CONT
60  FORMAT (110,1H,,5X,9A8,A1,A7)
      L2=8
70  READ(1,10) BUFFER
      FULL=.TRUE.
      IF (EOF(1).NE.0) GO TO 200
      L1=L2+1
      L2=L1+7
80  FORMAT (A1,A7,8A8,A1,A7)
      DECODE(80,80,BUFFER) CHAR,CONT2,(CARD(L),L=L1,L2),C,CONT3
      CALL LSHIFT(CARD(L1),8)
      IF (CHAR.NE.PLUS.AND.CHAR.NE.STAR) GO TO 120
      FULL=.FALSE.
      IF (DUP) WRITE(3,80) CHAR,CONT2,(CARD(L),L=L1,L2),C,CONT3
      IF (ECHO) PRINT 90,CHAR,CONT2,(CARD(L),L=L1,L2),C,CONT3
90  FORMAT (16X,A1,A7,8A8,A1,A7)
      IF (CONT2.EQ.CONT) GO TO 110
      PRINT 100,NO
100  FORMAT (48H *** CONTINUATION CARD OUT OF ORDER FOR CARD NO.,15)
      L1=L1-8
      L2=L2-8
      GO TO 70
110  CONT=CONT3
      GOTO 70
120  NCARD=L2-8
      IF (NCARD.LE.MXCARD) RETURN
      PRINT 130,NO
130  FORMAT (13H *** CARD NO.,15,10H TOO LONG)
      IF (,.NOT.ECHO) PRINT 60,NU,NAME,(CARD(I),I=1,8)
      L2=MXCARD

```

```
RETURN  
200 NTYPE=-1  
RETURN  
END
```

```

SUBROUTINE LSHIFT(CARD,N)
DIMENSION CARD(1)
INTEGER BLANK
DATA MASK/77 00 00 00 00 00 00 00 00 00B/
DATA BLANK/1L /
DO 20 I=1,N
DO 10 J=1,8
IF ((CARD(I).AND.MASK).NE.BLANK) GO TO 20
CARD(I)=SHIFT(CARD(I),6)
10 CONTINUE
20 CONTINUE
RETURN
END

```

```

      REAL FUNCTION DCODE (ITEM,TYPE,ERR)
      IMPLICIT INTEGER (A-Z)
      LOGICAL BAD,ERR
      DIMENSION CHAR(8)
      DATA PLUS,MINUS,ZERO,NINE,BLANK,E,PERIOD/
      .      1H+,1H-,1H0,1H9,1H ,1HE,1H./
      DECODE(8,5,ITEM) CHAR
5  FORMAT (8A1)
      GO TO (100,200,300),TYPE

C
C      INTEGER
C
100  L=-1
      BAD=.FALSE.
      DO 110 I=1,8
      C=CHAR(9-I)
      IF (C.NE.BLANK.AND.L.LT.0) L=1-I
      IF (C.GE.ZERO.AND.C.LE.NINE) GO TO 110
      IF (C.EQ.PLUS.OR.C.EQ.MINUS) GO TO 110
      IF (C.EQ.BLANK) GO TO 110
      BAD=.TRUE.
110  CONTINUE
      IF (BAD) GO TO 400
      IC=ITEM
      IF (L.GT.0) IC=SHIFT(ITEM,60-6*L)
      DECODE(8,120,IC) M
120  FORMAT (18)
      DCODE=M
      GO TO 400

C
C      FLOATING POINT
C
200  BAD=.FALSE.
      L=-1
      DO 210 I=1,8
      C=CHAR(9-I)
      IF (C.NE.BLANK.AND.L.LT.0) L=1-I
      IF (C.GE.ZERO.AND.C.LE.NINE) GO TO 210
      IF (C.EQ.PLUS.OR.C.EQ.MINUS) GO TO 210
      IF (C.EQ.E.OR.C.EQ.BLANK) GO TO 210
      IF (C.EQ.PERIOD) GO TO 210
      BAD=.TRUE.
210  CONTINUE
      IF (BAD) GO TO 400
      IC=ITEM
      IF (L.GT.0) IC=SHIFT(ITEM,60-6*L)
C      CROSS YOUR FINGERS
      DECODE(8,220,IC) DCODE
220  FORMAT (E8.0)
      GO TO 400

```



```

C
C      ALPHA
C
300  BAD=.FALSE.
      IF (CHAR(1).NE.BLANK) GO TO 315
      DO 310 I=2,8
      IF (CHAR(I).NE.BLANK) GO TO 320
310  CONTINUE
315  ENCODE(8,5,DCODE) CHAR
      GO TO 400
320  I1=I-1
      ENCODE(8,5,DCODE) (CHAR(J),J=1,8),(BLANK,J=1,I1)
400  ERR=ERR.OR.BAD
      RETURN
      END

```

```

REAL FUNCTION DCODE2(ITEM,TYPE,ERR)
IMPLICIT INTEGER (A-Z)
LOGICAL BAD,ERR
DIMENSION ITEM(2)
DIMENSION CHAR(16)
DIMENSION TEMP(2)
DATA PLUS,MINUS,ZERO,NINE,BLANK,E,PERIOD/
      1H+,1H-,1H0,1H9,1H ,1HE,1H./
      DECODE(18,10,ITEM) CHAR
10  FORMAT (8A1,2X,8A1)
      GO TO (100,200,300),TYPE

C
C      INTEGER
C
100  L=-1
      BAD=,FALSE,
      DO 110 I=1,16
      C=CHAR(17-I)
      IF (C,NE,BLANK,AND,L,LT,0) L=I-1
      IF (C,GE,ZERO,AND,C,LE,NINE) GO TO 110
      IF (C,EQ,PLUS,OR,C,EQ,MINUS) GO TO 110
      IF (C,EQ,BLANK) GO TO 110
      BAD=,TRUE,
110  CONTINUE
      DCODE2=0
      IF (L,EQ,-1) RETURN
      IF (BAD) GO TO 400
      DCODE2=0
      IF (L) 160,140,130
130  L16=16-L
      ENCODE(16,240,TEMP) (BLANK,I=1,L),(CHAR(I),I=1,L16)
      GO TO 150
140  ENCODE(16,240,TEMP) CHAR
150  DECODE(16,270,TEMP) DCODE2
160  GO TO 400

C
C      FLOATING POINTS
C
200  BAD=,FALSE,
      L=-1
      DO 210 I=1,16
      C=CHAR(17-I)
      IF (C,NE,BLANK,AND,L,LT,0) L=I-1
      IF (C,GE,ZERO,AND,C,LE,NINE) GO TO 210
      IF (C,EQ,PLUS,OR,C,EQ,MINUS) GO TO 210
      IF (C,EQ,BLANK,OR,C,EQ,E) GO TO 210
      IF (C,EQ,PERIOD) GO TO 210
      BAD=,TRUE,
210  CONTINUE
      IF (BAD) GO TO 400

```

```

        DCODE2=0
        IF (L) 280,250,230
230    L16=16-L
        ENCODE(16,240,TEMP) (BLANK,I=1,L),(CHAR(I),I=1,L16)
240    FORMAT (16A1)
        GO TO 260
250    ENCODE(16,240,TEMP) CHAR
260    DECODE(16,270,TEMP) DCODE2
270    FORMAT (E16.0)
280    GO TO 400
C
C      ALPHA
C
300    BAD=.FALSE.
        IF (CHAR(1).NE.BLANK) GO TO 315
        DO 310 I=2,16
        IF (CHAR(I).NE.BLANK) GO TO 320
310    CONTINUE
315    ENCODE(8,240,DCODE2) (CHAR(I),I=1,8)
        GO TO 400
320    IF (1.GT.9) GO TO 330
        17=1+7
        ENCODE(8,240,DCODE2) (CHAR(J),J=1,17)
        GO TO 400
330    19=1-9
        ENCODE(8,240,DCODE2) (CHAR(J),J=1,16),(BLANK,J=1,19)
400    ERR=ERR.OR.BAD
        RETURN
        END

```

```

SUBROUTINE SWING(A, ITYPE)
  DIMENSION A(3)
  IF (ITYPE.EQ.1) RETURN
  IF (ITYPE.NE.2) GO TO 10
  THETA=A(2)*3.1415927/180
  A(2)=A(1)*SIN(THETA)
  A(1)=A(1)*COS(THETA)
  RETURN
10 IF (ITYPE.NE.3) RETURN
  THETA=A(2)*3.1415927/180
  PHI=A(3)*3.1415927/180
  A(3)=A(1)*COS(THETA)
  P=A(1)*SIN(THETA)
  A(2)=P*SIN(PHI)
  A(1)=P*COS(PHI)
  RETURN
END

```



```

SUBROUTINE IMATRIX(T,A,B,C)
DIMENSION T(3,3),A(3),B(3),C(3)
DO 10 J=1,3
    T(J,2)=C(J)-A(J)
10  T(J,3)=B(J)-A(J)
    CALL NORMAL(T(1,3))
    CALL CROSS(T(1,3),T(1,2),T(1,2))
    CALL NORMAL(T(1,2))
    CALL CROSS(T(1,2),T(1,3),T(1,1))
RETURN
END

```

```
SUBROUTINE CROSS(A,B,C)
DIMENSION A(3),B(3),C(3)
C1=A(2)*B(3)-B(2)*A(3)
C2=-(A(1)*B(3)-B(1)*A(3))
C3=A(1)*B(2)-B(1)*A(2)
C(1)=C1
C(2)=C2
C(3)=C3
RETURN
END
```

```
SUBROUTINE NORMAL(X)
DIMENSION X(3)
E=SQRT(X(1)*X(1)+X(2)*X(2)+X(3)*X(3))
IF (E.EQ.0) RETURN
X(1)=X(1)/E
X(2)=X(2)/E
X(3)=X(3)/E
RETURN
END
```

```

SUBROUTINE CORD(TRANS)
C
C INTERPRETS CORD BULK DATA CARDS AND SETS UP COORDINATE
C TRANSFORMATION TABLES
C
  DIMENSION TRANS(3,5,1)
  COMMON/CONTRL/ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD,NCORE,
    .   WTMASS,NASTY
  LOGICAL ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD
  COMMON/CARD/NAME,CARD(400),MXCARD,NCARD,NTYPE,NSUB,NO
  COMMON/N/NCORD,NUMNP,NSHELL,NPSHEL,NPSHELX,NSHLTY,NBEAM,NPBEAM,
    .   NBM1Y,NMAT,NSPOINT,MAXGRD,NDUF
  DIMENSION AA(3),BB(3),CC(3)
  IF (NCARD.L1.9) GO TO 20
  NCORD=NCORD+1
  TRANS(1,1,NCORD)=DCODE(CARD,1,BAD)
  TRANS(2,1,NCORD)=NSUB
  DO 10 I=1,3
  AA(I)=DCODE(CARD(2+I),2,BAD)
  BB(I)=DCODE(CARD(5+I),2,BAD)
10 CC(I)=DCODE(CARD(8+I),2,BAD)
  J=DCODE(CARD(2),1,BAD)
  BAD=BAD.OR.J.NE.0
  IF (BAD) GO TO 20
  CALL TMRX(TRANS(1,3,NCORD),AA,BB,CC)
  CALL MUVER(AA,1,TRANS(1,2,NCORD),1,3)
  RETURN
20 CALL BUM
  RETURN
  END

```



```

SUBROUTINE GRID(CURD,GRD)
DIMENSION CURD(3,5,1),GRD(13,1)

C
C   INTERPRETS GRID CARDS AND SETS UP TABLES
C
COMMON/CONTRL/ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD,NCORE,
•   WTMASS,NASTY
LOGICAL ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD
COMMON/CARD/NAME,CARD(400),MXCARD,NCARD,NTYPE,NSUB,NO
COMMON/N/NCURD,NUMNP,NSHELL,NPSHEL,NPSHELX,NSHLTY,NBEAM,NPBEAM,
•   NBMTY,NMAT,NSPOINT,MAXGRD,NDUF
DIMENSION XYZ(3)
NUMNP=NUMNP+1
BAD=.FALSE.
IF (NSUB.EQ.2) GO TO 10

C
C   SINGLE FIELD
C
GRD(1,NUMNP)=DCODE(CARD(1),1,BAD)
GRD(2,NUMNP)=DCODE(CARD(2),1,BAD)
GRD(3,NUMNP)=DCODE(CARD(3),2,BAD)
GRD(4,NUMNP)=DCODE(CARD(4),2,BAD)
GRD(5,NUMNP)=DCODE(CARD(5),2,BAD)
GRD(9,NUMNP)=DCODE(CARD(6),1,BAD)
GRD(10,NUMNP)=DCODE(CARD(7),1,BAD)
GO TO 20

C
C   DOUBLE FIELD
C
10 GRD(1,NUMNP)=DCODE2(CARD(1),1,BAD)
GRD(2,NUMNP)=DCODE2(CARD(3),1,BAD)
GRD(3,NUMNP)=DCODE2(CARD(5),2,BAD)
GRD(4,NUMNP)=DCODE2(CARD(7),2,BAD)
GRD(5,NUMNP)=DCODE2(CARD(9),2,BAD)
GRD(9,NUMNP)=DCODE2(CARD(11),1,BAD)
GRD(10,NUMNP)=DCODE2(CARD(13),1,BAD)
20 IF (.NOT.BAD) GO TO 30
CALL BUM
RETURN
30 MAXGRD=MAX0(MAXGRD,IFIX(GRD(1,NUMNP)))
IF (NUMNP.EQ.1) GO TO 38
N1=NUMNP-1
DO 32 I=1,N1
IF (GRD(1,I).EQ.GRD(1,NUMNP)) GO TO 35
32 CONTINUE
GO TO 38
35 PRINT 36,IFIX(GRD(1,NUMNP))
36 FORMAT (19H *** DUPLICATE GRID,16)
ERR=.TRUE.
RETURN

```

```

38 DO 40 I=1,3
   GRD(1+5,NUMNP)=GRD(1+2,NUMNP)
40 GRD(1+10,NUMNP)=0
   GRD(2,NUMNP)=0
   GRD(9,NUMNP)=0
   IF (GRD(2,NUMNP).LE.0) GO TO 80
   JCORD=INDEX(CORD,NCORD,15,GRD(2,NUMNP))
   IF (JCORD.GT.0) GO TO 60
   M=GRD(1,NUMNP)
   N=GRD(2,NUMNP)
   PRINT 50,M,N
50 FORMAT ( 9H *** GRID,18,37H REFERENCES UNKNOWN COORDINATE SYSTEM,
   .      18)
   ERR=.TRUE.
   GO TO 80
C
C   CONVERT LOCATION TO BASIC COORDINATES
C
60 CALL MOVER(GRD(3,NUMNP),1,XYZ,1,3)
   CALL SWING(XYZ,IF1X(CORD(2,1,JCORD)))
   DO 70 I=1,3
   CALL SCPROD(3,3,1,CORD(1,3,JCORD),XYZ,XX)
   GRD(2,NUMNP)=JCORD
70 GRD(5+I,NUMNP)=XX+CORD(1,2)
80 IF (GRD(9,NUMNP).LE.0) RETURN
   JCORD=INDEX(CORD,NCORD,15,GRD(9,NUMNP))
   IF (JCORD.GT.0) GO TO 90
   IF (JCORD.GT.0) RETURN
   M=GRD(1,NUMNP)
   N=GRD(9,NUMNP)
   PRINT 50,M,N
   ERR=.TRUE.
   RETURN
90 GRD(9,NUMNP)=JCORD
   RETURN
   END

```

```

SUBROUTINE SEWGP(SEQ,GRD)
C
C
C
  DIMENSION SEQ(1),GRD(13,1)
  INTEGER SEQ
  COMMON/CONTRL/ECHO,ERR,MAIGEN,DYN,APPLD,PLOT,DBUG,BAD,NCORE,
    .      WIMASS,NASTY
  LOGICAL ECHO,ERR,MAIGEN,DYN,APPLD,PLOT,DBUG,BAD
  COMMON/CARD/NAME,CARD(400),MXCARD,NCARD,NTYPE,NSUB,NO
  COMMON/N/NCORD,NUMNP,NSHELL,NPSHEL,NPSHELX,NSHLTY,NBEAM,NPBEAM,
    .      NBMIY,NMAT,NSPINT,MAXGRD,NDOF
  BAD=.FALSE.
  DO 20 I=1,7,2
    GR=DCODE(CARD(I),1,BAD)
    SQ=DCODE(CARD(I+1),1,BAD)
    IF (GR.EQ.0) GO TO 20
    J=INDEX(GRD,NUMNP,13,GR)
    IF (J.NE.-1) GO TO 10
  5 BAD=.TRUE.
    GO TO 20
  10 IF (SQ.LT.0.OR.SQ.GT.NUMNP) GO TO 5
    SEQ(J)=SQ
  20 CONTINUE
    IF (BAD) CALL BUM
    RETURN
  END

```

```

SUBROUTINE ISHELL(SHELL,GRD)
C
C PROCESSES CLOUFB AND CLOUFB CARDS
C
  DIMENSION SHELL(12,1),GRD(13,1)
  COMMON/CUNTRL/ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD,NCORE,
    * WTMASS,NASTY
  LOGICAL ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD
  COMMON/CARD/NAME,CARD(400),MXCARD,NCARD,NTYPE,NSUB,NO
  COMMON/N/NCORD,NUMNP,NSHELL,NPSHEL,NPSHELX,NSHLTY,NBEAM,NPBEAM,
    * NBMTY,NMAT,NSPOINT,MAXGRD,NDOF
  BAD=.FALSE.
  NSHELL=NSHELL+1
  SHELL(9,NSHELL)=0
  SHELL(10,NSHELL)=0
  NNOD=8
  IF (NSUB.EQ.2) NNOD=6
  N2=NNOD+2
  BAD=NCARD.LT.N2
  DO 10 I=1,N2
10 SHELL(I,NSHELL)=DCODE(CARD(I),1,BAD)
  SHELL(11,NSHELL)=NSHELL
  SHELL(12,NSHELL)=0
  IF (NCARD.GE.N2+1) SHELL(12,NSHELL)=DCODE(CARD(N2+1),2,BAD)
  IF (.NOT.BAD) GO TO 15
  CALL BUM
  RETURN
C
C GET GRIDS
C
15 DO 30 I=3,N2
  L=SHELL(1,NSHELL)
  J=INDEX(GRD,NUMNP,13,FLOAT(L))
  SHELL(1,NSHELL)=J
  IF (J.NE.-1) GO TO 30
  PRINT 20,IFIX(SHELL(1,NSHELL)),L
20 FORMAT (12H *** ELEMENT,18,24H REFERENCES UNKNOWN GRID,110)
  ERR=.TRUE.
30 CONTINUE
  RETURN
  END

```



```

C      SUBROUTINE IBEAM(BEAM,GRD)
C
C      PROCESSES CLOJF3 CARDS
C
      DIMENSION BEAM(21,1),GRD(13,1)
      COMMON/CONTRL/ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD,NCORE,
      .      WTMASS,NASTY
      LOGICAL ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD
      COMMON/CARD/NAME,CARD(400),MXCARD,NCARD,NTYPE,NSUB,NO
      COMMON/N/NCORD,NUMNP,NSHELL,NPSHEL,NPSHELX,NSHLTY,NBEAM,NPBEAM,
      .      NBMTY,NMAT,NSPOINT,MAXGRD,NDOF
      DIMENSION F(3)
      BAD=.FALSE.
      NBEAM=NBEAM+1
      DO 10 I=1,8
10  BEAM(I,NBEAM)=DCODE(CARD(I),1,BAD)
      IF (BEAM(7,NBEAM).EQ.0) BEAM(7,NBEAM)=BEAM(6,NBEAM)
      IF (BEAM(8,NBEAM).EQ.0) BEAM(8,NBEAM)=BEAM(6,NBEAM)
      CALL MOVER(0,0,BEAM(16,NBEAM),1,6)
      BEAM(15,NBEAM)=NBEAM
      IF (BAD) CALL BUM
      IF (NCARD.LT.12) RETURN
      IF=DCODE(CARD(12),1,BAD)
      IF (IF.EQ.0) GO TO 45
      BAD=BAD.OR.(IF.NE.1.AND.IF.NE.2)
      IF (BAD) GO TO 100
      GO TO (20,30),IF
20  DO 22 I=1,3
22  F(I)=DCODE(CARD(8+I),2,BAD)
      IF (BAD) GO TO 100
      GO TO 40
30  L=DCODE(CARD(9),1,BAD)
      IF (BAD) GO TO 100
      J=INDEX(GRD,NGRID,13,L)
      IF (J.EQ.-1) GO TO 80
      DO 32 I=1,3
32  F(I)=GRD(5+I,J)
40  CALL SCALAR(F,F,FSQ)
      IF (FSQ.EQ.0.) GO TO 100
      FSQ=SQRT(FSQ)
      DO 42 I=1,3
42  BEAM(18+I,NBEAM)=F(I)/FSQ
45  CALL MOVER(0,0,BEAM(9,NBEAM),1,6)
      IF (NCARD.LE.16) GO TO 60
      DO 50 I=17,22
50  BEAM(I-8,NBEAM)=DCODE(CARD(I),2,BAD)
60  DO 70 I=3,5
      L=BEAM(1,NBEAM)
      J=INDEX(GRD,NUMNP,13,FLOAT(L))
      BEAM(1,NBEAM)=J

```

```

      IF (J.EQ.-1) GO TO 80
70  CONTINUE
      RETURN
80  PRINT 90,IFIX(BEAM(1,NBEAM)),L
90  FORMAT (12H *** ELEMENT,18,24H REFERENCES UNKNOWN GRID,110)
      ERR=.TRUE.
      RETURN
100 CALL BUM
      RETURN
      END

```

```

SUBROUTINE CONGR(GRD,SHELL,BEAM)
C
C
C
  DIMENSION GRD(13,1),SHELL(12,1),BEAM(18,1)
  COMMON/CUNTRL/ECHO,ERR,MAIGEN,DYN,APPLD,PLOT,DBUG,BAD,NCORE,
    •   WTMASS,NASTY
  LOGICAL ECHO,ERR,MAIGEN,DYN,APPLD,PLOT,DBUG,BAD
  COMMON/CARD/NAME,CARD(400),MXCARD,NCARD,NTYPE,NSUB,NO
  COMMON/N/NCORD,NUMNP,NSHELL,NPSHEL,NPSHELX,NSHLTY,NBEAM,NPBEAM,
    •   NBMTY,NMAT,NSPOINT,MAXGRD,NDOF
  EL1=DCODE(CARD(1),1,BAD)
  J1=INDEX(SHELL,NSHELL,12,EL1)
  IF (J1.LT.0) GO TO 30
  DO 10 I=2,NCARD
  EL=DCODE(CARD(I),1,BAD)
  IF (BAD) GO TO 10
  IF (EL.EQ.0) GO TO 10
  J=INDEX(SHELL,NSHELL,12,EL)
  BAD=BAD.OR.J.LT.0
  IF (.NOT.BAD) SHELL(11,J)=SHELL(11,J1)
10 CONTINUE
  IF (BAD) CALL BUM
  RETURN
30 J1=INDEX(BEAM,NBEAM,18,EL1)
  BAD=J1.LT.0
  DO 40 I=2,NCARD
  EL=DCODE(CARD(I),1,BAD)
  IF (BAD) GO TO 40
  IF (EL.EQ.0) GO TO 40
  J=INDEX(BEAM,NBEAM,18,EL)
  BAD=BAD.OR.J.LT.0
  IF (.NOT.BAD) BEAM(15,J)=BEAM(15,J1)
40 CONTINUE
  IF (BAD) CALL BUM
  RETURN
END

```

```

SUBROUTINE IPSHELL(PSHELL)
C
C PROCESSES PLOOF CARDS -- SHELL PROPERTIES
C
  DIMENSION PSHELL(3,1)
  COMMON/CUNTKL/ECHO,ERR,MAIGEN,DYN,APPLD,PLOT,DBUG,BAD,NCORE,
    • WIMASS,NASTY
  LOGICAL ECHO,ERR,MAIGEN,DYN,APPLD,PLOT,DBUG,BAD
  COMMON/CARD/NAME,CARD(400),MXCARD,NCARD,NTYPE,NSUB,NO
  COMMON/N/NCORD,NUMNP,NSHELL,NPSHEL,NPSHELX,NSHLTY,NBEAM,NPBEAM,
    • NBMTY,NMAT,NSPOINT,MAXGRD,NDUF
  BAD=.FALSE.
  NPSHEL=NPSHEL+1
  PSHELL(1,NPSHEL)=DCODE(CARD(1),1,BAD)
  PSHELL(2,NPSHEL)=DCODE(CARD(2),1,BAD)
  PSHELL(3,NPSHEL)=DCODE(CARD(3),2,BAD)
  IF (BAD) CALL BUM
  RETURN
END

```



```

SUBROUTINE IPSHELL(PSHELL)
C
C PROCESSES PLODFX CARDS -- ANISOTROPIC SHELL
C
  DIMENSION PSHELL(20,1)
  COMMON/CONTRL/ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD,NCORE,
    . WIMASS,NASTY
  LOGICAL ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD
  COMMON/N/NCORD,NUMNP,NSHELL,NPSHEL,NPSHELLX,NSHLTY,NBEAM,NPBEAM,
    . NBMTY,NMAT,NSPOINT,MAXGRD,NDOF
  COMMON/CARD/NAME,CARD(400),MXCARD,NCARD,NTYPE,NSUB,NO
  BAD=.FALSE.
  IF (NCARD.LT.7) GO TO 40
  NPSHELLX=NPSHELLX+1
  PSHELL(1,NPSHELLX)=DCODE(CARD(1),1,BAD)
  DO 10 I=2,7
10 PSHELL(1,NPSHELLX)=DCODE(CARD(1),2,BAD)
  PSHELL(20,NPSHELLX)=DCODE(CARD(9),2,BAD)
  CALL MOVER(0,0,PSHELL(8,NPSHELLX),1,12)
  IF (NCARD.LT.15) GO TO 35
  DO 20 I=10,15
20 PSHELL(1-2,NPSHELLX)=DCODE(CARD(1),2,BAD)
  IF (NCARD.LT.23) GO TO 35
  DO 30 I=18,23
30 PSHELL(1-4,NPSHELLX)=DCODE(CARD(1),2,BAD)
35 IF (.NOT.BAD) RETURN
40 CALL BUM
  RETURN
  END

```

```

SUBROUTINE IPBEAM(PBEAM)
C
C PROCESSES PLOTF3 CARDS -- BEAM PROPERTIES
C
  DIMENSION PBEAM(9,1)
  COMMON/CONTRL/ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD,NCORE,
    WTMASS,NASTY
  LOGICAL ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD
  COMMON/CARD/NAME,CARD(400),MXCARD,NCARD,NTYPE,NSUB,NO
  COMMON/N/NCORD,NUMNP,NSHELL,NPSHEL,NPSHELX,NSHLTY,NBEAM,NPBEAM,
    NBMTY,NMAT,NSPOINT,MAXGRD,NDOF
  NPBEAM=NPBEAM+1
  BAD=.FALSE.
  PBEAM(1,NPBEAM)=DCODE(CARD(1),1,BAD)
  DO 10 I=2,8
10 PBEAM(I,NPBEAM)=DCODE(CARD(I),2,BAD)
  PBEAM(9,NPBEAM)=0
  IF (NCARD.GE.9) PBEAM(9,NPBEAM)=DCODE(CARD(9),2,BAD)
  BAD=BAD.OR.PBEAM(2,NPBEAM).LE.0
  IF (BAD) CALL BUM
  RETURN
END

```

```

SUBROUTINE IMAT(XMAT)
  DIMENSION XMAT(6,1)
  COMMON/CNTRL/ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD,NCURE,
    • WTMASS,NASTY
  LOGICAL ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD
  COMMON/CARD/NAME,CARD(400),MXCARD,NCARD,NTYPE,NSUB,NU
  COMMON/N/NCORD,NUMNP,NSHELL,NPSHEL,NPSHELX,NSHLTY,NBEAM,NPBEAM,
    • NBMTY,NMAT,NSPUINT,MAXGRD,NDOF
  REAL NU
  NMAT=NMAT+1
  BAD=.FALSE.
  XMAT(1,NMAT)=DCODE(CARD(1),1,BAD)
  IF (NSUB.EQ.2) GO TO 10
  E=DCODE(CARD(2),2,BAD)
  G=DCODE(CARD(3),2,BAD)
  NU=DCODE(CARD(4),2,BAD)
  RHO=DCODE(CARD(5),2,BAD)
  IF (G.EQ.0) G=E/(2*(1+NU))
  IF (NU.EQ.0) NU=E/(2*G)-1
  BAD=BAD.OR.NU.GE..5.OR.NU.LE.-1
  IF (.NOT.BAD) GO TO 5
  CALL BUM
  RETURN
5  XMAT(2,NMAT)=E/(1-NU*NU)
  XMAT(3,NMAT)=NU*XMAT(2,NMAT)
  XMAT(4,NMAT)=XMAT(2,NMAT)
  XMAT(5,NMAT)=G
  XMAT(6,NMAT)=RHO
  RETURN
10 DO 20 I=2,6
20 XMAT(1,NMAT)=DCODE(CARD(1),2,BAD)
  IF (BAD) CALL BUM
  RETURN
END

```

```

SUBROUTINE ILOAD(GRD)
C
C PROCESSES PLOAD4 CARDS -- PRESSURES BY GRID POINTS
C
  DIMENSION GRD(13,1)
  COMMON/CONTRL/ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD,NCORE,
    . WTMASS,NASTY
  LOGICAL ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD
  COMMON/CARD/NAME,CARD(400),MXCARD,NCARD,NTYPE,NSUB,NO
  COMMON/N/NCORD,NUMNP,NSHELL,NPSHEL,NPSHELX,NSHLTY,NBEAM,NPBEAM,
    . NBM1Y,NMAT,NSPOINT,MAXGRD,NDOF
  BAD=.FALSE.
  VALUE=DCODE(CARD(2),2,BAD)
  THRU=DCODE(CARD(4),3,BAD)
  IF (THRU.EQ.4HTHRU) GO TO 100
  DO 30 I=3,8
  GR=DCODE(CARD(I),1,BAD)
  IF (.NOT.BAD) GO TO 5
  CALL BUM
  GO TO 30
5 IF (GR.EQ.0) GO TO 30
  J=INDEX(GRD,NUMNP,13,GR)
  IF (J.NE.-1) GO TO 20
  PRINT 10,IFIX(GR)
10 FORMAT (40H *** PLOAD4 CARD REFERENCES UNKNOWN GRID,18)
  ERR=.TRUE.
  RETURN
20 GRD(12,J)=VALUE
30 CONTINUE
  RETURN
100 GR1=DCODE(CARD(3),1,BAD)
  GR2=DCODE(CARD(5),1,BAD)
  BAD=BAD.OR.GR2.LT.GR1
  IF (.NOT.BAD) GO TO 110
  CALL BUM
  RETURN
110 J1=INDEX(GRD,NUMNP,13,GR1)
  J2=INDEX(GRD,NUMNP,13,GR2)
  GR=GR+1
  IF (J1.GT.0) GO TO 120
  PRINT 10,J1
  BAD=.TRUE.
120 IF (J2.GT.0) GO TO 130
  PRINT 10,J2
  BAD=.TRUE.
130 IF (.NOT.BAD) GO TO 140
  CALL BUM
  RETURN
140 J=INDEX(GRD,NUMNP,13,GR)
  IF (J.EQ.-1) GO TO 150

```



```

SUBROUTINE ITHICK(GRD)
DIMENSION GRD(13,1)

C
C
C
PROCESS THLOOF CARDS (SHELL THICKNESS BY NODE POINT)

COMMON/CONTRL/ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD,NCORE,
    WIMASS,NASTY
    LOGICAL ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD
COMMON/CARD/NAME,CARD(400),MXCARD,NCARD,NTYPE,NSUB,NO
COMMON/N/NCORD,NUMNP,NSHELL,NPSHEL,NPSHELX,NSHLTY,NBEAM,NPBEAM,
    NBMTY,NMAT,NSPOINT,MAXGRD,NDOF
    BAD=.FALSE.
    VALUE=DCODE(CARD(2),2,BAD)
    THRU=DCODE(CARD(4),3,BAD)
    IF (THRU.EQ.4HTHRU) GO TO 100
    DO 30 I=3,8
    GR=DCODE(CARD(1),1,BAD)
    IF (,NOT,BAD) GO TO 5
    CALL BUM
    GO TO 30
5 IF (GR.EQ.0) GO TO 30
J=INDEX(GRD,NUMNP,13,GR)
IF (J,NE,-1) GO TO 20
PRINT 10,IFIX(GR)
10 FORMAT (40H *** THLOOF CARD REFERENCES UNKNOWN GRID,18)
ERR=.TRUE.
GO TO 30
20 GRD(12,J)=VALUE
30 CONTINUE
RETURN
100 GR1=DCODE(CARD(3),1,BAD)
GR2=DCODE(CARD(5),1,BAD)
BAD=BAD.OR.GR2.LT.GR1
IF (,NOT,BAD) GO TO 110
CALL BUM
RETURN
110 J1=INDEX(GRD,NUMNP,13,GR1)
J2=INDEX(GRD,NUMNP,13,GR2)
IF (J1,GT.0) GO TO 120
PRINT 10,J1
BAD=.TRUE.
120 IF (J2,GT.0) GO TO 130
PRINT 10,J2
BAD=.TRUE.
130 IF (,NOT,BAD) GO TO 140
CALL BUM
RETURN
140 J=INDEX(GRD,NUMNP,13,GR)
IF (J,EQ,-1) GO TO 150
GRD(12,J)=VALUE

```

```

SUBROUTINE PARAM
COMMON/CARD/NAME,CARD(400),MXCARD,NCARD,NTYPE,NSUB,NO
COMMON/CONTRL/ECHO,ERR,MAIGEN,DYN,APPLD,PLOT,DBUG,BAD,NCORE,
    • WIMASS,NASTY
LOGICAL ECHO,ERR,MAIGEN,DYN,APPLD,PLOT,DBUG,BAD
BAD=.FALSE.
A=DCODE(CARD(1),3,BAD)
IF (A.NE.6HWIMASS) GO TO 10
WIMASS=DCODE(CARD(2),2,BAD)
10 IF (BAD) CALL BUM
RETURN
END

```

```

SUBROUTINE ISPNT
COMMON/CONTRL/ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD,NCORE,
      WTMASS,NASTY
      LOGICAL ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD
COMMON/N/NCORD,NUMNP,NSHELL,NPSHEL,NPSHELX,NSHLTY,NBEAM,NPBEAM,
      NBM1Y,NMAT,NSPOINT,MAXGRD,NDOF
COMMON/CARD/NAME,CARD(400),MXCARD,NCARD,NTYPE,NSUB,NO
BAD=.FALSE.
IF (CARD(2).EQ.4HTHRU) GO TO 20
DO 10 I=2,8
J=DCODE(CARD(I),1,BAD)
IF (J.EQ.0) GO TO 10
BAD=BAD.OR.J.LE.MAXGRD
NSPOINT=NSPOINT+1
10 CONTINUE
IF (BAD) CALL BUM
RETURN
20 J1=DCODE(CARD(1),1,BAD)
J2=DCODE(CARD(3),1,BAD)
BAD=BAD.OR.J2.LT.J1.OR.J1.LE.MAXGRD
IF (BAD) GO TO 30
NSPOINT=NSPOINT+J2-J1+1
RETURN
30 CALL BUM
RETURN
END

```

```

SUBROUTINE PROCESS(GRID,SHELL,BEAM,ISEQ,IPSPC)
DIMENSION GRID(13,1),SHELL(12,1),BEAM(21,1),ISEQ(1),IPSPC(1)
COMMON/CONTROL/ECHO,ERR,MAIGEN,DYN,APPLD,PLOT,DBUG,BAD,NCORE,
    WIMASS,NASTY
LOGICAL ECHO,ERR,MAIGEN,DYN,APPLD,PLOT,DBUG,BAD
COMMON/N/NCORD,NUMNP,NSHELL,NPSHEL,NPSHELX,NSHLTY,NBEAM,NPBEAM,
    NBMTY,NMAT,NSPUINT,MAXGRD,NDOF
COMMON/LUC/LCORD,LGRID,LSEQ,LSHELL,LBEAM,LPSHEL,LPSHELX,LPBEAM,
    LMAT,LPSPC
COMMON/TITLE/TITLE(8)
COMMON A(1)
CALL MOVER(-1,0,IPSPC,1,NUMNP)

C
C SEQUENCE GRID POINTS
C
IF (ISEQ(1).NE.0) GO TO 10
LTEMP=NCORE
CALL MORCOR(NUMNP)
CALL SURT(GRID,A(LTEMP),ISEQ)
NCORE=LTEMP
10 IF (NSHELL.EQ.0) GO TO 90
C
C SHELL ELEMENTS
C
NSHLTY=0
DO 80 I=1,NSHELL
C
C FIGURE OUT IF THIS IS A QUAD OR TRIANGLE
C
LNODZ=8
IF (SHELL(10,I).EQ.0) LNODZ=6
IF (SHELL(11,I).EQ.1) NSHLTY=NSHLTY+1
LK=LNODZ+2
KL=LNODZ+1
C
C IDENTIFY LOOF NODES
C
DO 50 K=3,KL,2
LL=SHELL(K,I)
IF (LL.EQ.0) GO TO 50
MM=SHELL(K+1,I)
NN=SHELL(K+2,I)
IF (K.EQ.KL) NN=SHELL(3,I)
IF (,.NOT.PLOT) GO TO 35
IP1=10*SHELL(1,I)+K-2
IP2=IP1+1
LLL=GRID(1,LL)
MMM=GRID(1,MM)
NNN=GRID(1,NN)
WRITE(3,20) IP1,LLL,MMM,IP2,MMM,NNN

```



```

20      FORMAT (6HPLUTEL,19,2(318,8X))
35      IF (IPSPC(LL).NE.6) GO TO 40
        WRITE(6,30) IFIX(GRID(1,LL))
30      FORMAT (4H *** GRID,18,
        39H USED AS BOTH LOUF NODE AND CORNER NODE)
        ERR=.TRUE.
40      IPSPC(LL)=456
50      CONTINUE
        DO 70 K=4,LK,2
          LL=SHELL(K,1)
          IF (LL.EQ.0) GO TO 70
          IF (IPSPC(LL).NE.456) GO TO 60
          WRITE(6,30) IFIX(GRID(1,LL))
          ERR=.TRUE.
60          IPSPC(LL)=6
70          CONTINUE
80      CONTINUE
C
C      BEAM ELEMENTS
C
90      IF (NBEAM.EQ.0) GO TO 150
        NBMTY=0
        DO 140 I=1,NBEAM
          IF (BEAM(15,I).EQ.1) NBMTY=NBMTY+1
C
C      IDENTIFY LOUF NODES
C
        DO 110 K=3,5,2
          LL=BEAM(K,1)
          IF (LL.EQ.0) GO TO 110
          IF (IPSPC(LL).NE.6) GO TO 100
          WRITE(6,30) IFIX(GRID(1,LL))
          ERR=.TRUE.
100         IPSPC(LL)=0
110        CONTINUE
          K=4
          LL=BEAM(K,1)
          IF (LL.EQ.0) GO TO 130
          IF (IPSPC(LL).NE.456.AND.IPSPC(LL).NE.0) GO TO 120
          WRITE(6,30) IFIX(GRID(1,LL))
          ERR=.TRUE.
120         IPSPC(LL)=6
130         CONTINUE
140        CONTINUE
150      WRITE(7) TITLE,NUMNP,NSHELL,NPSHEL,NSHLTY,NBEAM,NPBEAM,
        NBMTY,NMAT
        IF (NSHELL.GT.0) WRITE(7) ((SHELL(J,1),J=1,12),I=1,NSHELL)
        IF (NBEAM.GT.0) WRITE(7) ((BEAM(J,1),J=1,15),I=1,NBEAM)
        DO 190 I=1,NUMNP
          GRID(11,I)=1SEQ(I)

```

```

      IF (IPSPC(1).NE.-1) GO TO 170
      PRINT 160,IFIX(GRID(1,1))
160    FORMAT (19H ** WARNING -- GRID,18,
      .      31H HAS NO LOUF ELEMENTS CONNECTED )
      IPSPC(1)=0
170    WRITE(3,180) IFIX(GRID(1,1)),IFIX(GRID(2,1)),GRID(3,1),
      .      GRID(4,1),1,1,GRID(5,1),IFIX(GRID(9,1)),IPSPC(1)
180    FORMAT (5HGRID*,19,2116,2E16,9,2H+G,13/2H*G,13,19,E16,9,2116)
190    CONTINUE
      IF (MATGEN) WRITE(3,195)
195    FORMAT (17HPARAM STIFMAT 1)
      IF (DYN) WRITE(3,200)
200    FORMAT (17HPARAM MASSMAT 1)
      IF (APPLD) WRITE(3,210)
210    FORMAT (17HPARAM LOADVEC 1)
      IF (NSPOINT.GT.0) WRITE(3,220) NSPOINT
220    FORMAT (16HPARAM NSPOINT ,18)
      WRITE(3,230)
230    FORMAT (7HENDDATA)
      RETURN
      END

```

```

SUBROUTINE SORT(GRID,IGRD,ISEQ)
DIMENSION GRID(13,1),IGRD(1),ISEQ(1)
C
C   SORT GRID POINTS ACCORDING TO EXTERNAL G.P. NO.
C
COMMON/N/NCORD,NUMNP,NSHELL,NPSHEL,NPSHELX,NSHLTY,NBEAM,NPBEAM,
      NBM1Y,NMAT,NSPOINT,MAXGRD,NDUF
LOGICAL SWAP
DO 5 I=1,NUMNP
  IGRD(I)=GRID(1,I)
S   ISEQ(I)=1
  N1=NUMNP-1
10  SWAP=.FALSE.
  DO 20 I=1,N1
    IF (IGRD(I).LT.IGRD(I+1)) GO TO 20
    L=ISEQ(I)
    ISEQ(I)=ISEQ(I+1)
    ISEQ(I+1)=L
    L=IGRD(I)
    IGRD(I)=IGRD(I+1)
    IGRD(I+1)=L
    SWAP=.TRUE.
20  CONTINUE
    IF (SWAP) GO TO 10
  DO 30 I=1,NUMNP
30  IGRD(ISEQ(I))=I
    CALL MUVER(IGRD,1,ISEQ,1,NUMNP)
    RETURN
  END

```

```

SUBROUTINE XSHELL(GRD,SHELL,PSHELL,PSHELLX,XMAT,ELSTIF,ELMASS,
.      ELOAD,
.      ELSTRES,SINDEX,SHAPE,SEQ)
C
C      DRIVES SHELL ELEMENT MATRIX GENERATOR
C
      DIMENSION GRD(13,1),SHELL(12,1),PSHELL(3,1),PSHELLX(20,1),
.      XMAT(6,1)
      DIMENSION ELSTIF(32,32),ELMASS(32,32)
      DIMENSION ELOAD(32),ELSTRES(6,32)
      DIMENSION SINDEXT(1),SEQ(1)
      INTEGER SEQ
      DIMENSION LNODS(8),PRESS(8)
      INTEGER SINDEXT
      COMMON/CUNTRL/ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD,NCORE,
.      WTMASS,NASTY
      LOGICAL ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD
      COMMON/N/NCORD,NUMNP,NSHELL,NPSHEL,NPSHELLX,NSHLTY,NBEAM,NPBEAM,
.      NBMTY,NMAT,NSPOINT,MAXGRD,NDOF
      DIMENSION ELXYZ(9,4)
      DIMENSION AA(6),BB(6),DD(6)
      IF (NSHELL.EQ.0) RETURN
C
C      LOOP ON ELEMENTS
C
      DO 200 I=1,NSHELL
C
C      GET MATRIX INDEX AND SEE IF MATRICES FOR THIS CONGRUENCE
C      GROUP HAVE BEEN GENERATED YET
C
      LSEQ=SHELL(11,I)
      IF (SINDEXT(LSEQ+1).NE.0) GO TO 200
      LPROP=INDEX(PSHELL,NPSHEL,3,SHELL(2,I))
      IF (LPROP.LE.0) GO TO 20
      LTYPE=1
      LMAT=INDEX(XMAT,NMAT,6,PSHELL(2,LPROP))
      THIK=PSHELL(3,LPROP)
      CALL MUVER(XMAT(1,LMAT),1,AA,1,6)
      DENS=XMAT(6,LMAT)
      IF (LMAT.GT.0) GO TO 40
      WRITE(6,10) IFIX(PSHELL(1,LPROP)),IFIX(PSHELL(3,LPROP))
10  FORMAT (19H *** SHELL PROPERTY,16,
.      28H REFERENCES UNKNOWN MATERIAL,16)
      ERR=.TRUE.
      GO TO 200
20  LPROP=INDEX(PSHELLX,NPSHELLX,20,SHELL(2,I))
      IF (LPROP.GT.0) GO TO 30
      WRITE(6,25) IFIX(SHELL(1,I)),IFIX(SHELL(2,I))
25  FORMAT (18H *** SHELL ELEMENT,18,
.      28H REFERENCES UNKNOWN PROPERTY,18)

```



```

      ERR=.TRUE.
      GO TO 200
30  LTYPE=2
      CALL MOVER(PSHELLX(2,LPROP),1,AA,1,6)
      CALL MOVER(PSHELLX(8,LPROP),1,BB,1,6)
      CALL MOVER(PSHELLX(14,LPROP),1,DD,1,6)
      DENS=PSHELLX(20,LPROP)
      THIK=1.0
C
C      GATHER UP EIGHT (SIX) CORNER-POINT PROPERTIES
C
40  CALL MOVER(0,0,ELXYZ(9,1),9,4)
      LNODS(7)=0
      LNODS(8)=0
      DO 60 J=1,8
          NOD=SHELL(J+2,1)
          IF (NOD.EQ.0) GO TO 65
          SHELL(J+2,1)=SEQ(NOD)
          LNODS(J)=NOD
          PRESS(J)=GRD(12,NOD)
          DO 50 K=1,3
50      ELXYZ(J,K)=GRD(K+2,NOD)
60      ELXYZ(J,4)=GRD(13,NOD)
C
C      THICKNESS ASSOCIATED WITH SHELL OVERRIDES G.P. VALUE
C
65  IF (THIK.NE.0) CALL MOVER(THIK,0,ELXYZ(1,4),1,8)
C
C      HERE WE GO
C
      CALL QSHELL(ELXYZ,LNODS,LTYPE,DENS,AA,BB,DD,SHELL(12,1),
      .      PRESS,ELSTIF,ELMASS,ELOAD,ELSTRES,SHAPE,LVABZ)
      IF (.NOT.BAD) GO TO 67
      WRITE(6,66) IFIX(SHELL(1,1))
66  FORMAT (18H *** SHELL ELEMENT,16,17H HAS BAD GEOMETRY)
      ERR=.TRUE.
      GO TO 200
C
C      DUMP OUT MATRICES TO DISK
C
67  CALL WRITMS(98,ELSTIF,32*32,LSEQ)
      IF (DYN) CALL WRITMS(99,ELMASS,32*32,LSEQ)
      IF (APPLD) CALL WRITMS(97,ELOAD,32,LSEQ)
      IF (.NOT.DBUG) GO TO 200
C
C      DEBUG PRINTOUT
C
      WRITE(6,70) I
70  FORMAT (*OSTIFFNESS MATRIX FOR SHELL ELEMENT*,14)
      DO 80 J=1,LVABZ

```

```

80 WRITE(6,90) J,(ELSTIF(J,K),K=1,LVABZ)
90 FORMAT (* ROW*,13/(8E15,6))
   WRITE(6,100) I
100 FORMAT (*STRESS MATRIX FOR SHELL ELEMENT*,14)
   DO 110 J=1,6
110 WRITE(6,90) J,(ELSTRES(J,K),K=1,LVABZ)
   IF (.NOT.DYN) GO TO 140
   WRITE(6,120) I
120 FORMAT (*UMASS MATRIX FOR SHELL ELEMENT*,14)
   DO 130 J=1,LVABZ
130 WRITE(6,90) J,(ELMASS(J,K),K=1,LVABZ)
140 IF (.NOT.APPLD) GO TO 200
   WRITE(6,150) I,(ELOAD(J),J=1,LVABZ)
150 FORMAT (*LOAD VECTOR FOR SHELL ELEMENT*,13/(8E15,6))
200 CONTINUE
   RETURN
   END

```

```

SUBROUTINE WSHLL(XYZI, LNODS, LTYPE, DENS, AA, BB, DD, ANGLE, PRESS,
.   ELSTIF, ELMASS, ELUAD, ELSTRES, WSHLL, LVABZ)
DIMENSION XYZI(6,4), LNODS(8), AA(6), BB(6), DD(6), PRESS(8),
1   ELSTIF(32,32), ELMASS(32,32), ELUAD(32)
DIMENSION ELSTRES(6,32)
COMMON/GAUSS/NGAUS, GFACI
COMMON/N/NCORD, NUMNP, NSHELL, NPSHEL, NPSHELX, NSHLTY, NBEAM, NPBEAM,
.   NBMTY, NMAT, NSPOINT, MAXGRD, NDOF
COMMON/CONTRL/ECHO, ERR, MATGEN, DYN, APPLD, PLOT, DEBUG, BAD, NCORE,
.   WTMASS, NASTY
LOGICAL ECHO, ERR, MATGEN, DYN, APPLD, PLOT, DEBUG, BAD
DIMENSION ABD(6,6)
DIMENSION CGAUS3(3)
DATA P1/3.1415927/
DATA XGAUS3, XGAUS2/.77549 66692 4148, .57735 02691 8962/
DATA CGAUS3/.555555555555556, .888888888888889, .555555555555556/
LNODZ=8
IF (LNODS(7).EQ.0) LNODZ=6
LVABZ=4*LNODZ
DATA TH1K/0./
CALL MOVER(0,0,ELSTIF,1,32*32)
CALL MOVER(0,0,ELMASS,1,32*32)
CALL MOVER(0,0,ELUAD,1,32)

C
C FILL IN STRESS-STRAIN MATRIX
C
CALL MOVER(0,0,ABD,1,36)
IF (LTYPE.EQ.2) GO TO 110
ABD(1,1)=AA(2)
ABD(1,2)=AA(3)
ABD(2,1)=AA(3)
ABD(2,2)=AA(4)
ABD(3,3)=AA(5)
GO TO 130
110 ABD(1,1)=AA(1)
ABD(1,2)=AA(2)
ABD(2,1)=AA(2)
ABD(1,3)=AA(3)
ABD(3,1)=AA(3)
ABD(2,2)=AA(4)
ABD(2,3)=AA(5)
ABD(3,2)=AA(5)
ABD(3,3)=AA(6)
ABD(1,4)=BB(1)
ABD(1,5)=BB(2)
ABD(2,4)=BB(2)
ABD(1,6)=BB(3)
ABD(3,4)=BB(3)
ABD(2,5)=BB(4)
ABD(2,6)=BB(5)

```

```

      ABD(3,5)=BB(5)
      ABD(5,6)=BB(6)
      DO 120 I=1,3
      DO 120 J=1,3
120  ABD(I+3,J)=ABD(J,1+3)
      ABD(4,4)=DD(1)
      ABD(4,5)=DD(2)
      ABD(5,4)=DD(2)
      ABD(4,6)=DD(3)
      ABD(6,4)=DD(3)
      ABD(5,5)=DD(4)
      ABD(5,6)=DD(5)
      ABD(6,5)=DD(5)
      ABD(6,6)=DD(6)
130  IF (ANGLE.EQ.0) GO TO 150
      C=COS(ANGLE*180/PI)
      S=SIN(ANGLE*180/PI)
      C2=C*C
      S2=S*S
      C4=C2*C2
      S4=S2*S2
      DO 140 I=1,4,3
      DO 140 J=1,4,3
      A11=ABD(1,J)
      A12=ABD(1,J+1)
      A22=ABD(1+1,J+1)
      ABD(1,J)=A11*C4+A12*S2*C2+A22*S4
      ABD(1,J+1)=(A11+A22-2*A12)*S2*C2
      ABD(J+1,1)=ABD(1+1,J)
140  ABD(1+1,J+1)=A11*S4+A12*S2*C2+A22*C4
C
150  IF (LNODZ.EQ.6) GO TO 40
      IF (NGAUS.EQ.3) GO TO 20
C
C QUADRILATERAL, 4-POINT GAUSS INTEGRATION
C
      IF (GFACT.EQ.0) WRITE(7) 4,LPRUP
      IF (GFACT.NE.0) WRITE(7) 5,LPRUP
      DO 10 IX=1,2
      DO 10 JY=1,2
      CALL ZHELL(XGAUS2*(3-2*IX),XGAUS2*(3-2*JY),
      .        XYZ1,ABD,LTYPE,PRESS,DENS,WSHEL,ELSTRES,ELSTIF,ELMASS,ELOAD,
      .        LNODS,LNODZ,LVABZ,1.-GFACT/4.)
      IF (BAD) RETURN
10  CONTINUE
C
C OPTIONAL EXTRA POINT AT CENTER (QUAD)
C
      IF (GFACT.NE.0.) CALL ZHELL(0.,0.,
      .        XYZT,ABD,LTYPE,PRESS,DENS,WSHEL,ELSTRES,ELSTIF,ELMASS,ELOAD,

```



```

        LNODS, LNODZ, LVABZ, GFACT)
    IF (BAD) RETURN
    GO TO 55
C
C QUADRILATERAL, 9-POINT GAUSS INTEGRATION
C
20 WRITE(7) 9, LPROP
    DO 30 IX=1,3
    DO 30 JY=1,3
    CALL ZHELL(XGAUS3*(2-IX), XGAUS3*(2-JY),
    .          XYZI, ABD, LTYPE, PRESS, DENS, WSHEL, ELSTRES, ELSTIF, ELMASS, ELOAD,
    .          LNODS, LNODZ, LVABZ, CGAUS3(IX)*CGAUS3(JY))
    IF (BAD) RETURN
30 CONTINUE
    GO TO 55
C
C 3-POINT GAUSS INTEGRATION ON TRIANGLE
C
40 WRITE(7) 3, LPROP
    DO 50 JT=1,3
    CALL ZHELL(((4-JT)*(JT-1))/4., MOD(JT,2)/2.,
    .          XYZI, ABD, LTYPE, PRESS, DENS, WSHEL, ELSTRES, ELSTIF, ELMASS, ELOAD,
    .          LNODS, LNODZ, LVABZ, 1.)
    IF (BAD) RETURN
50 CONTINUE
55 DO 60 J=1, LVABZ
    DO 60 I=1, J
    ELSTIF(J,I)=ELSTIF(I,J)
60 ELMASS(J,I)=ELMASS(I,J)
    RETURN
    END

```

```

SUBROUTINE ZHELL(X1,ETA,XYZT,ABD,LTYPE,PRESS,DENS,WSHEL,
.   ELSTRES,ELSTIF,ELMASS,ELLOAD,LNODS,LNODZ,LVABZ,FACT)
DIMENSION XYZT(9,4),LNODS(6),ABD(6,6),ELSTRES(6,32),
.   ELSTIF(32,32),ELMASS(32,32),ELLOAD(32)
COMMON/N/NCURD,NUMNP,NSHELL,NPSHEL,NPSHELX,NSHLTY,NBEAM,NPBEAM,
.   NBMTY,NMAT,NSPOINT,MAXGRD,NDOF
COMMON/CONTRL/ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD,NCORE,
.   WTMASS,NASTY
LOGICAL ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD
DIMENSION D(6,6),FRAM(3,3),POINT(3),XITA(2),WSHEL(13,45),B(6,45)
XITA(1)=X1
XITA(2)=ETA
6 CALL HALOUF(1,AREA,XYZT,FRAM,POINT,SIDE,THIK,WSHEL,XITA,LNODS,
*   LNODZ,LVABZ,LNODZ)
BAD=.FALSE.
IF (AREA.EQ.-69.) GO TO 60

```

C

C B MATRIX

C

```

DO 8 N=1,LVABZ
B(1,N)=WSHEL(4,N)
B(2,N)=WSHEL(7,N)
B(3,N)=WSHEL(5,N)+WSHEL(6,N)
B(4,N)=WSHEL(10,N)
B(5,N)=WSHEL(12,N)
8 B(6,N)=2.0*WSHEL(11,N)

```

C

C D MATRIX

C

```

CALL MUVER(ABD,1,D,1,36)
IF (LTYPE.EQ.2) GO TO 15
DO 10 I=1,3
DO 10 J=1,3
D(1,J)=D(1,J)*THIK
10 D(I+3,J+3)=D(1,J)*THIK*THIK/12.

```

C

C STRESS AND STIFFNESS MATRICES

C

```

15 DO 30 J=1,LVABZ
DO 20 K=1,6
CALL SCPROD(6,1,1,D(1,K),B(1,J),XX)
20 ELSTRES(K,J)=XX*AREA
DO 30 I=1,J
CALL SCPROD(6,1,1,ELSTRES(1,J),B(1,1),BDB)
30 ELSTIF(1,J)=ELSTIF(1,J)+FACT*BDB
WRITE(7) ELSTRES,WSHEL,POINT,FRAM,THIK

```

C

C LOAD VECTOR

C

```

IF (APPLD)

```

```

      .CALL REDCLP(LVABZ,3,1,-AREA*PRESS*FACT,WSHEL(3,1),ELOAD)
C
C MASS MATRIX
C
      IF (.NOT.DYN) GO TO 50
      AA=DENS*AREA*THIK*WTMASS
      DO 40 J=1,LVABZ
      DO 40 I=1,J
      CALL SCPRUD(3,1,1,WSHEL(1,I),WSHEL(1,J),BB)
40  ELMASS(1,J)=ELMASS(1,J)+AA*BB
50  RETURN
60  BAD=.TRUE.
      RETURN
      END

```

```

SUBROUTINE XBEAM(GRID,BEAM,PBEAM,XMAT,
.     ELSIF,ELMASS,ELOAD,ELSTRES,SINDEX,SHAPE,ISEQ)
.     DIMENSION GRID(13,1),BEAM(21,1),PBEAM(9,1),XMAT(6,1),
.     ELSIF(17,17),ELMASS(17,17),ELOAD(17),SINDEX(1),SHAPE(1),ISEQ(1)
.     (1)
C
C     DRIVER FOR BEAM ELEMENT MATRIX GENERATOR
C
COMMON A(1)
COMMON/CONTRL/ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD,NCORE,
.     WTMASS,NASTY
LOGICAL ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD
COMMON/N/NCORD,NUMNP,NSHELL,NPSHEL,NPSHELX,NSHLTY,NBEAM,NPBEAM,
.     NBMTY,NMAT,NSPULN1,MAXGRD,NDUF
DIMENSION ELXYZ(3,3),AREA(3),X11(3),X12(3),XJ(3),ZNORM(3),ZBIN(3),
.     PRESS(3),LNODB(3)
INTEGER SINDEX
IF (NBEAM.EQ.0) RETURN
LETRANS=NCORE
CALL MURCOR(17*17)
LQ=NCORE
CALL MURCOR(17*17)
DO 130 I=1,NBEAM
LSEQ=BEAM(15,I)
IF (SINDEX(LSEQ+NSHLTY+1).NE.0) GO TO 130
LMAI=INDEX(XMAT,NMAT,1,BEAM(2,I))
IF (LMAI.GT.0) GO TO 20
WRITE(6,10) IFIX(BEAM(1,1))
10 FORMAT (17H *** BEAM ELEMENT,16,
.     28H REFERENCES UNKNOWN MATERIAL,16)
ERR=.TRUE.
GO TO 130
20 BAD=.FALSE.
DO 50 J=1,3
NOD=BEAM(J+2,1)
BEAM(J+2,1)=ISEQ(NOD)
LNODB(J)=NOD
IPB=INDEX(PBEAM,NPBEAM,9,BEAM(5+J,1))
BAD=.FALSE.
DO 30 K=1,3
30 ELXYZ(J,K)=GRID(2+K,LNODB(J))
AREA(J)=PBEAM(2,IPB)
X11(J)=PBEAM(3,IPB)
X12(J)=PBEAM(4,IPB)
XJ(J)=PBEAM(5,IPB)
ZNORM(J)=BEAM(7+2*J,1)
ZBIN(J)=BEAM(8+2*J,1)
IF (IPB.GT.0) GO TO 50
WRITE(6,40) IFIX(BEAM(1,1)),IFIX(BEAM(5+J,1))
40 FORMAT (17H *** BEAM ELEMENT,16,28H REFERENCES UNKNOWN PROPERTY,16)

```



```

      )
      BAD=.TRUE.
50    CONTINUE
      ERR=ERR.OR.BAD
      IF (BAD) GO TO 130
      CALL ZBEAM(ELXYZ,XMAT(1,LMAT),AREA,X11,X12,XJ,PRESS,ZNORM,ZBIN,
      .    BEAM(19,1),LNODB,
      .    ELSTIF,ELMASS,ELOAD,ELSTRES,A(LETRANS),A(LG),SHAPE)
      CALL WRITMS(98,ELSTIF,17*17,LSEQ+NSHLTY)
      IF (DYN) CALL WRITMS(99,ELMASS,17*17,LSEQ+NSHLTY)
      IF (APPLD) CALL WRITMS(97,ELOAD,17,LSEQ+NSHLTY)
      IF (.NOT.BAD) GO TO 55
      WRITE(6,54) IFIX(BEAM(1,1))
54    FORMAT (17H *** BEAM ELEMENT,16,17H HAS BAD GEOMETRY)
      ERR=.TRUE.
      GO TO 130
55    IF (.NOT.DBUG) GOTO 130
      WRITE(6,60) 1
60    FORMAT (*0STIFFNESS MATRIX FOR BEAM ELEMENT*,I4)
      DO 70 J=1,17
70    WRITE(6,80) J,(ELSTIF(J,K),K=1,17)
80    FORMAT (* ROW*,I3/(8E15,6))
      IF (.NOT.DYN) GO TO 110
      WRITE(6,90) 1
90    FORMAT (*0MASS MATRIX FOR BEAM ELEMENT*,I4)
      DO 100 J=1,17
100   WRITE(6,80) J,(ELMASS(J,K),K=1,17)
110   IF (.NOT.APPLD) GO TO 130
      WRITE(6,120) 1,ELOAD
120   FORMAT (*0LOAD VECTOR FOR BEAM ELEMENT*,I3/(8E15,6))
130   CONTINUE
C
C     GIVE BACK BORROWED CORE
C
      NCORE=LETRANS
      RETURN
      END

```

```

SUBROUTINE ZBEAM(XYZI,XMAT,ZAREA,ZX11,ZX12,ZXJ,ZPRESS,ZNORM,ZBIN,
.   ORIENT,LNUDB,
.   ELSTIF,ELMASS,ELOAD,ELSTRES,ETRANS,Q,WBEAM)
DIMENSION XYZI(3,3),XMAT(6),
.   ZAREA(3),ZX11(3),ZX12(3),ZXJ(3),ZPRESS(3),ZNORM(3),ZBIN(3),
.   ,ORIENT(3),LNUDB(3),
.   ELSTIF(17,17),ELMASS(17,17),ELOAD(17),ELSTRES(6,32),ETRANS(17,17),
.   ,17),
.   Q(17,17),WBEAM(13,45)

```

C
C
C

```

GETS BEAM STIFFNESS, MASS, AND/OR LOAD MATRICES

```

```

DIMENSION XGAUS(5),CGAUS(5),B(6,32),T(3,3,3),OFFSET(3,3)
REAL NU
DIMENSION FRAM(3,3),D(6),POINT(3)
COMMON/N/NCORD,NUMNP,NSHELL,NPSHEL,NPSHELX,NSHLTY,NBEAM,NPBEAM,
.   NBMTY,NMAT,NSPOINT,MAXGRD,NDOF
COMMON/CONTRL/ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD,NCORE,
.   WTMASS,NASTY
LOGICAL ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD
DATA XGAUS/-.577350269,0.577350269,-.7745966692,.7745966692,0./
DATA CGAUS/2*1.,2*.555555556,.688888889/
DATA LVABZ/17/
NU=XMAT(3)/XMAT(2)
E=(1-NU*NU)*XMAT(2)
G=XMAT(5)
DENS=XMAT(6)
CALL MUVER(0,0,ELSTIF,1,17*17)
CALL MUVER(0,0,ELMASS,1,17*17)
CALL MUVER(0,0,ELOAD,1,17)

```

C
C
C

```

LOOP ON GAUSS POINTS, END POINTS, AND MIDDLE POINT

```

```

DO 70 IGAUS=1,5
XI=XGAUS(IGAUS)
CALL GCBSF1(XI,POINT)

```

C
C
C

```

INTERPOLATE CROSS-SECTION PROPERTIES

```

```

AREA=0.
XI1=0.
XI2=0.
XJ=0.
PRESS=0.
DO 10 J=1,3
  AREA=AREA+ZAREA(J)*POINT(J)
  XI1=XI1+ZXI1(J)*POINT(J)
  XI2=XI2+ZXI2(J)*POINT(J)
  PRESS=PRESS+ZPRESS(J)*POINT(J)
10  XJ=XJ+ZXJ(J)*POINT(J)

```

10

```

C
C      INDICATE CROSS-SECTION ORIENTATION
C
      CALL MUVER(0,0,FRAM,1,9)
      CALL MUVER(ORIENT,1,FRAM,1,3)
C
C      GET SHAPE FUNCTIONS
C
      CALL LOFBEM(1,LVABZ,SIDE,XYZI,FRAM,POINT,WBEAM,XI,3,3,3,LNODB)
      BAD=.FALSE.
      IF (SIDE.EQ.-69.) GO TO 120
      IF (XI.EQ.0.) CALL MUVER(FRAM,1,I(1,1,2),1,9)
      FACMOD=E*SIDE*CGAUS(1GAUS)
      FACSHE=G*FACMOD/E
C
C      FORM STRAIN-DISPLACEMENT TERMS
C
      CALL MUVER(WBEAM(4,1),13,B(1,1),6,LVABZ)
      CALL MUVER(WBEAM(8,1),13,B(2,1),6,LVABZ)
      CALL MUVER(WBEAM(9,1),13,B(3,1),6,LVABZ)
      CALL MUVER(WBEAM(11,1),13,B(4,1),6,LVABZ)
      CALL MUVER(WBEAM(12,1),13,B(5,1),6,LVABZ)
      CALL MUVER(WBEAM(10,1),13,B(6,1),6,LVABZ)
C
C      STIFFNESS TERMS
C
      D(1)=FACMOD*AREA
      D(2)=FACMOD*X11
      D(3)=FACMOD*X12
      D(4)=FACSHE*AREA/1.2
      D(5)=FACSHE*AREA/1.2
      D(6)=FACSHE*XJ
C
C      LOAD TERMS
C
      IF (1GAUS.LT.3.AND.APPLD.AND.PRESS.NE.0)
        CALL REDCLP(LVABZ,1,3,1,-SIDE*PRESS,ELoad,WBEAM(3,1))
C
C      STIFFNESS MATRIX
C
      IF (1GAUS.GT.2) GO TO 30
      DO 20 J=1,LVABZ
      DO 20 I=1,J
20  ELSTIF(I,J)=ELSTIF(I,J)+B(1,J)*D(1)*B(1,I)+
      )          B(4,J)*D(4)*B(4,I)+B(5,J)*D(5)*B(5,I)
      GO TO 50
30  DO 40 J=1,LVABZ
      DO 40 I=1,J
40  ELSTIF(I,J)=ELSTIF(I,J)+B(2,J)*D(2)*B(2,I)+
      )          B(3,J)*D(3)*B(3,I)+B(6,J)*D(6)*B(6,I)

```

```

      GO TO 70
C
C   MASS MATRIX
C
50  IF (.NOT.DYN) GO TO 70
    AA=DENS*SIDE*AREA*WTMASS
    DO 60 J=1,LVABZ
      DO 60 I=1,J
        CALL SCPROD(3,1,1,WBEAM(1,I),WBEAM(1,J),BB)
    60  ELMASS(1,J)=ELMASS(1,J)+AA*BB
    70  CONTINUE
C
C   FILL IN SYMMETRIC TERMS
C
      DO 80 J=1,LVABZ
        DO 80 I=1,J
          ELSTIF(J,I)=ELSTIF(I,J)
    80  ELMASS(J,I)=ELMASS(I,J)
C
C   HANDLE OFFSET TRANSFORMATION
C
      DO 90 I=1,3
        IF (ZNORM(I).NE.0.) GO TO 100
        IF (ZBIN(I).NE.0.) GO TO 100
    90  CONTINUE
      RETURN
    100 CALL LOFBEM(1,LVABZ,SIDE,XYZI,T(1,1,1),POINT,WBEAM,-1.,3,3,3,LNODB)
        )
        CALL LOFBEM(1,LVABZ,SIDE,XYZI,T(1,1,3),POINT,WBEAM,1.,3,3,3,LNODB)
        DO 110 I=1,3
          DO 110 J=1,3
    110  OFFSET(I,J)=T(J,2,1)*ZNORM(I)+T(J,3,1)*ZBIN(1)
        CALL TRANS(ETRANS,OFFSET)
        CALL TRIPLE(ETRANS,ELSTIF,ETRANS,17,17,Q)
        RETURN
    120 BAD=.TRUE.
        RETURN
      END

```



```

SUBROUTINE TRANS(E,A)
DIMENSION E(17,17),A(3,3)
CALL MUVER(0,0,E,1,17*17)
CALL MUVER(1.,0,E,18,17)
E(1,13)=-A(1,3)
E(1,14)=A(1,2)
E(2,12)=A(1,3)
E(2,14)=-A(1,1)
E(3,12)=-A(1,2)
E(3,13)=A(1,1)
E(4,7)=-A(2,3)/2.
E(4,8)=-A(2,3)/2.
E(6,7)=A(2,1)/2.
E(6,8)=A(2,1)/2.
E(9,16)=-A(3,3)
E(9,17)=A(3,2)
E(10,15)=A(3,3)
E(10,17)=-A(3,1)
E(11,15)=-A(3,2)
E(11,16)=A(3,1)
RETURN
END

```

```

SUBROUTINE ASSY(NAME,IN,BIG,SHELL,BEAM,CORD,GRID,SEQ,
.      IPSPC,SMAT,BMAT,NBIG)
.      DIMENSION BIG(NBIG,6),SHELL(12,1),BEAM(21,1),CORD(3,5,1),
.      GRID(13,1),SEQ(1),IPSPC(1),SMAT(32,32),BMAT(17,17)
.      INTEGER SEQ
C
C      ASSEMBLE MASTER STIFFNESS OR MASS MATRIX
C
COMMON/N/NCORD,NUMNP,NSHELL,NPSHEL,NPSHELX,NSHLTY,NBEAM,NPBEAM,
.      NBMTY,NMAT,NSPOINT,MAXGRD,NDOF
COMMON/CONTRL/ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD,NCORE,
.      KTMASS,NASTY
LOGICAL ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD
DIMENSION T11(3,3),T12(3,3),TJ1(3,3),TJ2(3,3),Q(3,3)
DIMENSION MANE(2)
C
C      MATRICES ISH AND IBM GIVE INDICES INTO SHELL AND BEAM ELEMENT
C      STIFFNESS MATRICES GIVEN LOCAL NODE NUMBER (SHELL = 1 TO 8,
C      BEAM = 1 TO 3) AND DOF NUMBER (1 TO 6)
C
COMMON/ASS/ISH(6,8),IBM(6,3)
DATA ISH/
1      1,2,3,0,0,0,
2      4,5,6,7,8,0,
3      9,10,11,0,0,0,
4      12,13,14,15,16,0,
5      17,18,19,0,0,0,
6      20,21,22,23,24,0,
7      25,26,27,0,0,0,
8      28,29,30,31,32,0/
DATA IBM/
1      1,2,3,12,13,14,
2      4,5,6,7,8,0,
3      9,10,11,15,16,17/
GO TO (10,20),NASTY
10 MANE(1)=NAME
MANE(2)=4H
CALL OU2INI(4,MANE,NDOF,NDOF,6,ICNTR,10000,NDOF,1,1,101)
GO TO 25
20 CALL OU4INI(4,NDOF,NDOF,6,1,NAME)
C
C      PRIMARY LOOP OVER NODES
C
25 NCUL=0
DO 210 I=1,NUMNP
C
C      CLEAR SPACE FOR 6 COLUMNS FOR THIS NODE (ALL ROWS)
C
CALL MOVER(0,0,BIG,1,6*NDOF)
IF (NSHELL.EQ.0) GO TO 100

```

```

C
C SEARCH FOR SHELL ELEMENTS THAT TOUCH THIS NODE
C
DO 70 K=1,NSHELL
  LNODZ=8
  IF (SHELL(9,K).EQ.0) LNODZ=6
  DO 60 L=1,LNODZ
    LNOD=SHELL(L+2,K)
    IF (LNOD.NE.1) GO TO 60

C
C FOUND ONE (CORNER L OF ELEMENT K), GET ELEMENT MATRIX
C
CALL READMS(IN,SMAT,32*32,ifix(SHELL(11,K)))

C
C LOOP ON 6 DOF FOR THAT NODE
C
DO 50 LL=1,6

C
C SEE IF THIS DOF IS ACTIVE FOR THIS NODE
C IF SO GET INDEX INTO ELEMENT MATRIX
C
  LLL=ISH(LL,L)
  IF (LLL.EQ.0) GO TO 50

C
C LOOP ON ALL 8 (6) NODES FOR COUPLING TERMS
C
DO 40 M=1,LNODZ
  DO 30 MM=1,6
    MMM=ISH(MM,M)
    IF (MMM.EQ.0) GO TO 30

C
C GET ROW NO. FOR MASTER MATRIX
C MNOD=SHELL(M+2,K)
C
  IBIG=6*(MNOD-1)+MM
  BIG(IBIG,LL)=BIG(IBIG,LL)+SMAT(LL,MMM)
30 CONTINUE
40 CONTINUE
50 CONTINUE
60 CONTINUE
70 CONTINUE

C
C SAME SONG + DANCE FOR BEAM ELEMENTS
C
IF (NBEAM.EQ.0) GO TO 180
100 DO 170 K=1,NBEAM
  DO 160 L=1,3
    LNOD=BEAM(L+2,K)
    IF (LNOD.NE.1) GO TO 160
    CALL READMS(IN,BMAT,17*17,ifix(BEAM(15,K))+NSHLTY)

```

```

        DO 150 LL=1,6
          LLL=IBM(LL,L)
          IF (LLL.EQ.0) GO TO 150
          DO 140 M=1,3
            DO 130 MM=1,6
              MMM=IBM(MM,M)
              IF (MMM.EQ.0) GO TO 130
              MNOD=BEAM(M+2,K)
              IBIG=6*(MNOD-1)+MM
              BIG(IBIG,LL)=BIG(IBIG,LL)+BMAT(LLL,MMM)
130          CONTINUE
140        CONTINUE
150      CONTINUE
160    CONTINUE
170  CONTINUE

C
C    TRANSFORM TO OUTPUT COORDINATES
C
180    IG=SEQ(I)
        CALL MOVER(0,0,TI1,1,9)
        CALL MOVER(1.,0,TI1,4,3)
        CALL MOVER(0,0,TI2,1,9)
        CALL MOVER(1.,0,TI2,4,3)
        ICORD=GRID(9,IG)
        IF (ICORD.GT.0) CALL MOVER(CORD(1,3,ICORD),1,TI1,1,9)
        IF (ICORD.GT.0.AND.IPSPC(IG).NE.6)
          CALL MOVER(CORD(1,3,ICORD),1,TI2,1,9)
        DO 190 J=1,NUMNP
          JG=SEQ(J)
          CALL MOVER(0,0,TJ1,1,9)
          CALL MOVER(1.,0,TJ1,4,3)
          CALL MOVER(0,0,TJ2,1,9)
          CALL MOVER(1.,0,TJ2,4,3)
          JCORD=GRID(9,JG)
          IF (JCORD.EQ.0.AND.JCORD.EQ.0) GO TO 190
          CALL MOVER(CORD(1,3,JCORD),1,TJ1,1,9)
          IF (IPSPC(JG).NE.6)
            CALL MOVER(CORD(1,3,JCORD),1,TJ2,1,9)
          CALL TRIPLE(TJ1,BIG(6*J-5,1),TI1,NDOF,3,Q)
          CALL TRIPLE(TJ1,BIG(6*J-5,4),TI2,NDOF,3,Q)
          CALL TRIPLE(TJ2,BIG(6*J-2,1),TI1,NDOF,3,Q)
          CALL TRIPLE(TJ2,BIG(6*J-2,4),TI2,NDOF,3,Q)
190    CONTINUE
C
C    WE NOW HAVE 6 COLS OF THE MASTER MATRIX ASSEMBLED
C    TIME TO WRITE THEM OUT
C
        DO 200 J=1,6
          NCOL=NCOL+1
          IF (NASTY.EQ.1)

```



```

      • CALL DD2COL(4,BIG(1,J),NDOF,NDOF,ICNTR,IRPOS,MANE)
        IF (NASTY.EQ.2)
      • CALL DD4COL(4,BIG(1,J),NDOF,NDOF,NCOL,NAME)
200  CONTINUE
210  CONTINUE
      RETURN
      END

```

```

SUBROUTINE ASSL(NAME,IN,BIG,SHELL,BEAM,CORD,GRID,SEQ,
    IPSPC,SLOAD,BLOAD)
    DIMENSION BIG(1),SHELL(12,1),BEAM(21,1),CORD(3,5,1),GRID(13,1),
    SEQ(1),IPSPC(1),SLOAD(32),BLOAD(32)
    INTEGER SEQ
    COMMON/N/NCORD,NUMNP,NSHELL,NPSHEL,NPSHELX,NSHLTY,NBEAM,NPBEAM,
    NBMTY,NMAT,NSPOINT,MAXGRD,NDOF
    COMMON/CONTRL/ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD,NCORE,
    WTMASS,NASTY
    LOGICAL ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD
    DIMENSION Q(3,3)
    DIMENSION MANE(2)
    COMMON A(1)
    COMMON/ASS/ISH(6,8),IBM(6,3)
    GO TO (4,5),NASTY
4  MANE(1)=NAME
    MANE(2)=0
    CALL OU2INI(4,MANE,NDOF,1,2,1CNTR,10000,NDOF,1,1,101)
    GO TO 8
5  CALL OU4INI(4,NDOF,1,2,1,NAME)
8  CALL MOVER(0,0,BIG,1,NDOF)
    DO 200 I=1,NUMNP
        IF (NSHELL.EQ.0) GO TO 50
        DO 40 K=1,NSHELL
            DO 30 L=1,8
                LNOD=SHELL(L+2,K)
                IF (GRID(11,LNOD).NE.1) GO TO 30
                CALL READMS(IN,SLOAD,32,IFIX(SHELL(11,K)))
                DO 10 LL=1,6
                    LLL=ISH(LL,L)
                    IF (LLL.EQ.0) GO TO 10
                    IBIG=6*(GRID(11,LNOD)-1)+LL
                    BIG(IBIG)=BIG(IBIG)+SLOAD(LL)
10                 CONTINUE
30            CONTINUE
40        CONTINUE
        IF (NBEAM.EQ.0) GO TO 150
50        DO 140 K=1,NBEAM
            DO 130 L=1,3
                LNOD=BEAM(L+2,K)
                IF (GRID(11,LNOD).NE.1) GO TO 130
                CALL READMS(IN,BLOAD,17,IFIX(BEAM(15,K))+NSHLTY)
                DO 110 LL=1,6
                    LLL=IBM(LL,L)
                    IF (LLL.EQ.0) GO TO 110
                    IBIG=6*(GRID(11,LNOD)-1)+LL
                    BIG(IBIG)=BIG(IBIG)+BLOAD(LL)
110                CONTINUE
130            CONTINUE
140        CONTINUE

```

```

150      IG=SEQ(1)
        ICORD=GRID(9,IG)
        CALL MUVER(0,0,T1,1,9)
        CALL MUVER(1,0,T1,4,3)
        CALL MUVER(0,0,T2,1,9)
        CALL MUVER(1,0,T2,4,3)
        IF (ICORD.EQ.0) GO TO 200
        CALL MUVER(CORD(1,3,ICORD),1,T1,1,9)
        IF (IPSPC(IG).NE.6)
            CALL MUVER(CORD(1,3,ICORD),1,T2,1,9)
        CALL DOUBLE(BIG(6*1-5),T1,NDOF,3,Q)
        IF (NASTY.EQ.2) CALL OU4COL(4,BIG,NDOF,1,1,NAME)
200      CONTINUE
        IF (NASTY.EQ.1) CALL OU2COL(4,BIG,NDOF,1,ICNTR,IRPOS,MANE)
        IF (NASTY.EQ.2) CALL OU4COL(4,BIG,NDOF,1,1,NAME)
        END

```

```

SUBROUTINE PARVEC
COMMON/CONTROL/ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD,NCORE,
      WIMASS,NASTY
LOGICAL ECHO,ERR,MATGEN,DYN,APPLD,PLOT,DBUG,BAD
COMMON/N/NCORD,NUMNP,NSHELL,NPSHEL,NPSHELX,NSHLTY,NBEAM,NPBEAM,
      NBMTY,NMAT,NSPOINT,MAXGRD,NDOF
DIMENSION NAME(2)
COMMON A(1)
DATA NAME/4HPARV,2HEC/
NBUF=NCORE
CALL MURCOR(NDOF)
CALL MOVER(0,0,A(NBUF),1,6*NUMNP)
CALL MOVER(1.0,0,A(NBUF+6*NUMNP),1,NSPOINT)
CALL OUZINI(4,NAME,1,NDOF,2,1CNTR,10000,NDOF,2,1,101)
CALL OUZCOL(4,A(NBUF),1,NDOF,1CNTR,IRPOS,NAME)
NCORE=NBUF
RETURN
END

```



```

SUBROUTINE UU4COL(NRU,A,NROWS,NCOLS,NCOLN,NAME)
C      UU4COL 030375-1525
C      REVISED FOR MSC LEVEL 20
      DIMENSION A(1)
      DO 20 I=1,NROWS
      IF(A(I).NE.0.)GO TO 30
20     CONTINUE
      IF (NCOLN.LT,NCOLS) RETURN
      GO TO 70
30     NSTART=1
      DO 60 I=NSTART,NROWS
      IF(A(I).NE.0.)NSTOP=I
60     CONTINUE
      NW=NSTOP-NSTART+1
      WRITE(NRU) NCOLN,NSTART,NW,(A(I),I=NSTART,NSTOP)
      IF(NCOLN.LT,NCOLS)RETURN
70     NDUM1=NCOLS+1
      NDUM2=1
      DUM=-69.
      WRITE(NRU) NDUM1,NDUM2,NDUM2,DUM
      PRINT 1000,NAME,NRU
1000  FORMAT(21H UU4COL***DATA BLOCK ,A10,
+17H WRITTEN TO UNIT ,I2)
      RETURN
      END

```

```

SUBROUTINE UU4INI(NRU,NROWS,NCOLS,NFORM,NTYPE,NAME)
C   UU4INI 030375-1520
C   REVISED FOR MSC LEVEL 26
      WRITE(NRU) NCOLS,NROWS,NFORM,NTYPE
      PRINT 100,NAME,NRU
100  FORMAT(32H UU4INI***HEADER FOR DATA BLOCK ,A10,
+17H WRITTEN TO UNIT ,I2)
      RETURN
      END

```

```

SUBROUTINE UU2COL(NRU,A,NCOLS,NROWS,ICNTR,IRPOS,MNAME)
UU2COL 092374-1605

C
C
C   WRITE REAL S.P. NON-ZERO VALUES OF A MATRIX COLUMN IN
C   OUTPUT2 FORMAT-MUST CALL UU2INI BEFORE THIS TO
C   WRITE HEADER ON OUTPUT2 MATRIX FILE
C
C   9/18/74... GET
C
C   DIMENSION A(1),ISAME(4),IRPOS(1),MNAME(2)
C
C   DATA ISAME/1,0,0,264201/,IEOR/131071/
C
C   ICNTR=ICNTR-1
C   IF(IABS(ICNTR+2).GT.NCOLS) GO TO 40
C   WRITE(NRU) NROWS
C   WRITE(NRU) (A(I),I=1,NROWS)
C   WRITE(NRU)ICNTR
30  IF(IABS(ICNTR+2).NE.NCOLS) RETURN
C   NWRDS=0
C   WRITE(NRU) NWRDS
C   PRINT 60,MNAME,NRU
60  FORMAT(22H UU2COL*** DATA BLOCK ,2A4,
+17H WRITTEN TO TAPE ,I2)
C   RETURN
40  PRINT 50,(MNAME(I),I=1,2)
50  FORMAT(51H UU2COL*** ATTEMPT TO WRITE MORE COLUMNS THAN EXIST,
+12H FOR MATRIX ,2A4)
C   END

```

```

SUBROUTINE UU2INI(NRU,MNAME,NCOLS,NROWS,IFORM,ICNTR,IDNSTY,MCOL,
+ITYPE,ISKIP,IPOS)
C      UU2INI 092374-1600
C
C      WRITE HEADER INFORMATION IN OUTPUT2 FORMAT
C      CALL UU2COL AFTER THIS TO OUTPUT2 A MATRIX
C      CALL UU2SET TO OUTPUT2 USET
C
C      9/18/74...GET
C      '
C      DIMENSION NAST(7),MNAME(2)
C
C      DATA NAST/4HNAS1,4HRAN ,4HFORT,4H TAP,4HE ID,4H CUD,4HE - /
C      DATA LABL/4HXXXX/
C
C      901 FORMAT(3(1X,12))
C
C      IF (ISKIP.NE.0) GO TO 10
C      REWIND NRU
C      NWRDS=3
C      WRITE(NRU)NWRDS
C      CALL DATE (DAT)
C      DECODE(9,901,DAT)IMO,IDA,1YR
C      WRITE(NRU)IMO,IDA,1YR
C      NWRDS=7
C      WRITE(NRU)NWRDS
C      WRITE(NRU)(NAST(I),I=1,7)
C      NWRDS=2
C      WRITE(NRU)NWRDS
C      WRITE(NRU)LABL,LABL
C      ICNTR=-1
C      WRITE(NRU)ICNTR
C      ICNTR=0
C      WRITE(NRU)ICNTR
10  NWRDS=2
C      WRITE(NRU)NWRDS
C      WRITE(NRU)(MNAME(I),I=1,2)
C      ICNTR=-1
C      WRITE(NRU)ICNTR
C      NWRDS=8
C      WRITE(NRU)NWRDS
C      WRITE(NRU)IPOS,NCOLS,NROWS,IFORM,ITYPE,MCOL,IDNSTY
C      ,1
C      ICNTR=-2
C      WRITE(NRU)ICNTR
C      PRINT 20,MNAME,NRU
20  FORMAT(33H UU2INI*** HEADER FOR DATA BLOCK ,2A4,
+17H WRITTEN TO TAPE ,12)
C      RETURN
C      END

```


CCCC

```

SUBROUTINE DOUBLE(X,A,N,M,Q)
DIMENSION X(N,M),A(M,M),Q(M,M)
C
C   COMPUTE MATRIX PRODUCT Q=XA, STORE Q IN X
C
DO 10 I=1,M
  DO 10 K=1,M
    Q(I,K)=0.
    DO 10 J=1,M
      Q(I,K)=Q(I,K)+X(I,J)*A(J,K)
10  DO 20 I=1,M
    DO 20 J=1,M
      X(I,J)=Q(I,J)
20  X(I,J)=Q(I,J)
RETURN
END

```

```

      SUBROUTINE HALOOF(NEL, AREA, ELXYZT, FRAM, POINT, SIDE, THIK,
      , WSHEL, XITA, LNODZ, LNODZ, LVABZ, NUDEL)
C
C*** TO CREATE SHAPE FUNCTION ARRAY WSHEL, FOR SEMILOOF SHELL ELEMENT.
C
C*** WRITTEN BY BRUCE IRONS, JULY 1972, WASHINGTON D.C.
C
      DIMENSION AREAV(3), FRAME(3,3), GENSID(6,4),
1     SHEAR(11,43), SIGT(3), SWOP(6), THIKDD(3,3), TRANS(2,2),
2     VLOOF(3,36), WCORN(10,3), WLOOF(10,3), XGAUS(4,4), XILOOF(9,4),
3     XLOCAL(2), XYZDD(3,3), XYZPRE(8,4)
      DIMENSION ELXYZT(9,4), FRAM(3,3), POINT(3), WSHEL(13,45),
      , XITA(2), LNODS(NUDEL,NEL)
      EQUIVALENCE (T11,TRANS(1,1)), (T12,TRANS(1,2)),
1     (T21,TRANS(2,1)), (T22,TRANS(2,2))
      DATA GENSID/1., -1., 0., 3*-.5, 0., 1., -1., 4*1., 0., -1.,
1     4*0., 1., 0., -1., 2*1./, XILOOF/.211324866, 2*.788675134,
2     .211324866, 2*0., .3333333333, 4*0., .211324866, 2*.788675134,
3     .211324866, .3333333333, 2*0., -.577350269, .577350269, 2*1.,
4     .577350269, -.577350269, 2*-1., 0., 2*-1., -.577350269,
5     .577350269, 2*1., .577350269, -.577350269, 0. /,
6     XGAUS/0., 4*.5, 0., 2*.5, -.577350269, 2*.577350269,
7     3*-.577350269, 2*.577350269/, XYZPRE/32*0.0/, NOZPRE/0/
C
C*** GENERATE NSTAGE TO DEFINE PATH THROUGH HALOOF.
C
      N     ERROR = 1
      IF(LNODZ.NE.6. AND .LNODZ.NE.8) GO TO 99
      NSTAGE = 4
      IF(LNODZ.NE.NOZPRE) NSTAGE = 2
      NOZPRE = LNODZ
      DO 2 LNOD = 1,LNODZ
      DO 2 NX = 1,4
      IF(ELXYZT(LNOD,NX).NE.XYZPRE(LNOD,NX)) NSTAGE = 2
2     CONTINUE
      IF(NSTAGE.EQ.4) GO TO 18
C
C*** INITIALIZATION FOR NEW ELEMENT, NSTAGE = 1. FIND CENTRE COORDINATES
C
      LIMZ = (3*LNODZ)/2 - 1
      LVABZ = 4*LNODZ
      LNODZA = LNODZ + 1
      LVABZA = LVABZ + 1
      LVABZZ = LVABZ + LIMZ
      DO 3 L = LNODZA,LIMZ
      DO 3 J = 1,LVABZZ
3     SHEAR(L,J) = 0.0
      DO 5 NX = 1,4
      GASH = 0.0
      LNODZ1 = LNODZ/2

```

```

      DO 4 KURN = 1, LNODZ1
      DO 4 K = 1,2
4     GASH = GASH
      1   + 8.0*ELXYZT(2*KURN+K-2,NX)/FLUAT(216*K-408-LNODZ*(21*K-41))
      5 ELXYZT(9,NX) = GASH
C
C*** DIAGNOSTICS FOR A NEW ELEMENT. RELATE COORDINATES TO CENTRE.
C
      DO 10 I = 1, LNODZ
      N   ERROR = 2
      IF(ELXYZT(1,4).LE.0.0) GO TO 99
      IF(1.EQ.LNODZ) GO TO 9
      JA = 1 + 1
      DO 8 J = JA, LNODZ
      N   ERROR = 3
      IA = IABS(LNODS(I,NEL))
      IB = IABS(LNODS(J,NEL))
      IF(IA .EQ. IB) GO TO 99
      DO 7 K = 1,3
      IF(ELXYZI(I,K).NE.ELXYZT(J,K)) GO TO 8
7     CONTINUE
      N   ERROR = 4
      GO TO 99
8     CONTINUE
9     DO 10 NX = 1,4
      IF(NX.NE.4) ELXYZT(1,NX) = ELXYZT(1,NX) - ELXYZT(9,NX)
10    XYZPRE(1,NX) = ELXYZT(1,NX)
C
C*** CREATE SWOP = 1.0 OR -1.0, TO IMPLEMENT SIGN CHANGES AT LOOF NODES.
C
C*** ALSO INTERPOLATE TO ESTIMATE NORMAL THICKNESSES AT LOOF NODES.
C
      VLOOF(1,LVABZA) = ELXYZT(9,4)
      DO 12 NSIDE = 1,6
12     SWOP(NSIDE) = 1.0
      LAST = LNODZ - 1
      DO 14 NEXT = 1, LNODZ, 2
      MID = LAST + 1
      IA = IABS(LNODS(NEXT,NEL))
      IB = IABS(LNODS(LAST,NEL))
      IF(IA .LT. IB) SWOP(MID/2) = -1.0
      VLOOF(1,4*LAST-3) = .455341801*ELXYZT(LAST,4)
      1   + .666666667*ELXYZT(MID,4) - .122008468*ELXYZT(NEXT,4)
      VLOOF(1,4*MID-3) = -.122008468*ELXYZT(LAST,4)
      1   + .666666667*ELXYZT(MID,4) + .455341801*ELXYZT(NEXT,4)
C
C*** ALSO CHECK THAT MIDSIDE NODES ARE REASONABLY CENTRAL.
C
      GASH = 0.0
      GISH = 0.0

```



```

      GUSH = 0.0
      DO 13 I = 1,5
        ELMID = ELXYZT(MID,I)
        GASH = GASH + (ELXYZT(NEXT,I)-ELMID)**2
        GISH = GISH + (ELXYZT(LAST,I)-ELMID)**2
      13 GUSH = GUSH + (ELXYZT(LAST,I)+ELXYZT(NEXT,I)-ELMID-ELMID)**2
      N      ERROR = 5
C      IF (ABS(GASH-GISH).GT.0.040*(GASH+GISH)) GO TO 99
      N      ERROR = 6
C      IF (GUSH.GT.0.25*(GASH+GISH)) GO TO 99
      14 LAST = NEXT
C      WRITE(6,602) SWOP, (VLOOF(1,I), I = 1, LVABZA, 4)
C602  FORMAT(/7H SWOP =,6F6.3//26H THICKNESSES AT LOOF NODES/1X,9F13.8)
C
C*** ORGANISE LOOP AROUND LOOF NODES, FOR NSTAGE = 2
C
C*** DO 76 NSTAGE = 2,4 (IN EFFECT)
      15 NLOOF = 0
      16 NLOOF = NLOOF + 1
C*** DO 67 NLOOF = 1, LNODZ+1 IF NSTAGE = 2,
C*** OR DO 67 NLOOF = 1, (3*LNODZ)/2 IF NSTAGE = 3.
      DO 17 I = 1,2
        IF(NSTAGE.EQ.2. OR .NLOOF.LE.LNODZ)
          1  XLOCAL(I) = X1LOOF(NLOOF, LNODZ+I-6)
C
C*** AND ALSO AROUND INTEGRATING POINTS IF NSTAGE = 3.
C
      IF(NSTAGE.EQ.3. AND .NLOOF.GT.LNODZ)
        1  XLOCAL(I) = XGAUS(NLOOF-LNODZ, LNODZ+I-6)
      17 CONTINUE
      GO TO 23
C
C*** OTHERWISE, ORGANISE SINGLE-SHOT OPTION, FOR NSTAGE = 4.
C
C*** TEST WHETHER INPUT POINT IS A LOOF NODE, PLUS OR MINUS 0.0001.
C
      18 DO 19 I = 1,2
        19 XLOCAL(I) = X1TA(I)
        NLOOF = LNODZA
        DO 22 MAYBE = 1, LNODZ
          DO 20 I = 1,2
            IF (ABS(XLOCAL(I)-X1LOOF(MAYBE, LNODZ+I-6)).GT.0.0001) GO TO 22
          20 CONTINUE
          NLOOF = MAYBE
        22 CONTINUE
        IF(NLOOF.LE.LNODZ) WRITE(6,604) NLOOF
      604  FORMAT(/36H INPUT POINT RECOGNISED AS LOOF NODE,I3)
C
C*** CREATE VALUES AND XI,ETA DERIVATIVES OF X,Y,Z IN XYZDD, T IN THIKDD
C

```

```

23 CONTINUE
606 FORMAT(/13H *** NSTAGE =,12,10H, NLOOF =,13,7H, X1 =,F12.8,
1 8H, ETA =,F12.8)
CALL SFR(XLOCAL, WCURN, WLOOF, NSTAGE, LNODZ)
K = 0
DO 27 I = 1,3
DO 26 J = 1,3
GASH = 0.0
DO 24 L = 1, LNODZ
24 GASH = GASH + WCURN(L+K,1)*ELXYZT(L,J)
XYZDD(J,1) = GASH
IF(NSTAGE.EQ.2) GO TO 26
GASH = 0.0
DO 25 L = 1, LNODZA
25 GASH = GASH + WLOOF(L+K,1)*VLOOF(J,4*L-1)
THIKDD(J,1) = GASH
26 CONTINUE
27 K = 1
C WRITE(6,608) XYZDD
608 FORMAT(/6H XYZDD/(1X,3F15.10))
C IF(NSTAGE.EQ.3) WRITE(6,610) THIKDD
610 FORMAT(/7H THIKDD/(1X,3F15.10))
C
C*** CREATE VECTOR AREA = VAREA, AT GIVEN POINT X1, ETA.
C
CALL VECTOR(XYZDD(1,2), XYZDD(1,3), AREAV(1))
CALL SCALAR(AREAV(1), AREAV(1), AREASQ)
N ERROR = 7
IF(AREASQ.EQ.0.0) GO TO 99
AREA = SQRT(AREASQ)
C WRITE(6,612) AREA, AREAV
612 FORMAT(/7H AREA =,F13.10,10X,13HAREA VECTOR =,3F13.10)
C
C*** NORMALISE VECTOR AREA INTO FRAME, COL.3, AS LOCAL UNIT NORMAL Z.
C
C*** COLUMN 2 OF FRAME BECOMES UNIT Y AROUND EDGE.
C
DO 30 I = 1,3
FRAME(I,3) = AREAV(I)/AREA
GASH = 0.0
DO 29 J = 1,2
29 GASH = GASH + GENSID((NLOOF+1)/2, LNODZ+J-6)*XYZDD(1,J+1)
30 FRAME(1,2) = GASH
C
C*** NORMALISE Y, AND IMPLEMENT SWOP BY REVERSING SIGN OF Y.
C
C*** PUT APPROXIMATE VECTOR THICKNESS ETC. INTO VLOOF, FOR NSTAGE = 2
C
N ERROR = 8
CALL SCALAR(FRAME(1,2), FRAME(1,2), SIDESQ)

```

```

      IF(SIDESQ.EQ.0.0) GO TO 99
      SIDE = SQRT(SIDESQ)
      DO 31 I = 1,3
      FRAME(1,2) = FRAME(1,2)*SWOP((NLOOF+1)/2)/SIDE
      IF(NSTAGE.NE.2) GO TO 31
      VLOOF(1,4*NLOOF-2) = FRAME(1,2)
      VLOOF(1,4*NLOOF-1) = FRAME(1,3)*VLOOF(1,4*NLOOF-3)
      VLOOF(1,4*NLOOF) = FRAME(1,3)
31 CONTINUE
C
C*** AND COLUMN 1 IS UNIT X, THE OUTWARD POINTING IN-PLANE NORMAL.
C
      CALL VECTOR(FRAME(1,2), FRAME(1,3), FRAME(1,1))
      IF (ABS(XITA(1)).EQ.1.,OR,ABS(XITA(2)).EQ.1.) RETURN
C
      WRITE(6,614) ((FRAME(J,I), I = 1,3), J = 1,3)
C614  FORMAT(/44H COLS OF FRAME ARE UNIT LOCAL CARTESIAN AXES//
C      1  (1X,3F13,10))
C
C*** CHECK THAT NORMALS ARE REASONABLY PARALLEL, WHILE NSTAGE = 2.
C
      IF(NSTAGE.GT.2) GO TO 35
      IF(NLOOF.EQ.1) GO TO 67
      KZ = 4*NLOOF-4
      DO 32 K = 4, KZ, 4
      CALL VECTOR(VLOOF(1,4*NLOOF), VLOOF(1,K), POINT(1))
      CALL SCALAR(POINT(1), POINT(1), COSSQ)
      N ERROR = 9
C
      IF(COSSQ.GT.0.75) GO TO 99
32 CONTINUE
C
C*** PLACE CONTRIBUTION OF CENTRAL NUDE IN VLOOF (NSTAGE = 2 ONLY)
C
C*** COMPLETE LOOP NLOOF = 1 TO LNODZ+1 FOR NSTAGE = 2.
C
      IF(NLOOF.LE.LNODZ) GO TO 67
      TH1K = VLOOF(1,LVABZA)
      DO 33 I = 1,3
      DO 33 J = 1,2
33  VLOOF(I,LVABZ+J) = FRAME(I,J)*TH1K
      GO TO 67
C
C*** CREATE THE 2X2 JACOBIAN MATRIX, AND INVERT IT, (NSTAGE = 3 OR 4)
C
35  DO 36 J = 1,2
      DO 36 I = 1,2
      CALL SCALAR(FRAME(1,I), XYZDD(1,J+1), TRANS(J,1))
36 CONTINUE
C
      WRITE(6,616) TRANS
C616  FORMAT(/6H TRANS/(1X,2F13,10))
      GASH = T11

```

```

      T11 = T22/AREA
      T22 = GASH/AREA
      T12 = -T12/AREA
      T21 = -T21/AREA
C      WRITE(6,616) TRANS
C
C*** TRANSFORM WCORN AND WLOOF INTO LOCAL X,Y DERIVATIVES.
C
      DO 41 N = 1,LNODZA
      DO 41 I = 1,2
      GASH = 0.0
      GISH = 0.0
      DO 40 J = 1,2
      GASH = GASH + TRANS(I,J)*WCORN(N+11,J)
40    GISH = GISH + TRANS(I,J)*WLOOF(N+11,J)
      WCORN(N,I+1) = GASH
41    WLOOF(N,I+1) = GISH
C      WRITE(6,618)
C618    FORMAT(/16H WCORN AND WLOOF/)
C      DO 42 I = 1,3
C42    WRITE(6,620) (WCORN(N,I), N = 1,LNODZA)
C620    FORMAT(1X,9F13.10)
C      DO 43 I = 1,3
C43    WRITE(6,620) (WLOOF(N,I), N = 1,LNODZA)
C
C*** PUT THICKNESS AND DERIVATIVES INTO LOCAL COORDINATE SYSTEM.
C
      DO 45 I = 1,3
      DO 44 J = 1,2
      POINT(J) = 0.0
      DO 44 K = 1,2
44    POINT(J) = POINT(J) + TRANS(J,K)*THIKDD(1,K+1)
      DO 45 J = 1,2
45    THIKDD(1,J+1) = POINT(J)
      DO 48 J = 1,3
      DO 47 I = 1,3
      CALL SCALAR(THIKDD(1,J), FRAME(1,I), POINT(I))
47    CONTINUE
      DO 48 I = 1,3
48    THIKDD(1,I) = POINT(I)
C      WRITE(6,622) THIKDD
C622    FORMAT(/17H THICKNESS VECTOR,3F12.7//17H X-DERIVATIVES      ,3F12.7//
C      1  17H Y-DERIVATIVES      ,3F12.7)
      THIK = THIKDD(3,1)
      N      ERROR = 10
      IF(THIK.LE.0.0) GO TO 99
C
C*** FIND THE CHAVGE IN LOCAL X,Y DERIVATIVES ACROSS THICKNESS OF SHELL.
C
      DO 57 LNOD = 1,LNODZA

```



```

      IF(NSTAGE.NE.4) GO TO 51
      DO 50 I = 2,3
      GASH = 0.0
      DO 49 J = 1,2
49  GASH = GASH + THIKDD(J,1)*WCORN(LNOD,J+1)
50  POINT(I) = GASH
C
C*** CREATE WSHEL = SHAPE FUNCTION ARRAY, DISPLACEMENT TERMS FIRST.
C
      51 KORN = (LNOD+1)/2
      DO 54 K = 1,3
      KOL = 2*KORN + 3*LNOD + K - 5
      IF(LNOD.GT.LNODZ) KOL = 5*LNODZ + 2 + K
      DO 53 N = 1,3
      FACT = FRAME(K,N)
      WSHEL(N,KOL) = WCORN(LNOD,1)*FACT
      IF(NSTAGE.EQ.4. AND .N.EQ.3) FACT = 0.0
      DO 53 ND = 2,3
53  WSHEL(N+N+ND,KOL) = WCORN(LNOD,ND)*FACT
      DO 54 N = 1,2
      DO 54 ND = 2,3
      WSHEL(N+7,KOL) = WSHEL(N+7,KOL)
      1   = THIKDD(ND-1,1)*WSHEL(N+N+ND,KOL)/THIK
      IF(NSTAGE.EQ.4) WSHEL(N+N+ND+6,KOL) = (POINT(ND)*FRAME(K,N)
      1   + THIKDD(3,ND)*WCORN(LNOD,N+1)*FRAME(K,3))/THIK
54  CONTINUE
C
C*** INTRODUCE ROTATION TERMS WITH BENDING ACTION INTO WSHEL.
C
      DO 57 L = 1,2
      KOL = (L-1)*4*LNODZ + (2-L)*6*KORN + LNOD
      IF(LNOD.GT.LNODZ) KOL = 5*LNODZ + 3 - L
      DO 56 N = 1,2
      CALL SCALAR(VLUOF(1,4*LNOD+L-4), FRAME(1,N), FACT)
      WSHEL(N+7,KOL) = FACT*WLUOF(LNOD,1)/THIK
      IF(NSTAGE.NE.4) GO TO 56
      DO 55 ND = 2,3
55  WSHEL(N+N+ND+6,KOL) = FACT*WLUOF(LNOD,ND)/THIK
56  CONTINUE
      DO 57 NROW = 1,7
57  WSHEL(NROW,KOL) = 0.0
C
C*** COMBINE LAST THREE COLUMNS OF WSHEL TO CREATE NORMAL DEFLECTION.
C
      IF(LNODZ.EQ.6) GO TO 61
      IZ = 3*NSTAGE + 1
      DO 60 I = 1,IZ
      GASH = 0.0
      DO 59 K = 1,3
59  GASH = GASH + WSHEL(1,42+K)*VLUOF(K,4*LNODZ+4)

```

```

      DO WSHEL(1,43) = GASH
C      WRITE(6,624) (N, (WSHEL(K,N), K = 1,13), N = 1, LVABZZ)
C624  FORMAT(/15H WSHEL ORIGINAL/9X,1HU,8X,1HV,8X,1HW,8X,2HUX,7X,2HUY,
C      1 7X,2HVX,7X,2HUY,7X,2HUY,7X,2HUY,7X,2HUY,6X,3HUXZ,6X,3HUYZ,6X,3HVXZ,6X,
C      2 3HVYZ/68X,18HUR UZ+WX OR VZ+WY/(14,13F9.5))
      61 IF(NSTAGE.EQ.4) GO TO 66
C
C*** CREATE ARRAY SHEAR, FOR INTRODUCING THE CONSTRAINTS (NSTAGE = 3)
C
      IF(NLOOF.GT.LNODZ) GO TO 63
      DO 62 I = 1, LVABZZ
      SHEAR(NLOOF,I) = WSHEL(9,I)
      SHEAR(11,I) = SHEAR(11,I) + WSHEL(8,I)*SIDE*THIK*SWOP((NLOOF+1)/2)
62  CONTINUE
      GO TO 67
63  DO 66 KOL = 1, LVABZZ
      DO 66 NXY = 1, 2
      GASH = SHEAR(LNODZ+NXY,KOL)
      DO 65 MXY = 1, 2
      CALL SCALAR(FRAME(1,MXY), VLOOF(1,4*LNODZ+NXY), FACT)
65  GASH = GASH + WSHEL(MXY+7,KOL)*AREA*THIK*FACT
66  SHEAR(LNODZ+NXY,KOL) = GASH
C
C*** COMPLETE LOOP AROUND LOOF NODES ETC. TO CREATE VLOOF OR SHEAR.
67  IF(NLOOF.LE.LNODZ. OR .
      1 (NSTAGE.EQ.3. AND .NLOOF.LT.(3*LNODZ)/2)) GO TO 16
      IF(NSTAGE.NE.2) GO TO 76
C
C*** CREATE PLUS-MINUS SUM OF THICKNESS VECTORS AT LOOF NODES (NSTAGE=2)
C
      DO 70 I = 1, 3
      GASH = 0.0
      DO 68 N = 3, LVABZ, 4
68  GASH = -GASH + VLOOF(1,N)
      SIGT(1) = GASH
C
C*** AND THE 3X3 MATRIX ASSOCIATED WITH IT, STORED IN XYZDD.
C
      DO 70 J = 1, 3
      GASH = 0.0
      IF(1.EQ.J) GASH = FLOAT(LNODZ)
      DO 69 N = 2, LVABZ, 4
69  GASH = GASH - VLOOF(1,N)*VLOOF(J,N)
70  XYZDD(I,J) = GASH
C      WRITE(6,626) SIGT, XYZDD
C626  FORMAT(/24H PLUS-MINUS ERROR VECTOR,3F12.8//
C      1 25H MATRIX FOR CORRECTING IT,3F12.8,2(/25X,3F12.8))
C
C*** GET THE ADJUGATE OF THIS 3X3 SYMMETRIC POSITIVE DEFINITE MATRIX.
C

```

```

      K = 3
      DO 71 I = 1,3
      CALL VECTOR(XYZDD(1,I), XYZDD(1,6-I-K), FRAME(1,K))
71 K = 1
      CALL SCALAR(XYZDD(1,1), FRAME(1,1), DETERM)
      DO 73 I = 1,3
      CALL SCALAR(FRAME(1,I), SIGT(1), PROD)
73 POINT(1) = PROD/DETERM
C      WRITE(6,628) FRAME, POINT
C628 FORMAT(/9H ADJUGATE,3(/1X,3F12.8)//10H SOLUTIONS,3F12.8)
C
C*** CORRECT VECTOR THICKNESSES IN VLOOP.
C
      FACT = 1.0
      DO 75 N = 2, LVABZ, 4
      FACT = -FACT
      CALL SCALAR(POINT(1), VLOOP(1,N), PROD)
      DO 74 I = 1,3
74 VLOOP(I,N+1) = VLOOP(1,N+1) - FACT*(POINT(I)-PROD*VLOOP(1,N))
C
C*** CREATE DIFFERENTIAL DISPLACEMENT VECTORS TO DEFINE ROTATIONS.
C
C*** THIS COMPLETES WORK FOR NSTAGE = 2.
C
      TFIRST = VLOOP(1,N-1)
      CALL VECTOR(VLOOP(1,N), VLOOP(1,N+1), VLOOP(1,N-1))
      DO 75 I = 1,3
75 VLOOP(I,N) = VLOOP(I,N)*TFIRST
      NZ = 4*LNODZA
C      WRITE(6,630) (N, (VLOOP(1,N), I = 1,3), N = 1,NZ)
C630 FORMAT(/6H VLOOP/(1X,13,6X,3F15.10))
      NSTAGE = 3
      GO TO 15
C
C*** SHEAR HAS BEEN CREATED IN NLOOP LOOP FOR NSTAGE = 3.
C
C*** CHOOSE PIVOT FOR REDUCING ARRAY SHEAR, AND DO ROW INTERCHANGE.
C
76 CONTINUE
C      WRITE(6,632)
C632 FORMAT(/6H SHEAR)
C      DO 77 N = 1, LVABZZ
C77 WRITE(6,634) N, (SHEAR(1,N), I = 1,LIMZ)
C634 FORMAT(14,11F10.6)
      DO 83 LIM = 1,LIMZ
      KP = LVABZ + LIM
      PIVOT = 0.0
      DO 79 L = LIM,LIMZ
      IF (ABS(PIVOT).GT.ABS(SHEAR(L,KP))) GO TO 79
      LBIG = L

```



```

      PIVOT = SHEAR(LBIG,KP)
79  CONTINUE
      DO 80 K = 1,LVABZZ
        CHANGE = SHEAR(LBIG,K)
        SHEAR(LBIG,K) = SHEAR(LIM,K)
80  SHEAR(LIM,K) = CHANGE/PIVOT
C
C*** REDUCE ARRAY SHEAR TO CREATE CONSTRAINT MATRIX,
C
C*** THIS COMPLETES WORK FOR NSTAGE = 3.
C
      DO 82 NROW = 1,LIMZ
        FACT = SHEAR(NROW,KP)
        IF(NROW.EQ.LIM. OR .FACT.EQ.0.0) GO TO 82
        DO 81 KOL = 1,LVABZZ
91  SHEAR(NROW,KOL) = SHEAR(NROW,KOL) - FACT*SHEAR(LIM,KOL)
82  CONTINUE
83  CONTINUE
C
      WRITE(6,636)
C636  FORMAT(/22H SHEAR AFTER REDUCTION)
C
      DO 85 N = 1,LVABZZ
C
95  WRITE(6,634) N, (SHEAR(I,N), I = 1,LIMZ)
        NSTAGE = 4
        GO TO 18
C
C*** USE ARRAY SHEAR TO CONSTRAIN WSHEL AT THE GIVEN POINT XI,ETA.
C
86  DO 88 I = 1,LVABZ
      DO 88 J = 1,13
        GASH = WSHEL(J,I)
        DO 87 K = 1,LIMZ
87  GASH = GASH - WSHEL(J,K+LVABZ)*SHEAR(K,I)
88  WSHEL(J,I) = GASH
C
      WRITE(6,638)
C640  FORMAT(14,13F9.5)
C638  FORMAT(/18H WSHEL CONSTRAINED)
C
C*** IMPLEMENT SWOP TO EXCHANGE TWO NORMAL SLOPES.
C
      DO 92 N = 8, LVABZ, 8
        IF(SWOP(N/8).EQ.1.0) GO TO 92
        DO 91 J = 1,13
          CHANGE = WSHEL(J,N)
          WSHEL(J,N) = WSHEL(J,N-1)
91  WSHEL(J,N-1) = CHANGE
92  CONTINUE
C
      WRITE(6,642)
C642  FORMAT(/14H WSHEL SWOPPED)
C
      DO 94 N = 1,LVABZ
C
94  WRITE(6,640) N, (WSHEL(J,N), J = 1,13)

```



```

C
C*** ASSEMBLE UXZ, UYZ, VXZ, VYZ TO CREATE WXX, WXY, WYY.
C
C      WRITE(6,644)
C644  FORMAT(/30H WSHEL WITH SECOND DERIVATIVES)
      DO 96 N = 1,LVABZ
      WSHEL(10,N) = -WSHEL(10,N)
      WSHEL(11,N) = -0.5*(WSHEL(11,N)+WSHEL(12,N))
      WSHEL(12,N) = -WSHEL(13,N)
C      WRITE(6,640) N, (WSHEL(J,N), J = 1,12)
      96 CONTINUE
C
C*** PUT POINT, FRAM IN COMMON, ALSO AREA, SIDE WITH INTEGRATING FACTORS
C
      AREA = AREA*(FLOAT(LNODZ)-5.6)/2.4
      SIDE = SIDE*FLOAT(LNODZ-4)/4.0
      DO 98 I = 1,3
      POINT(I) = XYZDD(1,I) + ELXYZI(9,I)
      DO 98 J = 1,3
      98 FRAM(1,J) = FRAME(1,J)
      RETURN
C
C*** WRITE DIAGNOSTIC ERROR MESSAGE.
C
      99 WRITE(6,699) NERROR
C699  FORMAT(/6H ERROR,15,18H IN SEGMENT HALDOOF)
      AREA=-69.
      END

```

```

      SUBROUTINE SFR(XLOCAL, WCURN, WLOOF, NSTAGE, LNODZ)
C
C*** SHAPE FUNCTION SUBROUTINE TO SERVE HALOOF.
C
      DIMENSION MD(4), TERMV(46), WCURN(10,3), WLOOF(10,3), XLOCAL(2)
      COMMON/COEF/COEF(247)
      DATA MD/8, 43, 90, 171/
      DATA TERMV /46*0./
C
C*** INITIALIZE AND PREPARE TO CALCULATE TERMV = POLYNOMIAL TERMS.
C
      XI = XLOCAL(1)
      ETA = XLOCAL(2)
C
      WRITE(6,604) XI, ETA
C604  FORMAT(/5H XI =,F15.10,6X,5HETA =,F15.10)
      IA = 2
C
C*** CREATE POLYNOMIAL TERMS AND XI, ETA DERIVATIVES.
C
      TERMV(1) = 0.0
      TERMV(2) = 1.0
      NZ = (LNODZ+NSTAGE-3)/2
      DO 6 N = 1, NZ
      IAN = IA + N
      N2 = N + 15
      N3 = N + 30
      DO 4 J = IA, IAN
      TERMV(J+N) = TERMV(J)*XI
      TERMV(J+N2) = TERMV(J)*FLOAT(IAN-J)
4 TERMV(J+N3) = TERMV(J-1)*FLOAT(J-IA)
      IA = IAN
6 TERMV(IA+N) = TERMV(IA-1)*ETA
C
C*** CREATE SPECIAL COMBINATIONS FOR LOOF NODES, ETC.
C
      DO 8 I = 8, 38, 15
      IF(LNODZ.EQ.6) TERMV(I) =
1  2.0*(TERMV(I)-TERMV(I+3)) + 3.0*(TERMV(I+1)-TERMV(I+2))
      IF(LNODZ.EQ.8) TERMV(I) = TERMV(I+2)
      IF(LNODZ.EQ.8) TERMV(I+2) = TERMV(I+6)
8 CONTINUE
C
C*** USE TERMV TO FIND WCURN AND WLOOF AND XI, ETA DERIVATIVES.
C
      NFOISZ = (NSTAGE+1)/2
      DO 18 NFOIS = 1, NFOISZ
      NZ = (3*LNODZ)/2 + NFOIS - 4
      IF(NZ.LE.10) GO TO 12
      NZ = 9
      DO 10 I = 10,40,15

```

```

10 TERMV(1) = TERMV(1+3) - TERMV(1+5)
12 K = 0
   DO 10 I = 1,3
   DO 16 N = 1,NZ
   GASH = 0.0
   MDEL = MD(LNODZ+NFOIS-6) + N*NZ - 15*I
   MA = 16*I-14
   MZ = 15*I+NZ-14
   DO 14 M = MA, MZ
14  GASH = GASH + TERMV(M)*COEF(M+MDEL)
   IF(NFOIS.EQ.1) WCORN(N+K,1) = GASH
   IF(NFOIS.EQ.2) WLOOF(N+K,1) = GASH
16  CONTINUE
18  K = 1
C    WRITE(6,606)
C606  FORMAT(/16H WCORN AND WLOOF/)
C    DO 22 I = 1,3
C 22  WRITE(6,608) (WCORN(N,I), N = 1,LNODZ+1)
C    DO 24 I = 1,3
C 24  WRITE(6,608) (WLOOF(N,I), N = 1,LNODZ+1)
C608  FORMAT(1X,9F13.10)
      RETURN
C
C*** ERROR DIAGNOSTICS, IF POINT LIES OUTSIDE ELEMENT.
C
   99 WRITE(6,610) XI, ETA
  610 FORMAT(/30H ERROR 11 IN SEGMENT SFR, XI =,F15.9,3X,SETA =,F15.9)
      STOP
      END

```

BLOCK DATA

```

C
C*** TO INITIALIZE COEFFICIENTS FOR CORNER-MIDSIDE AND LOOF VERSIONS
C
C*** OF QUADRATIC TRIANGLE AND QUADRILATERAL FOR SUBROUTINE SFR.
C
  DIMENSION COEFA(166), COEFB(81)
  COMMON/COEF/COEF(247)
  EQUIVALENCE (COEF(1), COEFA(1)), (COEF(167), COEFB(1))
  DATA COEFA/ 1., -3., -3., 2., 4., 2., 0., 4., 0., -4., -4., 0., 0.,
1 -1., 0., 2., 0., 0., 0., 0., 0., 0., 4., 0., 0., 0., -1., 0., 0.,
2 2., 0., 0., 4., 0., -4., -4., 0.910683603, 1.577350269,
3 -6.041451884, -6.196152423, 2.464101615, 8.928203230, 1.732050808,
4 -0.244016936, 0.422649731, 2.041451884, 4.196152423, -4.464101615,
5 -4.928203230, -1.732050808, 0.333333333, -1.422649731, -2.577350269,
6 -1.464101615, 5.000000000, 5.464101615, 1.732050808, 0.333333333,
7 -2.577350269, -1.422649731, 5.464101615, 5.000000000, -1.464101615,
8 -1.732050808, -0.244016936, 2.041451884, 0.422649731, -4.928203230,
9 -4.464101615, 4.196152423, 1.732050807, 0.910683602, -6.041451884,
1 1.577350269, 8.928203230, 2.464101615, -6.196152422, -1.732050807,
2 -1., 6., 6., -6., -6., -6., 0., -25., 0., 0., 25., 25., 25., -25., -25., 0.,
3 .5, 0., -.5, -.5, 0., 0., 0., 5, 0., -.25, 0., 0., 25., -25., 25., 25., -25., 0.,
4 .5, .5, 0., 0., 0., -.5, -.5, 0., 0., -.25, 0., 0., 25., 25., 25., 25., 25., 0.,
5 .5, 0., .5, -.5, 0., 0., 0., -.5, 0., -.25, 0., 0., 25., -25., 25., -25., 25., 0.,
6 .5, -.5, 0., 0., 0., -.5, .5, 0., 0., 1., 0., 0., -1., 0., -1., 0., 0., 1./
  DATA COEFB/ 0.000000000, 0.216506351, -0.375000000, -0.093750000,
1 0.216506351, 0.281250000, -0.649519053, 0.375000000, -0.324759526,
2 -0.000000000, -0.216506351, -0.375000000, -0.093750000, -0.216506351,
3 0.281250000, 0.649519053, 0.375000000, 0.324759526, 0.000000000,
4 0.375000000, 0.216506351, 0.281250000, -0.216506351, -0.093750000,
5 -0.375000000, -0.649519053, -0.324759526, 0.000000000, 0.375000000,
6 -0.216506351, 0.281250000, 0.216506351, -0.093750000, -0.375000000,
7 0.649519053, 0.324759526, -0.000000000, -0.216506351, 0.375000000,
8 -0.093750000, 0.216506351, 0.281250000, 0.649519053, -0.375000000,
9 -0.324759526, 0.000000000, 0.216506351, 0.375000000, -0.093750000,
1 -0.216506351, 0.281250000, -0.649519053, -0.375000000, 0.324759526,
2 -0.000000000, -0.375000000, -0.216506351, 0.281250000, -0.216506351,
3 -0.093750000, 0.375000000, 0.649519053, -0.324759526, -0.000000000,
4 -0.375000000, 0.216506351, 0.281250000, 0.216506351, -0.093750000,
5 .375, -.649519053, .324759526, 1., 0., 0., -.75, 0., -.75, 0., 0., 0./
  END

```



```

      SUBROUTINE VECTOR(U, V, W)
C
C*** TO COMPUTE VECTOR PRODUCT U*V INTO AREA W.
C
      DIMENSION U(3), V(3), W(3)
      K = 3
      DO 2 I = 1,3
      W(6-I-K) = U(K)*V(I) - U(I)*V(K)
2 K = I
      RETURN
      END

```

```

      SUBROUTINE SCALAR(U, V, PROD)
C
C*** TO COMPUTE SCALAR PRODUCT OF VECTORS U AND V.
C
      DIMENSION U(3), V(3)
      PROD = 0.0
      DO 2 I = 1,3
2 PROD = PROD + U(I)*V(I)
      RETURN
      END

```

```

      SUBROUTINE LOFBEM(NEL, LVABZ, SIDE, ELXYZI, FRAM, POINT, WBEAM,
      .      XI, NODEL, NDIM, LNODZ, LNODS)
C
C*** WRITTEN BY BRUCE IRONS AND FERNANDO ALBUQUERQUE
C*** SWANSEA, 1/7/1973
C
C
C*** TO ERATE SHAPE FUNCTION ARRAY WBEAM, FOR SEMILOOF BEAM ELEMENT
C
      LOGICAL KNOWNR
      DIMENSION LNODS(NODEL,NEL)
      DIMENSION ELXYZI(NODEL,NDIM)
      DIMENSION POINT(3), FRAM(3,3), WBEAM(13,45)
      DIMENSION FRAME(3,3), LNUPRE(6), SHEAR(10,21), VLOOF(3,6)
      DIMENSION WCORN(3,2), WLOOF(4,2), XILOOF(5)
      DATA XILOOF /-.577350269, .577350269,
      2 -.7745966692, .7745966692, 0.0/, LNUPRE/6*0/
C
C*** GENERATE NSTAGE
C
      N ERROR = 1
      IF(LVABZ .NE. 17 .AND. LVABZ .NE. 19) GO TO 99
      IF( LVABZ .EQ. 17) KONST = 4
      IF( LVABZ .EQ. 19) KONST = 2
      NSTAGE = 3
      DO 2 N = 1, LNODZ
      IF(LNODS(N,NEL) .NE. LNUPRE(N)) NSTAGE = 1
      2 LNUPRE(N) = LNODS(N,NEL)
      3 XLOCAL = XI
      GO TO ( 4, 99, 24) , NSTAGE
C
C***  INITIALIZATION AND DIAGNOSTICS
C
      4 DO 10 I=2,3
      JJ = I-1
      DO 8 J = 1, JJ
      N ERROR = 2
      IA = IABS(LNODS(I,NEL))
      IB = IABS(LNODS(J,NEL))
      IF(IA .EQ. IB) GO TO 99
      8 CONTINUE
      10 CONTINUE
C
C*** CHECK THAT MIDSIDE IS NEAR THE MIDLE AND THAT ELEMENT IS
C      NOT TOO CURVED
C
      GASH = 0.0
      GISH = 0.0
      GUSH = 0.0
      DO 12 I = 1,3

```

```

      GASH = GASH + (ELXYZT(3,1) - ELXYZT(2,1))**2
      GISH = GISH + (ELXYZT(2,1) - ELXYZT(1,1))**2
12  GUSH = GUSH + (ELXYZT(3,1) - 2.0*ELXYZT(2,1) + ELXYZT(1,1))**2
C
C***  CREATE UNIT BINORMAL Z OUT OF PLANE OF ELEMENT = COL. 3 OF FRAME
C
      KNOWFR=.FALSE.
      DO 13 I=1,3
        FRAME(I,1)=FRAM(1,1)
13  KNOWFR=KNOWFR.OR.FRAME(1,1).NE.0.
      DO 14 I = 1,3
        IF (.NOT.KNOWFR) FRAME(I,1)=ELXYZT(1,1)-2.*ELXYZT(2,1)+ELXYZT(3,1)
14  FRAME(1,2) = ELXYZT(3,1) - ELXYZT(1,1)
C
      DO 16 I = 1,4
        IF (.NOT.KNOWFR) FRAME(I-1,1)=FRAME(I-1,1)+1.0
        CALL VECTOR(FRAME(1,1), FRAME(1,2), FRAME(1,3))
        CALL SCALAR ( FRAME(1,3), FRAME(1,3), ZSQ)
        IF(ZSQ .GT. 0.0) GO TO 18
16  CONTINUE
      N  ERROR = 5
      GO TO 99
18  DO 20 I = 1,3
20  FRAME(1,3) = FRAME(1,3)/SQRT(ZSQ)
C
C***  ORGANIZE LOOP AND END NODES TO CREATE VLOOP , THEN SHEAR
C
      21 NLOOP = 0
      22 NLCOF = NLOOP + 1
C
C***  DO 56 NLOOP = 1,2 IF NSTAGE= 1, BUT NLOOP = 1,5 IF NSTAGE = 2
C
      XLOCAL = XILOOP ( NLOOP )
      24 CALL SFR1(XLOCAL, WCORN, WLOOP)
C
C***  COMPLETE FRAME X IS ALONG TANGENT TO BEAM
C
      DO 28 I = 1,3
        GASH = 0.0
        DO 26 N = 1,3
26  GASH = GASH + ELXYZT(N,1)*WCORN(N,2)
28  FRAME(I,1) = GASH
        CALL SCALAR ( FRAME(1,1), FRAME(1,1), SIDESQ)
        N  ERROR= 6
        IF(SIDESQ .EQ. 0.0) GO TO 99
        SIDE = SQRT(SIDESQ)
        DO 30 I = 1,3
30  FRAME(I,1) = FRAME(I,1)/SIDE
C
C***  AND UNIT Y, IN COL 2 OF FRAME, IS NORMAL TO BOTH

```



```

C      CALL VECTOR( FRAME(1,3), FRAME(1,1), FRAME(1,2))
C      WRITE(6,608) ((FRAME(J,1), I = 1,3), J = 1,3)
C 608  FORMAT(/44H CULS OF FRAME ARE UNIT LOCAL CARTESIAN AXES,
C      1// (1X,3F13.10))
C      GO TO ( 31, 34, 34), NSTAGE
C
C***  APPEND FRAME TO VLOOF , HENCE ENDING TASK WITH NSTAGE = 1
C
C      31 DO 32 I = 1,3
C          DO 32 J = 1,3
C      32  VLOOF(1,3*(NLOOF-1)+J) = FRAME(I,J)
C          GO TO 56
C
C***  TRANSFORM DERIVATIVES OF WCORN AND WLOOF IN XILOOF
C
C      34 DO 36 N = 1,3
C      36  WCORN(N,2) = WCORN(N,2)/SIDE
C      WRITE(6,612)
C 612  FORMAT(/16H WCORN AND WLOOF/)
C      DO 37 I = 1,2
C      37  WRITE(6,614) (WCORN(N,I), N = 1,3)
C 614  FORMAT(1X, 4F13.10)
C      DO 38 N = 1,4
C      38  WLOOF(N,2) = WLOOF(N,2)/SIDE
C      DO 39 I = 1,2
C      39  WRITE(6,614) (WLOOF(N,I), N = 1,4)
C
C***  INITIALIZE WBEAM TO ZERO
C
C      DO 40 I = 1,12
C      DO 40 J = 1,24
C      40  WBEAM (I,J) = 0.0
C
C***  INTRODUCE TERMS FOR DISPLACEMENT INTO COL 1 TO 11 OF WBEAM
C
C      DO 42 NOD = 1,3
C      DO 42 IUVW = 1,3
C      KOL = 3*(NOD-1) + IUVW
C      IF ( NOD .EQ. 3) KOL = KOL+2
C      WBEAM (IUVW,KOL) = WCORN(NOD,1)
C      WBEAM(4,KOL) = WCORN(NOD,2)*FRAME(IUVW,1)
C      DO 42 NSHEAR = 1,2
C      42  WBEAM(NSHEAR+10,KOL) = WCORN(NOD,2)*FRAME(IUVW,NSHEAR+1)
C
C***  INTRODUCE ROTATION TERMS INTO WBEAM
C
C      DO 46 NOD = 1,4
C      DO 46 IUVW = 1,3
C      KOL = 3*NOD+IUVW+8

```

```

      IF( NOD .GE. 3) KOL = KOL+1
      DO 44 IXYZ = 1,3
44  WBEAM(IXYZ+4,KOL) = WLOOF(NOD,1) * FRAME(IUVW,IXYZ)
      WBEAM(8,KOL) = WLOOF(NOD,2)*FRAME(IUVW,3)
      WBEAM(9,KOL) = -WLOOF(NOD,2)*FRAME(IUVW,2)
      WBEAM(10,KOL) = WLOOF(NOD,2)*FRAME(IUVW,1)
      WBEAM(11,KOL) = -WLOOF(NOD,1)*FRAME(IUVW,3)
46  WBEAM(12,KOL) = WLOOF(NOD,1)*FRAME(IUVW,2)
C
C***  CREATE LOCAL ROTATION TERMS IN BEAM FOR NSTAGE = 2
C
      DO 51 NOD = 3,4
      DO 51 I = 1,3
      IF (I .EQ. 1) KOL = NOD+4
      IF(I .GT. 1) KOL = 2*NOD+I+10
      DO 51 NROW = 1,12
      GASH = 0.0
      DO 50 J = 1,3
50  GASH = GASH + WBEAM(NROW,3*NOD+J+9)*VLOOF(J,3*(NOD-3)+1)
51  WBEAM(NROW,KOL) = GASH
C  WRITE(6,617) (N, (WBEAM(K,N), K= 1,12), N = 1,24)
C 617 FORMAT(/15H WBEAM ORIGINAL/9X,1HU,8X,1HV,8X,1HW,
C      1 8X,2HDS,7X,4HROT,5X,4HROTY,5X,4HROTZ,
C      2 3X,6HCURVXY,3X,6HCURVXZ,2X,7HTORCURV,4X,5HSHEXY,4X,5HSHEXZ ,
C      3 /(14,12F9.5))
      GO TO ( 99,53,70), NSTAGE
C
C***  INTRODUCE THE BENDING CURVATURES INTO SHEAR
C
      53 DO 54 KOL = 1,21
      DO 54 IYZ = 1,2
      54 SHEAR(2*NLOOF+IYZ-2,KOL) = WBEAM(10+IYZ,KOL)
C***  COMPLETE LOOP FOR NLOOF , WITH NSTAGE = 1 OR 2
C
      56 IF(NLOOF .LT. 2) GO TO 22
      NSTAGE = NSTAGE + 1
      IF(NSTAGE .EQ. 2) GO TO 21
C  WRITE(6,618)
C 618 FORMAT(14H INITIAL SHEAR)
C  DO 57 N = 1,21
C  57 WRITE(6,620) N, (SHEAR(1,N) , I = 1,8)
C 620 FORMAT(14 ,10F10.6)
C
C***  FINALIZE THE MATRIX SHEAR , IN FIRST FOUR ROWS
C
      DO 58 KOL = 1,12
      DO 58 IYZ = 1 , 2
      SHEAR(IYZ,KOL) = SHEAR(IYZ,KOL)
58  SHEAR(IYZ+2,KOL) = SHEAR(IYZ+2,KOL)
C  WRITE(6,619)

```

```

C 619 FORMAT(12H FINAL SHEAR)
C      DO 59 N = 1,21
C 59 WRITE(6,620) N, (SHEAR(1,N), I = 1,10)
C
C***  CHOOSE PIVOT FOR REDUCING ARRAY SHEAR AND DO ROW INTERCHANGE
C
      DO 66 LIM = 1,KONST
      KP = LVABZ + LIM
      PIVOT = 0.0
      DO 60 L = LIM,KONST
      IF(ABS(PIVOT) .GT. ABS(SHEAR(L,KP))) GO TO 60
      LBIG = L
      PIVOT = SHEAR(LBIG,KP)
60 CONTINUE
      DO 61 K = 1,21
      CHANGE = SHEAR(LBIG,K)
      SHEAR(LBIG,K) = SHEAR(LIM,K)
61 SHEAR(LIM,K) = CHANGE/PIVOT
C
C***  REDUCE ARRAY SHEAR TO CREATE CONSTRAINT MATRIX
C
      DO 64 NROW = 1, KONST
      FACT = SHEAR(NROW,KP)
      IF(NROW .EQ. LIM .OR. FACT .EQ. 0.0) GO TO 64
      DO 62 KOL = 1,21
62 SHEAR(NROW,KOL) = SHEAR(NROW,KOL) - FACT*SHEAR(LIM,KOL)
64 CONTINUE
66 CONTINUE
C      WRITE(6,621)
C 621 FORMAT(/22H SHEAR AFTER REDUCTION)
C      DO 67 N = 1,21
C 67 WRITE(6,623) N, (SHEAR(1,N), I = 1,KONST)
C 623 FORMAT(14,4F10.6)
C      GO TO 3
C
C***  USE ARRAY SHEAR TO CONSTRAIN WBEAM AT THE GIVEN POINT XI
C
70 DO 74 I = 1, LVABZ
      DO 74 J = 1,12
      GASH = WBEAM(J,I)
      DO 72 K = 1,KONST
72 GASH = GASH - WBEAM(J,K+LVABZ)*SHEAR(K,I)
74 WBEAM(J,I) = GASH
C      WRITE(6,622)
C 622 FORMAT(/18H WBEAM CONSTRAINED)
C
C***  INTRODUCE SIGN CONVENTION FOR LOOF ROTATIONS
C
      IA = IABS(LNODS(3,NEL))
      IB = IABS(LNODS(1,NEL))

```

```

      IF (IA .GT. 18) GO TO 77
      DO 75 J = 1,12
      CHANGE = WBEAM(J,7)
      WBEAM(J,7) = -WBEAM(J,8)
      75 WBEAM(J,8) = -CHANGE
C      WRITE(6,626)
C 626 FORMAT(/14H WBEAM SWAPPED)
C      DO 76 N = 1,LVABZ
C 76 WRITE (6,624) N, (WBEAM(J,N), J= 1,12)
C 624 FORMAT(14,12F9.5)
C
C***  COMPUTE POINT , AND PASS FRAME INTO FRAM IN COMMON
C
      77 DO 80 I = 1,3
      GASH = 0.0
      DO 78 J = 1,3
      GASH = GASH + WCORN(J,1) * ELXYZT(J,1)
      78 FRAM(I,J) = FRAME(I,J)
      80 POINT(I) = GASH
      RETURN
      99 WRITE (6,699) NERROR
      699 FORMAT(/7H1 ERROR,15,18H IN SEGMENT LOFBEM)
      SIDE=-69.
      END

```



```

SUBROUTINE GCBSP1( X1, P)
C
C***  ROUTINE GIVES QUADRATIC INTERPOLATION FUNCTIONS IN X1.
C
  DIMENSION P(3)
  G = X1
  P(1) = .5 * G * (-1. + G )
  P(2) = 1. - G * G
  P(3) = .5 * G * ( 1. + G )
  RETURN
END

```

```

      SUBROUTINE SFRI(XI, WCORN, WLOOF)
C
C*** SHAPE FUNCTION SUBROUTINE TO SERVE LUBEM
      DIMENSION COEF(56), TERMV(4), WCORN(3,2), WLOOF(4,2)
C
C
C*** DEFINE COEFFICIENTS FOR SHAPE FUNCTIONS AND DERIVATIVES
C
      DATA COEF / 0., -.5, .5, 0.,
11., 0., -1., 0.,
2 0., .5, .5, 0.,
3-.5, 1., 0., 0.,
40., -2., 0., 0.,
5.5, 1., 0., 0.,
6 -0.25, 0.25, 0.75, -0.75,
7 -0.25, -0.25, 0.75, 0.75,
8.75, -1.2990381, -.75, 1.299038,
9.75, 1.2990381, -.75, -1.2990381,
1 0.25, 1.5, -2.25, 0.0,
2 -0.25, 1.5, 2.25, 0.0,
3-1.2990381, -1.5, 3.8971143, 0.,
41.2990381, -1.5, -3.8971143, 0./ , TERMV(1)/1.0/
C
C*** DIAGNOSTICS AND CREATE POWERS OF XI = XLOCAL
C
C 600 FORMAT(/9H XLOCAL =, F12.6)
C WRITE(6,600) XI
      IF(XI.LT.-1.0 .OR. XI .GT. 1.0) GO TO 99
      DO 2 I = 1,3
        2 TERMV(I+1) = TERMV(I)*XI
C
C*** CREATE WCORN WITH XI DERIVATIVES
C
      DO 10 IDIF = 1,2
        DO 6 NOD = 1,3
          IDEL = 12*IDIF + 4*NOD - 16
          GASH = 0.0
          DO 4 I = 1,3
            4 GASH = GASH+TERMV(I)*COEF(I+IDEL)
          6 WCORN(NOD,IDIF) = GASH
C
C*** CREATE WLOOF WITH XI DERIVATIVES
C
      DO 10 NOD = 1,4
        IDEL = 16*IDIF+4*NOD+4
        GASH = 0.0
        DO 8 I = 1,4
          8 GASH = GASH + TERMV(I) *COEF(I+IDEL)
        10 WLOOF(NOD,IDIF) = GASH
      RETURN

```

```
C
C***  ERROR DIAGNOSTICS
C
    99 WRITE(6,602)
    602 FORMAT(34H1ERROR IN SFR, X1 IS OUTSIDE RANGE)
    END
```

B.4 POSTLOOF LISTING

```

*DECK POSTLOO
PROGRAM POSTLOO(UT2,LSTRES,INPUT,OUTPUT,TAPE3,
.   TAPE1=LSTRES,TAPE2=UT2,TAPE5=INPUT,TAPE6=OUTPUT)

C
C   PROGRAM TO RECOVER STRESSES FOR SEMI-LOOF SHELL AND BEAM
C   ELEMENTS USING NASTRAN. SEMI-LOOF PRE-PROCESSOR IS THE
C   COMPANION PROGRAM
C
COMMON/N/NUMNP,NSHELL,NBEAM,NSHLTY,NBMTY,NGAUS
C   COMMON A(1)
COMMON A(20000)
DIMENSION TITLE(8)
REWIND 1
REWIND 2
READ(1) TITLE,NUMNP,NSHELL,NSHLTY,NBEAM,NBMTY,NGAUS
CALL DATE(DA)
CALL TIME(TI)
NINDEX=1
NU=NINDEX+9*NSHELL+1
CALL IN4INI(2,NCOLS,NROWS,NFORM,NTYPE,0,1)
NLNODS=NU+NROWS
NNP=NLNODS+10*NSHELL
NSEL=NLNODS+NSHLTY
NLAST=NSEL+NSHELL
CALL RFL(NLAST)
CALL INIT(A(NINDEX),A(NLNODS),A(NNP),A(NSEL))
DO 20 I=1,NCOLS
CALL IN4COL(2,A(NU),NCOLS,NROWS,11,NFIN,NERR,0)
20 CALL STRESS(TITLE,DA,TI,1,A(NU),A(NLNODS),A(NNP))
END

*DECK INIT
SUBROUTINE INIT(INDEX,LNODS,NP)
DIMENSION INDEX(1),LNODS(10,1),NP(1)
COMMON/N/NUMNP,NSHELL,NBEAM,NSHLTY,NBMTY,NGAUS
COMMON/BLOCK/BLOCK(790)
1 CALL OPENMS(3,INDEX,9*NSHLTY+1,0)
READ(1) ((LNODS(J,I),J=1,10),I=1,NSHELL)
DO 10 I=1,NSHLTY
READ(1) LT,NP(LT)
NPP=NP(LT)
DO 10 J=1,NPP
READ(1) BLOCK
10 CALL WRITMS(3,BLOCK,790,9*LT+J-9)
READ(1) ((LNODS(J,1),J=1,10),I=1,NSHELL)
RETURN
END

```



```

*DECK STRESS
SUBROUTINE STRESS(TITLE,DATE,TIME,ICASE,U,LNODS,NP)
DIMENSION U(1),LNODS(10,1),NP(1)
DIMENSION TITLE(8)
COMMON/BLOCK/STRESM(6,32),WSHEL(13,45),POINT(3),FRAM(3,3),THIK
COMMON/N/NUMNP,NSHELL,NBEAM,NSHLTY,NBMTY
DIMENSION ELDISP(32),S(3,6),FF(6)
DIMENSION ISH(6,8)
DATA ISH/1,2,3,0,0,0,
.      4,5,6,7,8,0,
.      9,10,11,0,0,0,
.      12,13,14,15,16,0,
.      17,18,19,0,0,0,
.      20,21,22,23,24,0,
.      25,26,27,0,0,0,
.      28,29,30,31,32,0/
DATA PI/3.1415927/
IF (NSHELL,EQ.0) GO TO 300

C
C  SHELL ELEMENTS
C
  1J=0
  DO 200 I=1,NSHELL

C
C    GET ELEMENT DISPLACEMENTS
C
    DO 115 J=1,8
      L=LNODS(J+2,1)
      IF (L,EQ.0) GO TO 115
      DO 110 K=1,6
        MDUF=6*(L-1)+K
        LDUF=ISH(K,J)
        IF (LDUF,NE.0) ELDISP(LDUF)=U(MDUF)
110      CONTINUE
115    CONTINUE

C
C    LOOP ON GAUSS POINTS
C
    NGAUS=NP(LNODS(2,1))
    DO 170 J=1,NGAUS
      CALL READMS(3,STRESM,790,9*LNODS(2,1)+J-9)
      1J=1J+1

C
C    GET 6 FORCE RESULTANTS IN LOCAL COORD
C
    DO 120 K=1,6
120    CALL SCPRUD(32,6,1,STRESM(K,1),ELDISP,FF(K))

C
C    LOOP ON Z=H/2, 0, -H/2 FOR COMBINED BENDING-STRETCHING
C

```

```

      Z=THIK/2
      T3=THIK**3
      DO 140 K=1,3
C
C      GET SX, SY, TXY
C
      DO 130 L=1,3
130      S(K,L)=FF(L)/THIK+FF(L+3)*12*Z/T3
C
C      GET PRINCIPAL STRESSES
C
      SS=(S(K,1)+S(K,2))/2,
      TT=(S(K,1)-S(K,2))/2,
      RR=SQRT(TT*TT+S(K,3)**2)
      S(K,4)=SS+RR
      S(K,5)=SS-RR
      S(K,6)=0,
      IF (S(K,3).NE.0..OR.TT.NE.0.)
      S(K,6)=0.5*ATAN2(-S(K,3),TT)*180/PI
140      Z=Z-THIK/2
      IF (MOD(IJ,8).EQ.1)
      WRITE(6,150) TITLE,ICASE,DATE,TIME
150      FORMAT (26H1SEM1-LOOP STRESS RECOVERY//1X,8A10/
      T100,4HCASE,13,2(2X,A10)//
      8H ELEMENT,T12,8HLOCATION,T31,11HLOCAL X/Y/Z,T55,
      23H** FORCE RESULTANTS **,190,
      40H** STRESSES SX/SY/TXY/SMAX/SMIN/ANGLE **/
      T55,9HFX/FY/FXY,169,9HMX/MY/MXY,T91,3HTOP,T105,6HMIDDLE,
      T120,6HBOTTOM)
      WRITE(6,160) I,POINT(1),(FRAM(K,1),K=1,3),FF(1),FF(4),
      (S(K,1),K=1,3)
160      FORMAT (/15,110,2HX=,E10.4,126,3HX=(,3F7.4,1H),
      153,E12.6,168,E12.6,T85,3E15.6)
      WRITE(6,161) POINT(2),(FRAM(K,2),K=1,3),FF(2),FF(5),
      (S(K,2),K=1,3)
161      FORMAT (T10,2HY=,E10.4,126,3HY=(,3F7.4,1H),
      T53,E12.6,T68,E12.6,T85,3E15.6)
      WRITE(6,162) POINT(3),(FRAM(K,3),K=1,3),FF(3),FF(6),
      (S(K,3),K=1,3)
162      FORMAT (T10,2HZ=,E10.4,T26,3HZ=(,3F7.4,1H),
      T53,E12.6,168,E12.6,T85,3E15.6)
      WRITE(6,163) ((S(L,K),L=1,3),K=4,6)
163      FORMAT (2(T85,3E15.6/),T82,3F15.2)
170      CONTINUE
200      CONTINUE
300      RETURN
      END

```

W

APPENDIX C
DETAILED DESCRIPTION OF SOFTWARE USED IN
STATISTICAL ENERGY ANALYSIS EXPERIMENT

C.1 GENERAL

The experimental investigation of the wave transmission method for implementing an SEA model was carried out using a minicomputer-based data analysis system. The computer and peripheral A/D converters replace a rather large amount of special-purpose analog gear which would otherwise be required and at the same time allow much greater flexibility in data analysis. However, to fully explain the procedure, some description of the software is needed. This description, along with source code listings, is provided in this appendix. It should be noted that the group of programs, while quite easy to use, does not represent a finished product. It is presented simply as the least ambiguous way of describing the experimental data processing.

The programs are written in Time Series Language (TSL), a proprietary language developed by Time/Data Division of Genrad. A working familiarity with the language, with standard vendor-supplied subroutines and with the DEC RT-11 operating system, is assumed. As written, the programs run on a Time/Data Model TDA-25 system and require a cartridge disc peripheral (DEC model RK05 or equivalent) and Tektronix 4000 series CRT terminal.

In the next section a brief functional description of each subroutine is given. Following that, a list of the subroutine groups which make up each load module is given along with the required subroutine call formats. Finally, source listings are given for all TSL subroutines. No listing for vendor-supplied machine language routines are given since these are not normally available to the user.

C.2 SUBROUTINE DESCRIPTIONS

The total program package is too large to fit in memory at one time. It is divided into groups of subroutines which are loaded and run as needed for various tasks. Since the TSL interpreter does not allow automatic handling of overlays, the groups must be loaded under manual control from the keyboard. In this section each subroutine is described individually.

Subroutine Name: ADSET

Description: Vendor-supplied machine language driver for A/D converters. In the current application it is accessed only through other vendor-supplied TSL subroutines which provide a simple operator interface.

Subroutine Name(s): SETADS, SETUP, PARAM, IO, IOALL

Description: Standard subroutines from the vendor-supplied TSL applications library which set, retrieve, store, and list all control parameters for the A/D peripheral. They are described in the TSL user manual for TDA-25 systems.

Subroutine Name: RTIO

Description: Vendor-supplied machine language subroutine to allow storage and retrieval of data on the system disc. It is described in the vendors application note.

Subroutine Name: DISPLY

Description: Vendor-supplied machine language routine to provide graphical display of data. It is described in the TSL manual.

Subroutine Name: CPHASE

Description: Vendor-supplied machine language subroutine which performs linear phase correction on cross-spectral density functions. It is used to compensate for non-simultaneous

sampling of the multiplexed analog signals and is described in the TDA-25 system documentation.

Subroutine Name: TRAN1

Description: Computes mobility function from acceleration signal on channel 2 and force signal on channel 1. Uses setup conditions previously stored under subroutine SETUP. Prompts operator for channel scale factors EU_A and EU_B in engineering units per volt and returns the mobility function in consistent engineering units (CEU's). Simultaneously computes coherence and channel 1 autopower spectral density.

Subroutine Name: NORM1

Description: Called by TRAN1 to provide proper normalization of mobility ($EU_A\text{-SEC}/EU_B$), coherence (nondimensional), and channel 1 autopower ($EU_A^2/\text{Hz.}$).

Subroutine Name: CLF

Description: Computes coupling loss factor by the wave transmission formula:

$$\omega\eta_{AB} = \frac{1}{m_A} \frac{R_e(\bar{H}_B)}{|\bar{H}_A + \bar{H}_B + H_C|^2}$$

where

$$\left. \begin{array}{l} \bar{H}_A = H_A(f_c) \\ \bar{H}_B = H_B(f_c) \end{array} \right\} \begin{array}{l} \text{complex coupling point mobilities of bodies} \\ \text{A and B previously measured under TRAN1,} \\ \text{frequency-averaged under AVGRB, and stored} \\ \text{on disc.} \end{array}$$

m_A = mass of body A

$$H_C(f) = \frac{i 2\pi f}{k_C}$$

k_C = coupling stiffness

$$i = \sqrt{-1}$$

f_c = center frequency of averaging band

Subroutine Name: CLFD

Description: Computes coupling loss factor by the same wave transmission formula as does CLF except that a direct connection between bodies A and B is assumed. This implies $k_c = \infty$ and $H_c = 0$.

Subroutine Name: AVGRB

Description: Performs a running average on a real or complex TSL block. When used on a complex block, the real and imaginary parts are smoothed separately. Averaging window width is specified in the call but must be $1/2^k$ times the width of the block. The value deposited in the destination block at a particular discrete frequency f_c is an approximation to

$$\bar{H}(f_c) = \frac{1}{\Delta f} \int_{f_c - \frac{\Delta f}{2}}^{f_c + \frac{\Delta f}{2}} H(f) df$$

The approximation comes about from the fact that $H(f)$ is obtained by measurement using discrete Fourier transform methods and is therefore known only as a discrete frequency series approximation to the true continuous function. The definite integration is approximated by a finite sum over a portion of the series.

Subroutine Name: DISPL2

Description: Displays unaveraged and averaged functions on the same plot scale as dotted and solid lines, respectively.

Subroutine Name: EBEA

Description: Computes predicted ratio of smoothed energy spectral densities $\bar{E}_A(f_c)/\bar{E}_B(f)$ from data files for $\omega\eta_{AB}$ and $\omega\eta_B$ previously computed and stored on disc. Uses the formula

$$\frac{\bar{E}_B(f_c)}{\bar{E}_A(f_c)} = \frac{\omega\eta_{AB}}{\omega\eta_B + \omega \frac{N_A}{N_B} \eta_{AB}}$$

Subroutine Name: PTRAN

Description: Computes transmitted power from force and acceleration signals at coupling point. Requests scale factors EU_A and EU_B for channels A and B in engineering units per volt. Expects force signal on channel A and returns the spectral density of transmitted power in $(EU_A - EU_B - \text{sec/Hz.})$.

Subroutine Name: LVXPS

Description: Called by PTRAN to compute the complex cross power spectral density between velocity and force by using signals proportional to acceleration and force.

Subroutine Name: NORM3

Description: Called by LVXPS to normalize the discrete cross power spectrum to account for channel scale factors, record length, number of records per ensemble, and double-sided DFT.

Subroutine Name: ENERGY

Description: Computes and stores a file of velocity power spectral densities and mass weighting coefficients which are later used by routine EADD to compute the total energy of a vibrating structure. Uses signals proportional to acceleration and computes for two points at a time. Prompts operator for channel scale factors and accelerometer locations.

Subroutine Name: PSDV2

Description: Called repeatedly by ENERGY to compute velocity power spectral density.

Subroutine Name: NORM2

Description: Called by PSDV2 to normalize velocity PSD to single-sided $(EU - \text{SEC})^2/\text{Hz}$. EU represents a scale factor in acceleration engineering units per volt for channel A or B as previously input under prompt from ENERGY.

Subroutine Name: EADD

Description: Uses an RTIO file of velocity PSD's previously stored by ENERGY to compute the energy spectral density for a component. Mass weighting coefficients are stored as the zero frequency entries of the velocity PSD's.

C.3 CALLING FORMATS

In this section the call formats are given for the various subroutines. The sequence in which they are presented represents a typical acquisition and processing run as described in section 3.4.3. It will be assumed that the system of units being used is force, length, and time in lbf., inches, and seconds. Consistent units for acceleration and force scale factors are thus (in/sec²)/volt and lbf/volt. In the following, prompts from computer are underlined and operator inputs are in capital letters. It is assumed that step 1 of the procedure is complete and that the RT-11 monitor and TSL interpreter are loaded and running. Storage of data is assumed to be done manually under calls to the subroutines of RTIO except for the multiple velocity PSD's measured in step 3 of section 3.4.3 which are stored under program control. RTIO usage is a standard TSL operation and will not be described further. Optional subroutine arguments are indicated by brackets. RT-11 file names and other text strings are enclosed in single quotes. All functions except point velocity PSD's are most conveniently stored in individual RT-11 files.

The description begins at step 2 of section 3.4.3.

<u>CALL</u>	<u>EXPLANATION</u>
<u>>LOAD</u> 'PTRANL'	Loads all required subroutines for measuring transmitted power.
<u>>SETUP</u>	Initiates a question-and-answer session to establish data acquisition conditions.
<u>>PTRAN</u>	Starts program to measure the spectral density of transmitted power.

INPUT FORCE (CH A) CEU/V Operator connects force signal to channel A input and responds with scale factor and carriage return (CR).

INPUT ACCEL (CH B) CEU/V Operator responds similarly for acceleration channel.

The spectral density of transmitted power is then measured, placed in a real floating point block B8, displayed on the system CRT, and control is returned to the operator. He may inspect and store the function as desired.

>CLEAR

>LOAD 'NRGL' Load all subroutines needed to measure component energy.

>SETUP Establish A/D conditions.

>ENERGY 'E_i Filename', 'text', recnum, mass ['N']

'E _i Filename'	RT-11 file for storage of velocity PSD's at various points on the structure. Point number is used as file record number.
'text'	Descriptive text string for E _i file.
recnum	Total number of desired records (new file) or next record to be filled (existing file).
mass	Mass coefficient to be stored with each velocity PSD. This version of the software assumes the same mass is associated with each point.

'N'	Denotes new file. Omit to resume filling an existing file.
<u>CH A EU/VOLT</u> <u>CH B EU/VOLT</u>	Operator inputs accelerometer scale factors for channels A and B each terminated by (CR).
<u>CH A TO POINT 1</u> <u>CH B TO POINT 2</u>	Operator mounts accelerometers and initiates measurement with (CR). Machine measures velocity PSD's, displays them, stores them in appropriate file records if ok'ed by the operator, and issues prompts to move accelerometers to the next two points.
<u>>CLEAR</u>	
<u>>LOAD 'EADDL'</u>	Loads subroutines to compute smoothed spectral density of component energy from point velocity PSD's and mass coefficients stored under ENERGY.
<u>>EADD 'E_i Filename', firstrec, lastrec</u>	
'E _i Filename'	File of point velocity PSD's and mass coefficients previously created under ENERGY.
firstrec	Point number of first velocity PSD to be included in mass weighted sum.
lastrec	Point number of last velocity PSD to be included in mass weighted sum.

Subroutine EADD will compute the spectral density of the total component energy as a mass-weighted sum of the point velocity PSD's and place it in real floating point block B8. The function is next frequency averaged.

>AVGRB B8, B9, avgfct

B8	Source block to be frequency averaged, in this case B8.
B9	Destination block where frequency averaged function is placed.
avgfct	Desired ratio of maximum DFT analysis frequency to averaging bandwidth.

The smoothed and unsmoothed component energy function may now be inspected and stored as desired under manual control. This completes step 3 of section 3.4.3.

The structures are next uncoupled and the shaker and transducing equipment rearranged to measure the admittance at the coupling point to body B.

>CLEAR

>LOAD 'HZACQL' Loads all subroutines required for mobility measurement. To save core, AVGRB is not yet loaded.

>SETUP Establish A/D parameters.

>TRAN1 Initiates program to measure $H_B(f)$.

INPUT FORCE (CH A) CEU/V Operator connects force signal to channel A input, inputs scale factor and (CR).

INPUT ACCEL (CH B) CEU/V Operator responds similarly for acceleration channel.

The program performs the measurements and computations to obtain $H_B(f)$. It is placed in a complex floating point block B8, displayed, and control is returned to the operator. The operator will normally inspect it and use RTIO calls to store it on disc as appropriate. The shaker and transducing are moved to body A and the process repeated to obtain $H_A(f)$. To obtain smoothed versions of H_A and H_B , the subroutine AVGRB is loaded.

>BLKCLR

>LOAD'AVGRB'

The functions $H_A(f)$ and $H_B(f)$ are read into memory by manual calls to RTIO and smoothed versions $\bar{H}_A(f_c)$ and $\bar{H}_B(f_c)$ are computed using AVGRB just as was done for component energy functions. AVGRB will handle either real or complex blocks. $\bar{H}_A(f_c)$ and $\bar{H}_B(f_c)$ are manually stored as usual. This completes step 6 of section 3.4.3.

>CLEAR

>LOAD 'CLFL'

>CLF 'H_AFile', 'H_BFile', k_c , m_A

'H_AFile'

File where $\bar{H}_A(f_c)$ was stored previously.

'H_BFile' File where H_B(f_c) was stored previously.

k_c Coupling stiffness

m_A Mass of body A.

The power transfer coefficient $2\pi f_c \eta_{AB}$ is computed from $H_A(f_c)$, $\bar{H}_B(f_c)$, k_c and m_A by CLF using equation 3.56 of this report. This is the predicted value of normalized transmitted power $\langle \pi_{AB} \rangle_t / \langle E_A \rangle_t$.

The actual value of frequency-dependent internal loss factor $\eta_B(f_c)$ is formed under manual control from the keyboard. The smoothed spectral densities of transmitted power and component B energy are read into memory and their ratio is formed. This is $2\pi f_c \eta_B(f_c)$ which is then written to disc.

The predicted energy ratio is next computed using the predicted coupling loss factor and actual internal loss factor.

>CLEAR

>LOAD 'EBEAL' Loads all subroutines needed to compute $\langle E_B \rangle_t / \langle E_A \rangle_t$

EBEA 'ωη_{AB}File', 'ωη_BFile', A_A/A_B

'ωη_{AB}File' File where predicted weighted coupling loss factor has previously been stored.

'ωη_BFile' File where actual weighted internal loss factor of plate B has been previously stored.

A_A/A_B Area ratio of plates

The predicted value of $\langle E_B \rangle_t / \langle E_A \rangle_t$ is computed using eq. 3.8 of this report. It is placed in a real floating block B8, displayed, and control is returned to the operator. He may then inspect and store the function as desired.

The actual values of normalized transmitted power $\langle \pi_{AB} \rangle_t / \langle E_A \rangle_t$ and normalized receptor energy $\langle E_B \rangle_t / \langle E_A \rangle_t$ are computed under manual control from the keyboard. The quantity $\langle \pi_{AB} \rangle_t$ is simply the smoothed spectral density of transmitted power times the smoothing bandwidth. A similar interpretation holds for $\langle E_B \rangle_t$ and $\langle E_A \rangle_t$. Since the same smoothing band was used throughout, the desired ratio is easily found. The previously measured and smoothed spectral densities are read into memory and divided. The ratio functions are then written back onto disc.

Predicted and measured values for all comparison values are now available. The standard TSL plot routine was used for making most of the figures of Chapter 3.0.

C.4 PROGRAM LISTING

Subroutine tables for each load module and source listing for each TSL subroutine are as follows.

```

LOAD 'HZACQL'
: PSTAB
IOP      63
IOALL    460
SETUP     70
PARAM    116
SETADS    169
HLET      34
ADSET    1415
CPHASE    199
NORM1     57
DISPLY    1800
OPEN      816
CLOSE     32
WRITEH    168
READH     362
WRITEB    53
READB     61
WAITIO    29
SPFUN     40
STORE     26
TRAN1     247
TOTAL=6217
>

```

```

LOAD 'CLFL'
: PSTAB
OPEN      816
CLOSE     32
WRITEH    168
READH     362
WRITEB    53
READB     61
WAITIO    29
SPFUN     40
DISPLY    1800
CLF       175
AUGRB     196
DISPL2    122
TOTAL=3854
>

```

```

LOAD 'PTRANL'
: PSTAB
OPEN      816
CLOSE     32
WRITEH    168
READH     362
WRITEB    53
READB     61
WAITIO    29
SPFUN     40
PARAM     116
IOALL     460
SETUP     70
ADSET     1415
SETADS    169
HLET      34
DISPLY    1800
IOP       63
CPHASE    199
LUXPS     330
NORM3     33
PTRAN     108
STORE     26
TOTAL=6384
>

```

```

LOAD 'NRGL'
: PSTAB
OPEN      816
CLOSE     32
WRITEH    168
READH     362
WRITEB    53
READB     61
WAITIO    29
SPFUN     40
DISPLY    1800
ADSET     1415
SETADS    169
HLET      34
SETUP     70
IOALL     460
PARAM     116
IOP       63
ENERGY    541
PSDU2     168
NORM2     81
TOTAL=6478
>

```

```

LOAD 'EADDL'
: PSTAB
OPEN      816
CLOSE     32
WRITEH    168
READH     362
WRITEB    53
READB     61
WAITIO    29
SPFUN     40
DISPLY    1800
LABEL     77
AUGRB     154
STORE     26
DISPL2    53
EADD      183
TOTAL=3854
>

```

```

LOAD 'EBEAL'
: PSTAB
OPEN      816
CLOSE     32
WRITEH    168
READH     362
WRITEB    53
READB     61
WAITIO    29
SPFUN     40
DISPLY    1800
DISPL3    122
EBEA      88
TOTAL=3571
>

```



```

LIST IOP
CREATE IOP
10 PRINT P0,': ',';'
20 GOTO I3,30,100
30 INPUT 'KB',P1
40 ARG I0
50 IF I0,4,110,110,60
60 IF P1,P3,10,70,70
70 IF P1,P4,110,110,10
100 PRINT P2,'-',P1
110 RETURN
END
>

```

```

LIST IOALL
CREATE IOALL
10 IOP 'FRAME SIZE',I15,'I15',128,8192
20 STACK 100,115,0,1,3,150,165,.5,165,100,0,1,2,1,1,165,150
30 IOP 'CHANNEL CODE 1=A 2=B 3=AB',I14,'I14',1,3
40 IOP 'BANDWIDTH',R15,'R15',0,80000.
50 IF R15,25000,80,80,60
60 PRINT 'NO FILTERS'
70 IOP 'SAMPLE RATE',R14,'R14',1,160000.
80 IOP 'CHANNEL A LEVEL, VOLTS',R13,'R13',0,16
90 IOP 'CHANNEL B LEVEL, VOLTS',R12,'R12',0,16
100 IOP 'BUFFER CODE 0=SINGLE 1=DOUBLE',I13,'I13',0,1
110 IOP 'TRIGGER 0=OFF 1=FIRST 2=EACH FRAME',I12,'I12',0,2
120 IF I12,170,170,130
130 IOP 'TRIGGER SOURCE 0=A 1=B 2=EXT',I11,'I11',0,2
140 IOP 'TRIGGER SLOPE 0=POS 1=NEG 2=+OR- 3=DIG',I10,'I10',0,3
150 IOP 'TRIGGER LEVEL X',R11,'R11',-99,99
160 IOP 'TRANSIENT CAPTURE X (0=OFF)',R10,'R10',0,99
170 IOP 'COUPLING 0=AC 1=DC',I9,'I9',0,1
180 IOP 'AVERAGE REMOVAL 0=OFF 1=REMOVE',I8,'I8',0,1
190 IOP 'ZERO INSERTION 0=OFF 1=ON',I7,'I7',0,1
200 IOP 'ERROR TRAPS 0=ENABLE 1=DISABLE',I6,'I6',0,1
210 IOP 'FRAME COUNT',I5,'I5'
220 IOP 'HANNING WINDOW 0=OFF 1=ON',I4,'I4',0,1
230 RETURN
END
>

```

```

LIST SETADS
CREATE SETADS
10 BLKCLR
20 STACK 100,101,102
30 HLET I0,I14,'XXA B AB'
40 STACK 25000.,R15,31,250
50 ADSET 'SU',B0,I15,I0,R0,R13,'UB',R12
60 IF R15,25000,80,80,70
70 ADSET 'SF',R14,'FL',0
80 HLET I0,I9,'ACDC'
90 HLET I1,I8,'DDED'
100 HLET I2,I7,'DZEZ'
110 ADSET 'CL',I0,I1,I2
120 HLET I0,I6,'ETDT'
130 HLET I1,I13,'SBDB'
140 ADSET I0,I1
150 IF I12,210,210,160
160 HLET I0,I12,'XXTFTE'
170 PROD I1,I11,4
180 SUM I1,I1,I10
190 HLET I1,I1,'APANABADBPBNBBBDEPENEBED'
200 ADSET I0,I1,R11,'TC',R10
210 STACK 152,151,150
220 RETURN
END
>

```

```

LIST SETUP
CREATE SETUP
10 ERASE
20 ARG I3
30 GOTO I3,40,60
40 IOALL
50 RETURN
60 ARG 0,P0,I3
70 IF I3,8,80,100,80
80 PRINT 'BAD FILE SPEC'
90 RETURN
100 INPUT P0,I15,I14,I13,I12,I11,I10,I9,I8,I7,I6,I5,I4,R15,R14,R13,R-
12,R11,R10
110 END
120 RETURN
END
>

```

```

LIST PARAM
CREATE PARAM
10 ARG I3
20 GOTO I3,30,90
30 ERASE
40 PRINT 'CURRENT SETUP'
50 PRINT
60 LET I3,1
70 IOALL
80 RETURN
90 ARG 0,P0,I3
100 IF I3,8,110,130,110
110 PRINT 'BAD FILE SPEC'
120 RETURN
130 WRITE P0,I15,',',I14,',',I13,',',I12,',',I11,',',I10,',',I9
140 WRITE P0,I8,',',I7,',',I6,',',I5,',',I4,',',R15,',',R14,',',R13,-
',',R12,',',R11,',',R10
150 END
160 RETURN
END
>

```

```

LIST EBEA
CREATE EBEA
10 REMARK "ETAAB .CLF","ETAB .NAD",AA/AB
20 OPEN P0
30 READH
40 READH B8
50 READB B8
60 OPEN P1
70 READH
80 READH B9
90 READB B9
100 CLOSE
105 MOVE B8,B10
110 MLCONR P2,B10
120 ADD B9,B10
130 DIU B10,B8
140 DISPLY B8,'EX',0,1000,'L25','SC',1.,'G','R'
150 RETURN
END
>

```

```

)LIST EADD
CREATE EADD
5 REMARK "EIFILE",FIRSTREC,LASTREC
10 BLKCLR
20 OPEN 0,P0,'B'
30 READH 0,B7
40 READB 0,B7,'R',0
50 MOVE B7,B8
60 ZERO B8
100 FOR I0,P1,P2
110 READB 0,B7,'R',I0
120 LET R6,B7,0
130 LET B7,0,0.
140 MLCONR R6,B7
150 ADD B7,B8
160 NEXT I0
170 MOVE B8,B7
180 BIBSET B7,5,I15
190 BIBSET B7,6,R1
200 QUOT R1,R1,2.
210 QUOT R1,R1,I15
215 MLCONR R1,B7
220 INTG B7
225 DIF I15,I15,1
230 LET R2,B7,I15
250 PRINT 'INPUT AUG. FACTOR'
260 INPUT R14
265 DISPLY B8,'YLAB','NRG/HZ','R'
270 AUGRB B8,B9,R14
280 BEAMP 550,700
290 PRINT 'TOTAL ENERGY = ',R2,' C.E.U.'
300 RETURN
END
>

```

```

LIST NORM2
CREATE NORM2
10 REMARK GIVES S.S. EU**2/HZ
20 BIBSET B7,6,R7
30 BIBSET B7,5,I3
40 QUOT R7,R7,4.
50 QUOT R7,I3,R7
60 QUOT R7,R7,I5
70 PROD R14,R8,R8
80 PROD R14,R7,R14
90 MLCONR R14,B7
100 PROD R14,R9,R9
110 PROD R14,R7,R14
120 MLCONR R14,B8
130 PROD R7,I5,R7
140 QUOT R7,2.,R7
150 RETURN
END
>

```

```

LIST PSDU2
CREATE PSDU2
10 LET I14,3
20 SETADS
30 FOR I0,1,15
40 ADSET 'SA'
50 DFT B0,I4
60 DFT B1,I4
70 ASPEC B0,B7
80 ASPEC B1,B8
90 NEXT I0
100 NORM2
110 PROD R1,6.283,R7
120 MOVE B7,B10
130 STCONR R1,B10
140 LET B10,0,0.
150 INTG B10
160 MUL B10,B10
170 LET B10,0,1.
180 FOR I0,0,2
190 LET B7,I0,0.
200 LET B8,I0,0.
210 NEXT I0
220 DIU B10,B7
230 DIU B10,B8
240 DISPLY B7,'YLAB','S(F)','R'
250 DISPLY B7,'GLAB','VELOCITY PSD','R'
260 DISPLY B7,'M','EX',0,R15,'G','R'
270 DISPLY B8,'NG','M','EX',0,R15,'G','R'
280 RETURN
END
>

```

```

LIST ENERGY
CREATE ENERGY
3 REMARK 'EFILE','TEXT',MAX OR
4 REMARK NEXTREC,MASS,C'N']
10 BLKCLR
20 PRINT 'CH A EU/VOLT'
30 INPUT R8
40 PRINT 'CH B EU/VOLT'
50 INPUT R9
60 ARG I0
70 IF I0,4,250,250,90
80 REMARK INITIALIZE NEW FILE
90 SETADS
100 ADSET 'SA'
110 DFT B0
120 ASPEC B0,B7
130 ZERO B7
140 OPEN P0,'N','B'
145 WRITEB B7,P1
150 SUM I2,P2,1
160 FOR I1,1,I2
180 WRITEB B7
190 NEXT I1
200 CLOSE
210 LET I1,1
220 SUM I2,I1,1
230 GOTO 270
240 REMARK OLD FILE:SET POINTERS
250 LET I1,P2
260 SUM I2,I1,1
270 REMARK ACQUIRE DATA FOR PSDU FILES
280 LET R6,P3
290 PRINT 'CH.A TO POINT ',I1
300 PRINT 'CH.B TO POINT ',I2
310 HOLIN

```

```

320 PRINT '(CR) TO BEGIN'
330 HOLIN I3
340 IF I3,13,330,350,330
350 PRINT 'NOW ACQUIRING'
360 PSDU2
370 REMARK STORE MASS AS DC VALUE
380 LET B7,0,R6
390 LET B8,0,R6
400 OPEN P0,'B'
420 WRITEB B7,I0,'R',I1
440 WRITEB B8,I0,'R',I2
460 IF I0,0,470,490,470
470 PRINT 'WRITEB ERROR'
480 RETURN
490 PRINT 'WHAT NOW?'
500 PRINT '    PROCEED (PR)'
510 PRINT '    REPEAT (RE)'
520 PRINT '    PAUSE (PA)'
530 PRINT '    EXIT (EX)'
540 HINPUT I3,'PR','RE','PA','EX'
550 IF I3,1,560,610,560
560 IF I3,2,570,290,570
570 IF I3,3,580,650,580
580 IF I3,4,590,640,590
590 PRINT 'WHAT?'
600 GOTO 540
610 SUM I1,I1,2
620 SUM I2,I1,1
630 GOTO 290
640 RETURN
650 PRINT 'DISPLAY A OR B?'
660 HINPUT I3,'A','B'
670 IF I3,1,680,710,680
680 IF I3,2,690,730,690
690 PRINT 'WHAT?'

```

```

700 GOTO 660
710 DISPLY B7,'M','EX',0,R15,'G'
720 GOTO 490
730 DISPLY B8,'M','EX',0,R15,'G'
740 GOTO 490
750 RETURN
END
>

```



```

LIST PTRAN
CREATE PTRAN
10 PRINT 'PTRAN MEASURES POWER '
15 PRINT 'FLOW INTO BODY 2 FROM'
17 PRINT 'FORCE AND ACCELERATION'
18 PRINT 'AT THE COUPLING POINT'
30 PRINT
35 PRINT 'INPUT FORCE (CH A) CEU/V'
50 INPUT R8
60 PRINT 'INPUT ACCEL (CH B) CEU/V'
75 INPUT R9
80 LUXPS
90 RETURN
END
>

```

```

LIST NORM3
CREATE NORM3
10 BIBSET B7,6,R7
12 QUOT R0,I15,R7
14 QUOT R0,R0,I5
20 PROD R0,R0,R8
30 PROD R0,R0,R9
100 MLCONR R0,B7
110 RETURN
END
>

```

```

LIST LUXPS
CREATE LUXPS
10 REMARK COMPUTE L-A XPS
20 LET I14,3
30 SETADS
40 FOR I0,1,I5
50 ADSET 'SA'
60 DFT B0,I4
70 DFT B1,I4
80 CSPEC B0,B1,B7
90 NEXT I0
100 CPHASE B7,1
110 NORM3
120 REMARK CREATE OMEGA IN B10
130 IMAG B7,B8
140 MOVE B8,B10
150 BIBSET B7,6,R7
160 QUOT R1,I15,R7
170 QUOT R1,6.283,R1
180 STCONR R1,B10
190 LET B10,0,0.
200 INTG B10
210 LET B10,0,1.
220 REMARK COMPUTE TRANSMITTED POWER
230 REMARK AND ITS SPECTRAL DENSITY
240 DIV B10,B8
250 QUOT R2,R1,6.283
260 LET B8,0,0.
270 LET B8,1,0.
280 MLCONR 2.,B8
290 MOVE B8,B7
300 MLCONR R2,B7
310 INTG B7
320 BIBSET B7,5,I0
330 DIF I0,I0,1

```

```

340 LET R3,B7,I0
350 REMARK DISPLAY OUTPUT
360 DISPLY B8,'EX',0,R15,'YLAB',' ','G','R'
370 BEAMP 20,600
380 PRINT 'TRANSMITTED'
390 BEAMP 0,580
400 PRINT 'POWER SPECTRAL'
410 BEAMP 40,560
420 PRINT 'DENSITY'
430 BEAMP 412,740
440 PRINT 'TRANSMITTED POWER = ',R3,'C.E.U.'
450 BEAMP 650,710
460 PROD R4,R3,1000000.
470 PRINT ' = ',R4,'MICRO CEU'
480 RETURN
END
>

```

```

LIST DISPL2
CREATE DISPL2
10 REMARK XMAX,YMAX,IFLOG
20 ARG I7
30 IF I7,3,40,70,70
40 DISPLY B8,'P','EX',0,P0,'SC',P1,'G','R'
50 DISPLY B9,'C','EX',0,P0,'SC',P1,'NG','G','R'
60 GOTO 90
70 DISPLY B8,'P','L50','EX',0,P0,'SC',P1,'G','R'
80 DISPLY B9,'C','NG','L50','EX',0,P0,'SC',P1,'G','R'
90 BEAMP 550,740
100 PRINT 'AUGNG BW. = ',R1,' HZ.'
110 RETURN
END
>

```

```

LIST AUGRB
CREATE AUGRB
10 REMARK SRCBLK,DSTBLK,AUGFCT
20 MOVE P0,P1
30 ZERO P1
40 BIBSET P0,5,I15
50 BIBSET P0,6,R1
60 QUOT R1,R1,2.
70 QUOT R1,R1,P2
80 QUOT I14,I15,P2
90 DIF I13,I15,I14
100 DIF I13,I13,1
110 QUOT I12,I14,2
120 FOR I11,0,I13
130 BIBSET P0,4,I8
140 BLKDEF B0,I14,I8,P0,I11
150 MOVE B0,B1
160 INTG B1
170 DIF I10,I14,1
180 IF I8,1,190,190,240
190 LET R0,B1,I10
200 QUOT R0,R0,I14
210 SUM I9,I12,I11
220 LET P1,I9,R0
230 GOTO 280
240 LET C0,B1,I10
245 QUOT R2,1.,I14
250 STACK 202,300,9,350
260 SUM I9,I12,I11
270 LET P1,I9,C0
280 NEXT I11
290 DISPLY P0,'P','G','R'
300 DISPLY P1,'NG','G','R'
310 BEAMP 550,740
320 PRINT 'AUGNG BDWTH = ',R1,' HZ.'
330 RETURN
END
>

```

```

LIST CLF
CREATE CLF
5 REMARK "HAFILE","HBFILE",KC,MA
10 OPEN P0
20 READH
30 READH B10
40 READB B10
50 CLOSE
60 PRINT
70 OPEN P1
80 READH
90 READH B11
100 READB B11
110 CLOSE
120 ASPEC B10,B8
130 ZERO B8
140 ASPEC B10,B9
150 ZERO B9
160 ADD B11,B10
170 BIBSET B11,6,R7
180 BIBSET B11,5,I0
190 QUOT R1,R7,I0
200 PROD R1,R1,3.141
210 MOVE B11,B12
220 STCONR R1,B9
230 ZERO B8
240 COMPLX B9,B8,B12
250 MOVE B12,B11
260 INTG B11
270 SUB B12,B11
280 QUOT R2,1.,P2
290 MLCONR R2,B11
300 MOVE B11,B12
310 ADD B11,B10
320 ASPEC B10,B9
330 OPEN P1
340 READH B11
350 READB B11
360 CLOSE
370 REAL B11,B8
380 DIV B9,B8
390 QUOT R2,1.,P3
400 MLCONR R2,B8
420 RETURN
END
>

```

```

LIST TRAN1
CREATE TRAN1
10 LET I14,3
20 PRINT 'INPUT FORCE (CH A) CEU/V'
30 INPUT R8
40 PRINT 'INPUT ACCEL (CH B) CEU/V'
50 INPUT R9
60 SETADS
70 FOR I0,1,I5
80 ADSET 'SA'
90 DFT B0,I4
100 DFT B1,I4
110 CSPEC B0,B1,B7
120 ASPEC B0,B5
130 ASPEC B1,B6
140 NEXT I0
150 CPHASE B7,1
160 NORM1
170 MOVE B5,B1
180 MOVE B5,B0
190 BIBSET B0,6,R7
200 BIBSET B0,5,I0
210 QUOT R1,R7,I0
220 QUOT R1,R1,2
230 PROD R1,R1,6.283
240 STCONR R1,B0
250 STCONR R1,B1
260 INTG B0
270 SUB B1,B0
280 LET B0,0,R1
290 MLCONC (0.,-1.),B7
300 DIU B0,B7
310 DIU B0,B6
320 DIU B0,B6
330 FOR I0,0,2
340 LET B5,I0,B5,3
350 LET B6,I0,B6,3
360 LET B7,I0,B7,3
370 NEXT I0
380 BIBSET B7,5,I0
390 BLKDEF B8,I0,2,B0,0
400 MOVE B7,B8
410 DIU B5,B8
420 MOVE B7,B9
430 ASPEC B9,B9
440 DIU B5,B9
450 DIU B6,B9
460 DISPLY B8,'EX',0,R15,'G','R'
470 RETURN
END
>

```

```

LIST NORM1
CREATE NORM1
10 BIBSET B7,6,R7
20 QUOT R0,I15,R7
30 QUOT R0,R0,I5
40 MLCONR R8,B5
50 MLCONR R8,B5
60 MLCONR R9,B6
70 MLCONR R9,B6
80 MLCONR R8,B7
90 MLCONR R9,B7
100 MLCONR R0,B5
110 MLCONR R0,B6
120 MLCONR R0,B7
130 RETURN
END
>

```

```

LIST CLFD
CREATE CLFD
5 REMARK "HAFILE","HBFIL",MA
8 REMARK "DIRECT CONNECTION" VERSION OF CLF
10 OPEN P0
20 READH
30 READH B10
40 READB B10
50 CLOSE
70 OPEN P1
80 READH
90 READH B11
100 READB B11
110 CLOSE
160 ADD B11,B10
320 ASPEC B10,B9
330 OPEN P1
340 READH B11
350 READB B11
360 CLOSE
370 REAL B11,B8
380 DIU B9,B8
390 QUOT R2,1.,P2
400 MLCONR R2,B8
420 RETURN
END
>

```

APPENDIX D

FUSELAGE MODEL LISTING

FULL FUSELAGE FREE VIBRATION

ID FUS,MODE

\$NASTRAN MSC

\$LUUFFACT 0.2

\$LUUFDYN YES

\$LUUFECNO YES

SOL 25,0

CHKPNT YES

TIME 20

DIAG 8

CEND

TITLE=FULL FUSELAGE MODEL

SUBTITLE=FIRST MODE

ECHO=NONE

METHOD=2

DYNRED=2

DISP=ALL

BEGIN BULK

CORD2R	1	0	0.	0.	0.	0.	0.	1.+CC1
+CC1	1.	0.	0.					
GRID	100	1	28.200	4.000	0.000			
GRID	1100	1	28.200	-4.000	0.000			
GRID	101	1	27.400	5.800	0.000			
GRID	1101	1	27.400	-5.800	0.000			
GRID	102	1	26.600	7.097	0.000			
GRID	1102	1	26.600	-7.097	0.000			
GRID	103	1	24.40	9.90	0.0			
GRID	1103	1	24.40	-9.900	0.0			
GRID	104	1	22.200	11.341	0.000			
GRID	1104	1	22.200	-11.341	0.000			
GRID	105	1	18.200	13.381	0.000			
GRID	1105	1	18.200	-13.381	0.000			
GRID	106	1	14.200	15.011	0.000			
GRID	1106	1	14.200	-15.011	0.000			
GRID	107	1	10.100	14.443	0.000			
GRID	1107	1	10.100	-14.448	0.000			
GRID	108	1	6.000	14.188	0.000			
GRID	1108	1	6.000	-14.188	0.000			
GRID	109	1	1.500	11.750	0.000			
GRID	1109	1	1.500	-11.750	0.000			
GRID	110	1	-3.408	7.787	0.000			
GRID	1110	1	-3.408	-7.787	0.000			
GRID	111	1	-4.120	5.003	0.000			
GRID	1111	1	-4.120	-5.003	0.000			
GRID	112	1	-5.240	0.000	0.000			
GRID	113	1	26.600	7.710	6.250			
GRID	1113	1	26.600	-7.710	6.250			
GRID	114	1	22.200	12.420	6.250			
GRID	1114	1	22.200	-12.420	6.250			
GRID	115	1	14.200	15.860	6.250			

GRID	1115	1	14.200	-15.860	6.250
GRID	116	1	6.000	15.410	6.250
GRID	1116	1	6.000	-15.410	6.250
GRID	117	1	-3.204	10.170	6.250
GRID	1117	1	-3.204	-10.170	6.250
GRID	118	1	-6.250	0.000	6.250
GRID	119	1	28.400	4.200	6.250
GRID	1119	1	28.400	-4.200	6.250
GRID	120	1	-4.750	7.590	6.250
GRID	1120	1	-4.750	-7.590	6.250
GRID	121	1	28.560	0.000	0.000
GRID	122	1	28.500	2.400	0.000
GRID	1122	1	28.500	-2.400	0.000
GRID	200	1	28.600	4.393	12.500
GRID	1200	1	28.600	-4.393	12.500
GRID	201	1	27.800	6.800	12.500
GRID	1201	1	27.800	-6.800	12.500
GRID	202	1	26.600	8.313	12.500
GRID	1202	1	26.600	-8.313	12.500
GRID	203	1	24.400	11.600	12.500
GRID	1203	1	24.400	-11.600	12.500
GRID	204	1	22.200	13.500	12.500
GRID	1204	1	22.200	-13.500	12.500
GRID	205	1	19.440	15.110	12.500
GRID	1205	1	19.440	-15.110	12.500
GRID	206	1	14.200	16.700	12.500
GRID	1206	1	14.200	-16.700	12.500
GRID	207	1	10.100	16.750	12.500
GRID	1207	1	10.100	-16.750	12.500
GRID	208	1	6.000	16.639	12.500
GRID	1208	1	6.000	-16.639	12.500
GRID	209	1	1.500	15.098	12.500
GRID	1209	1	1.500	-15.098	12.500
GRID	210	1	-3.000	12.558	12.500
GRID	1210	1	-3.000	-12.558	12.500
GRID	211	1	-5.120	10.345	12.500
GRID	1211	1	-5.120	-10.345	12.500
GRID	212	1	-7.240	0.000	12.500
GRID	213	1	26.600	9.240	19.500
GRID	1213	1	26.600	-9.240	19.500
GRID	214	1	22.200	14.200	19.500
GRID	1214	1	22.200	-14.200	19.500
GRID	215	1	14.200	16.840	19.500
GRID	1215	1	14.200	-16.840	19.500
GRID	216	1	4.550	16.720	19.500
GRID	1216	1	4.550	-16.720	19.500
GRID	217	1	-3.750	13.280	19.500
GRID	1217	1	-3.750	-13.280	19.500
GRID	218	1	-8.630	0.000	19.500
GRID	219	1	28.760	5.120	19.500

GRID	1219	1	28,760	-5,120	19,500
GRID	220	1	-7,200	2,750	12,500
GRID	1220	1	-7,200	-2,750	12,500
GRID	221	1	-7,200	1,750	12,500
GRID	1221	1	-7,200	-1,750	12,500
GRID	222	1	-8,580	3,000	19,500
GRID	1222	1	-8,580	-3,000	19,500
GRID	223	1	29,200	0,000	12,500
GRID	224	1	29,100	2,600	12,500
GRID	1224	1	29,100	-2,600	12,500
GRID	226	1	-5,80	3,0	19,50
GRID	1226	1	-5,80	-3,000	19,50
GRID	227	1	-6,80	3,0	19,50
GRID	1227	1	-6,80	-3,000	19,50
GRID	300	1	28,920	5,840	26,500
GRID	1300	1	28,920	-5,840	26,500
GRID	301	1	28,000	8,300	26,500
GRID	1301	1	28,000	-8,300	26,500
GRID	302	1	26,600	10,176	26,500
GRID	1302	1	26,600	-10,176	26,500
GRID	303	1	24,400	12,826	26,500
GRID	1303	1	24,400	-12,826	26,500
GRID	304	1	22,200	14,892	26,500
GRID	1304	1	22,200	-14,892	26,500
GRID	305	1	18,200	16,651	26,500
GRID	1305	1	18,200	-16,651	26,500
GRID	306	1	14,200	16,975	26,500
GRID	1306	1	14,200	-16,975	26,500
GRID	307	1	10,100	17,009	26,500
GRID	1307	1	10,100	-17,009	26,500
GRID	309	1	3,100	16,800	26,500
GRID	1309	1	3,100	-16,800	26,500
GRID	310	1	0,000	15,750	26,500
GRID	1310	1	0,000	-15,750	26,500
GRID	311	1	-4,500	14,000	26,500
GRID	1311	1	-4,500	-14,000	26,500
GRID	312	1	-10,000	0,000	26,500
GRID	313	1	26,600	10,890	33,000
GRID	1313	1	26,600	-10,890	33,000
GRID	314	1	21,600	15,500	33,000
GRID	1314	1	21,600	-15,500	33,000
GRID	315	1	14,200	17,500	33,000
GRID	1315	1	14,200	-17,500	33,000
GRID	316	1	4,550	17,330	33,000
GRID	1316	1	4,550	-17,330	33,000
GRID	318	1	-11,000	0,000	33,000
GRID	319	1	29,490	5,950	33,000
GRID	1319	1	29,490	-5,950	33,000
GRID	320	1	-9,950	3,250	26,500
GRID	1320	1	-9,950	-3,250	26,500

GRID	321	1	-9.000	7.000	26.500
GRID	1321	1	-9.000	-7.000	26.500
GRID	322	1	0.000	0.000	26.500
GRID	323	1	0.500	3.250	26.500
GRID	1323	1	0.500	-3.250	26.500
GRID	324	1	-7.980	9.170	33.000
GRID	1324	1	-7.980	-9.170	33.000
GRID	325	1	29.990	0.000	26.500
GRID	326	1	29.900	3.500	26.500
GRID	1326	1	29.900	-3.500	26.500
GRID	327	1	-2.00	3.25	26.50
GRID	1327	1	-2.00	-3.250	26.50
GRID	328	1	-7.25	3.25	26.50
GRID	1328	1	-7.25	-3.250	26.50
GRID	330	1	0.250	1.630	26.500
GRID	1330	1	0.250	-1.630	26.500
GRID	331	1	1.800	10.030	26.500
GRID	1331	1	1.800	-10.030	26.500
GRID	332	1	-4.500	8.630	26.500
GRID	1332	1	-4.500	-8.630	26.500
GRID	333	1	-2.250	1.630	26.500
GRID	1333	1	-2.250	-1.630	26.500
GRID	334	1	-4.500	3.250	26.500
GRID	1334	1	-4.500	-3.250	26.500
GRID	335	1	-10.000	1.630	26.500
GRID	1335	1	-10.000	-1.630	26.500
GRID	336	1	-7.250	1.630	26.500
GRID	1336	1	-7.250	-1.630	26.500
GRID	337	1	-5.000	0.000	26.500
GRID	338	1	3.20	17.20	33.00
GRID	13381		3.20	-17.200	33.00
GRID	339	1	0.000	0.000	33.000
GRID	340	1	-5.250	14.550	33.000
GRID	1340	1	-5.250	-14.550	33.000
GRID	341	1	0.250	1.630	33.000
GRID	1341	1	0.250	-1.630	33.000
GRID	400	1	30.050	6.060	39.500
GRID	1400	1	30.050	-6.060	39.500
GRID	401	1	28.000	9.200	39.500
GRID	1401	1	28.000	-9.200	39.500
GRID	402	1	26.500	11.600	39.500
GRID	1402	1	26.500	-11.600	39.500
GRID	403	1	24.400	13.565	39.500
GRID	1403	1	24.400	-13.565	39.500
GRID	404	1	21.000	16.100	39.500
GRID	1404	1	21.000	-16.100	39.500
GRID	405	1	18.200	17.296	39.500
GRID	1405	1	18.200	-17.296	39.500
GRID	406	1	14.200	18.027	39.500
GRID	1406	1	14.200	-18.027	39.500

GRID	407	1	10,100	18,116	39,500
GRID	1407	1	10,100	-18,116	39,500
GRID	408	1	6,000	17,859	39,500
GRID	1408	1	6,000	-17,859	39,500
GRID	409	1	4,500	17,768	39,500
GRID	1409	1	4,500	-17,768	39,500
GRID	410	1	3,200	17,592	39,500
GRID	1410	1	3,200	-17,592	39,500
GRID	411	1	-1,380	16,940	39,500
GRID	1411	1	-1,380	-16,940	39,500
GRID	412	1	-12,000	0,000	39,500
GRID	413	1	26,600	12,550	46,880
GRID	1413	1	26,600	-12,550	46,880
GRID	414	1	21,600	16,650	46,880
GRID	1414	1	21,600	-16,650	46,880
GRID	415	1	14,700	18,760	46,880
GRID	1415	1	14,700	-18,760	46,880
GRID	417	1	-6,500	15,570	46,880
GRID	1417	1	-6,500	-15,570	46,880
GRID	418	1	-13,000	0,000	46,880
GRID	419	1	30,325	6,580	46,880
GRID	1419	1	30,325	-6,580	46,880
GRID	420	1	-10,330	10,870	39,500
GRID	1420	1	-10,330	-10,870	39,500
GRID	421	1	-6,000	15,090	39,500
GRID	1421	1	-6,000	-15,090	39,500
GRID	422	1	0,000	0,000	39,500
GRID	423	1	-10,000	11,270	46,880
GRID	1423	1	-10,000	-11,270	46,880
GRID	426	1	14,450	17,710	46,880
GRID	1426	1	14,450	-17,710	46,880
GRID	427	1	10,350	17,630	46,880
GRID	1427	1	10,350	-17,630	46,880
GRID	428	1	3,100	15,870	46,880
GRID	1428	1	3,100	-15,870	46,880
GRID	429	1	-6,500	14,050	46,880
GRID	1429	1	-6,500	-14,050	46,880
GRID	430	1	0,000	0,000	46,880
GRID	431	1	1,600	8,800	39,500
GRID	1431	1	1,600	-8,800	39,500
GRID	500	1	30,600	7,100	54,250
GRID	1500	1	30,600	-7,100	54,250
GRID	501	1	28,600	11,000	54,250
GRID	1501	1	28,600	-11,000	54,250
GRID	502	1	26,600	13,500	54,250
GRID	1502	1	26,600	-13,500	54,250
GRID	503	1	24,400	16,114	54,250
GRID	1503	1	24,400	-16,114	54,250
GRID	504	1	22,200	17,200	54,250
GRID	1504	1	22,200	-17,200	54,250

GRID	505	1	18.200	19.000	54.250
GRID	1505	1	18.200	-19.000	54.250
GRID	506	1	15.200	19.500	54.250
GRID	1506	1	15.200	-19.500	54.250
GRID	507	1	14.700	17.400	54.250
GRID	1507	1	14.700	-17.400	54.250
GRID	508	1	3.070	14.143	54.250
GRID	1508	1	3.070	-14.143	54.250
GRID	509	1	-7.000	13.000	54.250
GRID	1509	1	-7.000	-13.000	54.250
GRID	510	1	-7.000	16.040	54.250
GRID	1510	1	-7.000	-16.040	54.250
GRID	511	1	-11.900	12.830	54.250
GRID	1511	1	-11.900	-12.830	54.250
GRID	512	1	-14.000	0.000	54.250
GRID	513	1	27.45	12.95	61.57
GRID	15131		27.45	-12.950	61.57
GRID	514	1	22.20	17.21	61.57
GRID	15141		22.20	-17.210	61.57
GRID	515	1	15.215	19.505	61.570
GRID	1515	1	15.215	-19.505	61.570
GRID	516	1	-14.520	9.455	61.570
GRID	1516	1	-14.520	-9.455	61.570
GRID	517	1	-8.250	16.570	61.570
GRID	1517	1	-8.250	-16.570	61.570
GRID	518	1	-15.000	0.000	61.570
GRID	519	1	31.150	6.950	61.570
GRID	1519	1	31.150	-6.950	61.570
GRID	520	1	-14.000	7.450	54.250
GRID	1520	1	-14.000	-7.450	54.250
GRID	521	1	0.000	0.000	54.250
GRID	522	1	-14.000	3.730	54.250
GRID	1522	1	-14.000	-3.730	54.250
GRID	528	1	31.300	4.300	54.250
GRID	1528	1	31.300	-4.300	54.250
GRID	529	1	31.500	0.000	54.250
GRID	530	1	8.700	15.800	54.250
GRID	1530	1	8.700	-15.800	54.250
GRID	531	1	14.950	18.450	54.250
GRID	1531	1	14.950	-18.450	54.250
GRID	532	1	-2.000	13.570	54.250
GRID	1532	1	-2.000	-13.570	54.250
GRID	533	1	-7.000	14.500	54.250
GRID	1533	1	-7.000	-14.500	54.250
GRID	534	1	14.850	16.700	61.570
GRID	1534	1	14.850	-16.700	61.570
GRID	535	1	2.890	14.240	61.570
GRID	1535	1	2.890	-14.240	61.570
GRID	536	1	-8.250	13.130	61.570
GRID	1536	1	-8.250	-13.130	61.570

GRID	537	1	1.540	7.070	54,250
GRID	1537	1	1.540	-7.070	54,250
GRID	538	1	-3.500	6.500	54,250
GRID	1538	1	-3.500	-6.500	54,250
GRID	539	1	-10.500	10.230	54,250
GRID	1539	1	-10.500	-10.230	54,250
GRID	540	1	-7.000	0.000	54,250
GRID	541	1	0.000	0.000	61,570
GRID	600	1	31.700	6.800	68,880
GRID	1600	1	31.700	-6.800	68,880
GRID	601	1	30.500	9.500	68,880
GRID	1601	1	30.500	-9.500	68,880
GRID	602	1	28.300	12.400	68,880
GRID	1602	1	28.300	-12.400	68,880
GRID	603	1	25.250	14.760	68,880
GRID	1603	1	25.250	-14.760	68,880
GRID	604	1	22.200	17.220	68,880
GRID	1604	1	22.200	-17.220	68,880
GRID	605	1	19.000	19.000	68,880
GRID	1605	1	19.000	-19.000	68,880
GRID	606	1	15.230	19.510	68,880
GRID	1606	1	15.230	-19.510	68,880
GRID	607	1	15.000	16.000	68,880
GRID	1607	1	15.000	-16.000	68,880
GRID	608	1	2.700	14.331	68,880
GRID	1608	1	2.700	-14.331	68,880
GRID	609	1	-9.500	13.250	68,880
GRID	1609	1	-9.500	-13.250	68,880
GRID	610	1	-9.500	17.100	68,880
GRID	1610	1	-9.500	-17.100	68,880
GRID	611	1	-13.640	15.320	68,880
GRID	1611	1	-13.640	-15.320	68,880
GRID	612	1	-16.000	0.000	68,880
GRID	613	1	28.950	12.300	76,760
GRID	1613	1	28.950	-12.300	76,760
GRID	614	1	22.200	18.110	76,760
GRID	1614	1	22.200	-18.110	76,760
GRID	615	1	14.720	20.230	76,760
GRID	1615	1	14.720	-20.230	76,760
GRID	616	1	-10.000	17.550	76,760
GRID	1616	1	-10.000	-17.550	76,760
GRID	617	1	-16.270	13.730	76,760
GRID	1617	1	-16.270	-13.730	76,760
GRID	618	1	-16.750	0.000	76,760
GRID	619	1	31.700	7.000	76,760
GRID	1619	1	31.700	-7.000	76,760
GRID	620	1	-16.000	3.000	68,880
GRID	1620	1	-16.000	-3.000	68,880
GRID	621	1	-15.040	11.460	68,880
GRID	1621	1	-15.040	-11.460	68,880

GRID	622	1	0.000	0.000	68,880
GRID	623	1	-16,270	9,320	76,760
GRID	1623	1	-16,270	-9,320	76,760
GRID	624	1	32,450	0.000	68,880
GRID	625	1	32,400	4,000	68,880
GRID	1625	1	32,400	-4,000	68,880
GRID	626	1	9,000	15,200	68,880
GRID	1626	1	9,000	-15,200	68,880
GRID	627	1	15,120	17,760	68,880
GRID	1627	1	15,120	-17,760	68,880
GRID	628	1	-3,400	13,790	68,880
GRID	1628	1	-3,400	-13,790	68,880
GRID	629	1	-9,500	15,180	68,880
GRID	1629	1	-9,500	-15,180	68,880
GRID	630	1	1,350	7,170	68,880
GRID	1630	1	1,350	-7,170	68,880
GRID	631	1	-4,750	6,630	68,880
GRID	1631	1	-4,750	-6,630	68,880
GRID	632	1	-12,270	12,360	68,880
GRID	1632	1	-12,270	-12,360	68,880
GRID	633	1	-8,000	0.000	68,880
GRID	634	1	14,550	16,700	76,760
GRID	1634	1	14,550	-16,700	76,760
GRID	635	1	0,300	14,920	76,760
GRID	1635	1	0,300	-14,920	76,760
GRID	636	1	-10,000	13,880	76,760
GRID	1636	1	-10,000	-13,880	76,760
GRID	637	1	10,53	16,25	76,76
GRID	1637	1	10,53	-16,250	76,76
GRID	638	1	4,38	15,42	76,76
GRID	1638	1	4,38	-15,420	76,76
GRID	700	1	31,700	7,200	84,630
GRID	1700	1	31,700	-7,200	84,630
GRID	701	1	30,800	9,900	84,630
GRID	1701	1	30,800	-9,900	84,630
GRID	702	1	29,600	12,200	84,630
GRID	1702	1	29,600	-12,200	84,630
GRID	703	1	26,400	15,800	84,630
GRID	1703	1	26,400	-15,800	84,630
GRID	704	1	22,200	19,000	84,630
GRID	1704	1	22,200	-19,000	84,630
GRID	705	1	18,200	20,300	84,630
GRID	1705	1	18,200	-20,300	84,630
GRID	706	1	14,200	21,000	84,630
GRID	1706	1	14,200	-21,000	84,630
GRID	708	1	-2,10	17,60	84,63
GRID	1708	1	-2,10	-17,600	84,63
GRID	709	1	6,050	20,800	84,630
GRID	1709	1	6,050	-20,800	84,630
GRID	710	1	-2,100	19,700	84,630

GRID	1710	1	-2,100	-19,700	84,630
GRID	711	1	-10,500	14,500	84,630
GRID	1711	1	-10,500	-14,500	84,630
GRID	712	1	-10,500	18,000	84,630
GRID	1712	1	-10,500	-18,000	84,630
GRID	713	1	29,950	12,200	93,630
GRID	1713	1	29,950	-12,200	93,630
GRID	714	1	22,200	19,500	93,630
GRID	1714	1	22,200	-19,500	93,630
GRID	715	1	14,200	21,750	93,630
GRID	1715	1	14,200	-21,750	93,630
GRID	716	1	6,000	21,550	93,630
GRID	1716	1	6,000	-21,550	93,630
GRID	717	1	-1,050	20,550	93,630
GRID	1717	1	-1,050	-20,550	93,630
GRID	718	1	-11,650	18,090	93,630
GRID	1718	1	-11,650	-18,090	93,630
GRID	719	1	32,500	7,000	93,630
GRID	1719	1	32,500	-7,000	93,630
GRID	720	1	14,100	17,400	84,630
GRID	1720	1	14,100	-17,400	84,630
GRID	721	1	6,05	18,65	84,63
GRID	1721	1	6,05	-18,650	84,63
GRID	722	1	-2,100	15,500	84,630
GRID	1722	1	-2,100	-15,500	84,630
GRID	723	1	-17,500	0,000	84,630
GRID	724	1	-2,100	0,000	84,630
GRID	725	1	-13,000	17,500	84,630
GRID	1725	1	-13,000	-17,500	84,630
GRID	726	1	-17,500	16,000	84,630
GRID	1726	1	-17,500	-16,000	84,630
GRID	727	1	-17,50	9,50	84,63
GRID	1727	1	-17,50	-9,500	84,63
GRID	728	1	10,100	21,000	84,630
GRID	1728	1	10,100	-21,000	84,630
GRID	729	1	6,050	16,500	84,630
GRID	1729	1	6,050	-16,500	84,630
GRID	730	1	10,100	17,150	84,630
GRID	1730	1	10,100	-17,150	84,630
GRID	731	1	1,900	20,200	84,630
GRID	1731	1	1,900	-20,200	84,630
GRID	732	1	1,98	16,00	84,63
GRID	1732	1	1,98	-16,000	84,63
GRID	733	1	-5,300	19,100	84,630
GRID	1733	1	-5,300	-19,100	84,630
GRID	734	1	-17,50	1,50	84,63
GRID	1734	1	-17,50	-1,500	84,63
GRID	735	1	-17,50	3,00	84,63
GRID	1735	1	-17,50	-3,000	84,63
GRID	736	1	-2,10	1,50	84,63

GRID	1736	1	-2.10	-1.500	84.63
GRID	737	1	32.740	0.000	84.630
GRID	738	1	32.400	4.300	84.630
GRID	1738	1	32.400	-4.300	84.630
GRID	739	1	-12.20	1.50	84.63
GRID	1739	1	-12.20	-1.500	84.63
GRID	740	1	-14.1	0.00	84.63
GRID	741	1	-2.100	3.000	92.630
GRID	1741	1	-2.100	-3.000	92.630
GRID	742	1	-2.100	0.000	93.130
GRID	743	1	-2.100	15.500	88.130
GRID	1743	1	-2.100	-15.500	88.130
GRID	744	1	-17.500	0.000	89.130
GRID	745	1	-17.500	0.000	93.630
GRID	746	1	-17.50	6.00	93.63
GRID	1746	1	-17.50	-6.000	93.63
GRID	747	1	-17.50	3.0	89.13
GRID	17471		-17.50	-3.000	89.13
GRID	748	1	-17.500	16.700	89.130
GRID	1748	1	-17.500	-16.700	89.130
GRID	749	1	-17.500	16.750	93.630
GRID	1749	1	-17.500	-16.750	93.630
GRID	750	1	-11.080	18.050	89.130
GRID	1750	1	-11.080	-18.050	89.130
GRID	751	1	-6.30	15.00	84.63
GRID	1751	1	-6.30	-15.000	84.63
GRID	752	1	-10.5	16.25	84.63
GRID	1752	1	-10.5	-16.250	84.63
GRID	753	1	14.150	19.200	84.630
GRID	1753	1	14.150	-19.200	84.630
GRID	754	1	-2.10	3.00	84.63
GRID	1754	1	-2.10	-3.000	84.63
GRID	755	1	-2.10	9.25	84.63
GRID	1755	1	-2.10	-9.250	84.63
GRID	756	1	-17.50	3.0	93.63
GRID	17561		-17.50	-3.000	93.63
GRID	757	1	-17.500	13.160	93.630
GRID	1757	1	-17.500	-13.160	93.630
GRID	758	1	-2.100	7.300	90.380
GRID	1758	1	-2.100	-7.300	90.380
GRID	759	1	-2.100	7.240	87.510
GRID	1759	1	-2.100	-7.240	87.510
GRID	760	1	-2.100	0.000	88.880
GRID	761	1	-16.58	12.05	93.63
GRID	1761	1	-16.58	-12.050	93.63
GRID	763	1	-15.80	0.0	84.63
GRID	764	1	-6.880	18.900	89.130
GRID	1764	1	-6.880	-18.900	89.130
GRID	765	1	-5.830	19.750	98.130
GRID	1765	1	-5.830	-19.750	98.130

GRID	766	1	-14,450	17,420	93,630
GRID	1766	1	-14,450	-17,420	93,630
GRID	767	1	-12,230	18,130	98,130
GRID	1767	1	-12,230	-18,130	98,130
GRID	768	1	-17,130	17,130	98,130
GRID	1768	1	-17,130	-17,130	98,130
GRID	769	1	-10,50	8,75	84,63
GRID	1769	1	-10,50	-8,750	84,63
GRID	770	1	-10,50	3,00	84,63
GRID	1770	1	-10,50	-3,000	84,63
GRID	771	1	-14,00	3,00	84,63
GRID	1771	1	-14,00	-3,000	84,63
GRID	772	1	-6,30	3,00	84,63
GRID	1772	1	-6,30	-3,000	84,63
GRID	773	1	-14,0	15,25	84,63
GRID	1773	1	-14,0	-15,250	84,63
GRID	774	1	-7,05	0,00	84,63
GRID	800	1	33,300	6,800	102,630
GRID	1800	1	33,300	-6,800	102,630
GRID	801	1	32,000	9,800	102,630
GRID	1801	1	32,000	-9,800	102,630
GRID	802	1	30,300	12,300	102,630
GRID	1802	1	30,300	-12,300	102,630
GRID	803	1	26,730	16,440	102,630
GRID	1803	1	26,730	-16,440	102,630
GRID	804	1	22,200	20,000	102,630
GRID	1804	1	22,200	-20,000	102,630
GRID	805	1	18,200	21,500	102,630
GRID	1805	1	18,200	-21,500	102,630
GRID	806	1	14,200	22,500	102,630
GRID	1806	1	14,200	-22,500	102,630
GRID	807	1	10,100	22,500	102,630
GRID	1807	1	10,100	-22,500	102,630
GRID	808	1	6,000	22,300	102,630
GRID	1808	1	6,000	-22,300	102,630
GRID	809	1	3,000	21,881	102,630
GRID	1809	1	3,000	-21,881	102,630
GRID	810	1	0,	21,4	102,630
GRID	18101		0,	-21,400	102,630
GRID	811	1	-6,4	20,	102,63
GRID	18111		-6,4	-20,000	102,63
GRID	812	1	-12,800	18,171	102,630
GRID	1812	1	-12,800	-18,171	102,630
GRID	813	1	29,150	14,050	112,130
GRID	1813	1	29,150	-14,050	112,130
GRID	814	1	21,750	20,400	112,130
GRID	1814	1	21,750	-20,400	112,130
GRID	815	1	14,200	22,700	112,130
GRID	1815	1	14,200	-22,700	112,130
GRID	816	1	6,000	22,650	112,130

GRID	1816	1	6,000	-22,650	112,130
GRID	817	1	0,000	21,800	112,130
GRID	1817	1	0,000	-21,800	112,130
GRID	818	1	-11,400	19,320	112,130
GRID	1818	1	-11,400	-19,320	112,130
GRID	819	1	33,050	6,350	112,130
GRID	1819	1	33,050	-6,350	112,130
GRID	820	1	-17,000	17,500	102,630
GRID	1820	1	-17,000	-17,500	102,630
GRID	821	1	-15,000	17,900	102,630
GRID	1821	1	-15,000	-17,900	102,630
GRID	826	1	-16,00	12,80	107,63
GRID	1826	1	-16,00	-12,800	107,63
GRID	822	1	-13,650	18,890	112,130
GRID	1822	1	-13,650	-18,890	112,130
GRID	823	1	34,000	0,000	102,630
GRID	824	1	33,800	4,500	102,630
GRID	1824	1	33,800	-4,500	102,630
GRID	825	1	-15,750	18,550	112,130
GRID	1825	1	-15,750	-18,550	112,130
GRID	900	1	32,800	5,900	121,630
GRID	1900	1	32,800	-5,900	121,630
GRID	901	1	31,100	11,100	121,630
GRID	1901	1	31,100	-11,100	121,630
GRID	902	1	28,000	15,800	121,630
GRID	1902	1	28,000	-15,800	121,630
GRID	903	1	24,400	18,900	121,630
GRID	1903	1	24,400	-18,900	121,630
GRID	904	1	21,300	20,800	121,630
GRID	1904	1	21,300	-20,800	121,630
GRID	905	1	18,200	22,000	121,630
GRID	1905	1	18,200	-22,000	121,630
GRID	906	1	14,200	22,900	121,630
GRID	1906	1	14,200	-22,900	121,630
GRID	907	1	10,100	23,200	121,630
GRID	1907	1	10,100	-23,200	121,630
GRID	908	1	6,000	23,000	121,630
GRID	1908	1	6,000	-23,000	121,630
GRID	909	1	3,000	22,700	121,630
GRID	1909	1	3,000	-22,700	121,630
GRID	910	1	0,000	22,200	121,630
GRID	1910	1	0,000	-22,200	121,630
GRID	911	1	-5,000	21,714	121,630
GRID	1911	1	-5,000	-21,714	121,630
GRID	912	1	-10,000	20,459	121,630
GRID	1912	1	-10,000	-20,459	121,630
GRID	920	1	-14,500	19,600	121,630
GRID	1920	1	-14,500	-19,600	121,630
GRID	921	1	-12,300	20,000	121,630
GRID	1921	1	-12,300	-20,000	121,630

GRID	922	1	33.300	0,000	121,630				
GRID	923	1	33,000	3,600	121,630				
GRID	1923	1	33,000	-3,600	121,630				
GRID	150	1	37,70	3,43	12,50				
GRID	151	1	30,2	-3,43	12,50				
SEQGP	150	584	151	585					
CLOUF8	1	1	100	101	102	113	202	201+CC1	
+CC1	200	119							
CLOUF8	1001	1001	1100	1101	1102	1113	1202	1201+CC1001	
+CC1001	1200	1119							
CLOUF8	2	1	102	103	104	114	204	203+CC2	
+CC2	202	113							
CLOUF8	1002	1001	1102	1103	1104	1114	1204	1203+CC1002	
+CC1002	1202	1113							
CLOUF8	3	1	104	105	106	115	206	205+CC3	
+CC3	204	114							
CLOUF8	1003	1001	1104	1105	1106	1115	1206	1205+CC1003	
+CC1003	1204	1114							
CLOUF8	4	1	106	107	108	116	208	207+CC4	
+CC4	206	115							
CLOUF8	1004	1001	1106	1107	1108	1116	1208	1207+CC1004	
+CC1004	1206	1115							
CLOUF8	5	1	108	109	110	117	210	209+CC5	
+CC5	208	116							
CLOUF8	1005	1001	1108	1109	1110	1117	1210	1209+CC1005	
+CC1005	1208	1116							
CLOUF8	6	1	110	111	112	118	212	221+CC6	
+CC6	220	120							
CLOUF8	1006	1001	1110	1111	112	118	212	1221+CC1006	
+CC1006	1220	1120							
CLOUF8	7	4	200	201	202	213	302	301+CC7	
+CC7	300	219							
CLOUF8	1007	1004	1200	1201	1202	1213	1302	1301+CC1007	
+CC1007	1300	1219							
CLOUF8	8	2	202	203	204	214	304	303+CC8	
+CC8	302	213							
CLOUF8	1008	1002	1202	1203	1204	1214	1304	1303+CC1008	
+CC1008	1302	1213							
CLOUF8	9	2	204	205	206	215	306	305+CC9	
+CC9	304	214							
CLOUF8	1009	1002	1204	1205	1206	1215	1306	1305+CC1009	
+CC1009	1304	1214							
CLOUF8	10	2	206	207	208	216	309	307+CC10	
+CC10	306	215							
CLOUF8	1010	1002	1206	1207	1208	1216	1309	1307+CC1010	
+CC1010	1306	1215							
CLOUF8	11	2	208	209	210	217	311	310+CC11	
+CC11	309	216							
CLOUF8	1011	1002	1208	1209	1210	1217	1311	1310+CC1011	
+CC1011	1309	1216							

CUOFF8	12	3	210	211	220	222	320	321+CC12
+CC12	311	217						
CUOFF8	1012	1003	1210	1211	1220	1222	1320	1321+CC1012
+CC1012	1311	1217						
CUOFF8	13	4	300	301	302	313	402	401+CC13
+CC13	400	319						
CUOFF8	1013	1004	1300	1301	1302	1313	1402	1401+CC1013
+CC1013	1400	1319						
CUOFF8	14	5	302	303	304	314	404	403+CC14
+CC14	402	313						
CUOFF8	1014	1005	1302	1303	1304	1314	1404	1403+CC1014
+CC1014	1402	1313						
CUOFF8	15	5	304	305	306	315	406	405+CC15
+CC15	404	314						
CUOFF8	1015	1005	1304	1305	1306	1315	1406	1405+CC1015
+CC1015	1404	1314						
CUOFF8	16	5	306	307	309	316	408	407+CC16
+CC16	406	315						
CUOFF8	1016	1005	1306	1307	1309	1316	1408	1407+CC1016
+CC1016	1406	1315						
CUOFF8	19	5	320	335	312	318	412	420+CC19
+CC19	421	324						
CUOFF8	1019	1005	1320	1335	312	318	412	1420+CC1019
+CC1019	1421	1324						
CUOFF8	20	4	400	401	402	413	502	501+CC20
+CC20	500	419						
CUOFF8	1020	1004	1400	1401	1402	1413	1502	1501+CC1020
+CC1020	1500	1419						
CUOFF8	21	5	402	403	404	414	504	503+CC21
+CC21	502	413						
CUOFF8	1021	1005	1402	1403	1404	1414	1504	1503+CC1021
+CC1021	1502	1413						
CUOFF8	22	5	404	405	406	415	506	505+CC22
+CC22	504	414						
CUOFF8	1022	1005	1404	1405	1406	1415	1506	1505+CC1022
+CC1022	1504	1414						
CUOFF8	23	5	421	420	412	418	512	522+CC23
+CC23	520	423						
CUOFF8	1023	1005	1421	1420	412	418	512	1522+CC1023
+CC1023	1520	1423						
CUOFF8	24	4	500	501	502	513	602	601+CC24
+CC24	600	519						
CUOFF8	1024	1004	1500	1501	1502	1513	1602	1601+CC1024
+CC1024	1600	1519						
CUOFF8	25	7	502	503	504	514	604	603+CC25
+CC25	602	513						
CUOFF8	1025	1007	1502	1503	1504	1514	1604	1603+CC1025
+CC1025	1602	1513						
CUOFF8	26	7	504	505	506	515	606	605+CC26
+CC26	604	514						

CLUOF8	1026	1007	1504	1505	1506	1515	1606	1605+CC1026
+CC1026	1604	1514						
CLUOF8	27	11	510	511	520	516	621	611+CC27
+CC27	610	517						
CLUOF8	1027	1011	1510	1511	1520	1516	1621	1611+CC1027
+CC1027	1610	1517						
CLUOF8	28	4	600	601	602	613	702	701+CC28
+CC28	700	619						
CLUOF8	1028	1004	1600	1601	1602	1613	1702	1701+CC1028
+CC1028	1700	1619						
CLUOF8	29	7	602	603	604	614	704	703+CC29
+CC29	702	613						
CLUOF8	1029	1007	1602	1603	1604	1614	1704	1703+CC1029
+CC1029	1702	1613						
CLUOF8	30	7	604	605	606	615	706	705+CC30
+CC30	704	614						
CLUOF8	1030	1007	1604	1605	1606	1615	1706	1705+CC1030
+CC1030	1704	1614						
CLUOF8	31	2	610	611	621	617	726	725+CC31
+CC31	712	616						
CLUOF8	1031	1002	1610	1611	1621	1617	1726	1725+CC1031
+CC1031	1712	1616						
CLUOF8	32	2	621	620	612	618	723	734+CC32
+CC32	735	623						
CLUOF8	1032	1002	1621	1620	612	618	723	1734+CC1032
+CC1032	1735	1623						
CLUOF8	33	11	712	725	726	748	749	766+CC33
+CC33	718	750						
CLUOF8	1033	1011	1712	1725	1726	1748	1749	1766+CC1033
+CC1033	1718	1750						
CLUOF8	37	4	700	701	702	713	802	801+CC37
+CC37	800	719						
CLUOF8	1037	1004	1700	1701	1702	1713	1802	1801+CC1037
+CC1037	1800	1719						
CLUOF8	38	3	702	703	704	714	804	803+CC38
+CC38	802	713						
CLUOF8	1038	1003	1702	1703	1704	1714	1804	1803+CC1038
+CC1038	1802	1713						
CLUOF8	39	3	704	705	706	715	806	805+CC39
+CC39	804	714						
CLUOF8	1039	1003	1704	1705	1706	1715	1806	1805+CC1039
+CC1039	1804	1714						
CLUOF8	40	3	706	728	709	716	808	807+CC40
+CC40	806	715						
CLUOF8	1040	1003	1706	1728	1709	1716	1808	1807+CC1040
+CC1040	1806	1715						
CLUOF8	42	3	709	731	710	717	810	809+CC42
+CC42	808	716						
CLUOF8	1042	1003	1709	1731	1710	1717	1810	1809+CC1042
+CC1042	1808	1716						

CLOUF8	43	11	718	760	749	768	820	821+CC43
+CC43	812	767						
CLOUF8	1043	1011	1718	1766	1749	1768	1820	1821+CC1043
+CC1043	1812	1767						
CLOUF8	47	28	724	736	754	772	770	739+CC47
+CC47	740	774						
CLOUF8	1047	1028	724	1736	1754	1772	1770	1739+CC1047
+CC1047	740	774						
CLOUF8	48	19	754	755	722	751	711	769+CC48
+CC48	770	772						
CLOUF8	1048	1019	1754	1755	1722	1751	1711	1769+CC1048
+CC1048	1770	1772						
CLOUF8	49	3	220	221	212	218	312	335+CC49
+CC49	320	222						
CLOUF8	1049	1003	1220	1221	212	218	312	1335+CC1049
+CC1049	1320	1222						
CLOUF8	51	4	800	801	802	813	902	901+CC51
+CC51	900	819						
CLOUF8	1051	1004	1800	1801	1802	1813	1902	1901+CC1051
+CC1051	1900	1819						
CLOUF8	52	11	520	522	512	518	612	620+CC52
+CC52	621	516						
CLOUF8	1052	1011	1520	1522	512	518	612	1620+CC1052
+CC1052	1621	1516						
CLOUF8	53	19	722	708	710	733	712	752+CC53
+CC53	711	751						
CLOUF8	1053	1019	1722	1708	1710	1733	1712	1752+CC1053
+CC1053	1711	1751						
CLOUF8	54	10	802	803	804	814	904	903+CC54
+CC54	902	813						
CLOUF8	1054	1010	1802	1803	1804	1814	1904	1903+CC1054
+CC1054	1902	1813						
CLOUF8	55	10	804	805	806	815	906	905+CC55
+CC55	904	814						
CLOUF8	1055	1010	1804	1805	1806	1815	1906	1905+CC1055
+CC1055	1904	1814						
CLOUF8	56	10	806	807	808	816	908	907+CC56
+CC56	906	815						
CLOUF8	1056	1010	1806	1807	1808	1816	1908	1907+CC1056
+CC1056	1906	1815						
CLOUF8	57	10	808	809	810	817	910	909+CC57
+CC57	908	816						
CLOUF8	1057	1010	1808	1809	1810	1817	1910	1909+CC1057
+CC1057	1908	1816						
CLOUF8	58	10	810	811	812	818	912	911+CC58
+CC58	910	817						
CLOUF8	1058	1010	1810	1811	1812	1818	1912	1911+CC1058
+CC1058	1910	1817						
CLOUF8	61	16	521	538	509	539	520	522+CC61
+CC61	512	540						

CLOUF8	1061	1016	521	1538	1509	1539	1520	1522+CC1061
+CC1061	512	540						
CLOUF8	68	13	607	634	720	753	706	615+CC68
+CC68	606	627						
CLOUF8	1068	1013	1607	1634	1720	1753	1706	1615+CC1068
+CC1068	1606	1627						
CLOUF8	41	12	720	730	729	721	709	728+CC41
+CC41	706	753						
CLOUF8	1041	1012	1720	1730	1729	1721	1709	1728+CC1041
+CC1041	1706	1753						
CLOUF8	70	26	608	635	722	751	711	636+CC70
+CC70	609	628						
CLOUF8	1070	1026	1608	1635	1722	1751	1711	1636+CC1070
+CC1070	1609	1628						
CLOUF8	71	13	609	636	711	752	712	616+CC71
+CC71	610	629						
CLOUF8	1071	1013	1609	1636	1711	1752	1712	1616+CC1071
+CC1071	1610	1629						
CLOUF8	81	6	408	409	410	428	508	530+CC81
+CC81	507	427						
CLOUF8	1081	1006	1408	1409	1410	1428	1508	1530+CC1081
+CC1081	1507	1427						
CLOUF8	82	25	410	428	508	532	509	429+CC82
+CC82	421	411	90.					
CLOUF8	1082	1025	1410	1428	1508	1532	1509	1429+CC1082
+CC1082	1421	1411						
CLOUF8	86	13	507	534	607	627	606	515+CC86
+CC86	506	531						
CLOUF8	1086	1013	1507	1534	1607	1627	1606	1515+CC1086
+CC1086	1506	1531						
CLOUF8	87	26	507	534	607	626	608	535+CC87
+CC87	508	530						
CLOUF8	1087	1026	1507	1534	1607	1626	1608	1535+CC1087
+CC1087	1508	1530						
CLOUF8	88	26	508	535	608	628	609	536+CC88
+CC88	509	532						
CLOUF8	1088	1026	1508	1535	1608	1628	1609	1536+CC1088
+CC1088	1509	1532						
CLOUF8	89	13	509	536	609	629	610	517+CC89
+CC89	510	533						
CLOUF8	1089	1013	1509	1536	1609	1629	1610	1517+CC1089
+CC1089	1510	1533						
CLOUF8	118	14	422	431	410	338	309	331+CC118
+CC118	323	341	90.					
CLOUF8	1118	1014	422	1431	1410	1338	1309	1331+CC1118
+CC1118	1323	1341						
CLOUF8	119	14	521	537	508	428	410	431+CC119
+CC119	422	430	90.					
CLOUF8	1119	1014	521	1537	1508	1428	1410	1431+CC1119
+CC1119	422	430						

CLOUF8	120	14	622	630	608	535	508	537+CC120
+CC120	521	541	90.					
CLOUF8	1120	1014	622	1630	1608	1535	1508	1537+CC1120
+CC1120	521	541						
CLOUF8	131	15	742	741	758	759	754	736+CC131
+CC131	724	760						
CLOUF8	1131	1015	742	1741	1758	1759	1754	1736+CC1131
+CC1131	724	760						
CLOUF8	136	11	745	756	746	747	735	734+CC136
+CC136	723	744						
CLOUF8	1136	1011	745	1756	1746	1747	1735	1734+CC1136
+CC1136	723	744						
CLOUF8	137	11	746	757	749	748	726	727+CC137
+CC137	735	747						
CLOUF8	1137	1011	1746	1757	1749	1748	1726	1727+CC1137
+CC1137	1735	1747						
CLOUF8	169	11	323	331	309	310	311	332+CC169
+CC169	334	327						
CLOUF8	1169	1011	1323	1331	1309	1310	1311	1332+CC1169
+CC1169	1334	1327						
CLOUF8	176	32	309	338	410	411	421	340+CC176
+CC176	311	310						
CLOUF8	1176	1032	1309	1338	1410	1411	1421	1340+CC1176
+CC1176	1311	1310						
CLOUF8	180	12	729	732	722	708	710	731+CC180
+CC180	709	721						
CLOUF8	1180	1012	1729	1732	1722	1708	1710	1731+CC1180
+CC1180	1709	1721						
CLOUF8	181	19	740	739	770	771	755	734+CC181
+CC181	723	763						
CLOUF8	1181	1019	740	1739	1770	1771	1735	1734+CC1181
+CC1181	723	763						
CLOUF8	182	18	622	631	609	632	621	620+CC182
+CC182	612	633	55.					
CLOUF8	1182	1018	622	1631	1609	1632	1621	1620+CC1182
+CC1182	612	633						
CLOUF8	185	19	770	769	711	773	726	727+CC185
+CC185	735	771						
CLOUF8	1185	1019	1770	1769	1711	1773	1726	1727+CC1185
+CC1185	1735	1771						
CLOUF6	17	32	309	316	408	409	410	338
CLOUF6	1017	1032	1309	1316	1408	1409	1410	1338
CLOUF6	18	5	311	321	320	324	421	340
CLOUF6		1018	1005	1311	1321	1320	1324	1421
CLOUF6		34	29	812	818	912	921	920
CLOUF6		1034	1029	1812	1818	1912	1921	1920
CLOUF6		62	16	509	533	510	511	520
CLOUF6		1062	1016	1509	1533	1510	1511	1520
CLOUF6	63	16	521	537	508	532	509	538
CLOUF6		69	27	720	730	729	637	607
								634

CLOOF6	1069	1027	1720	1730	1729	1637	1607	1634
CLOOF6	75	31	110	117	210	211	220	120
CLOOF6	1075	1031	1110	1117	1210	1211	1220	1120
CLOOF6	80	25	406	426	507	427	408	407
CLOOF6	1080	1025	1406	1426	1507	1427	1408	1407
CLOOF6	83	13	421	429	509	533	510	417
CLOOF6	1083	1013	1421	1429	1509	1533	1510	1417
CLOOF6	84	13	406	426	507	531	506	415
CLOOF6	1084	1013	1406	1426	1507	1531	1506	1415
CLOOF6	116	20	220	227	334	328	320	222
CLOOF6	1116	1020	1220	1227	1334	1328	1320	1222
CLOOF6	117	14	422	341	323	330	322	339
CLOOF6	1117	1014	422	1341	1323	1330	322	339
CLOOF6	121	5	421	423	520	511	510	417
CLOOF6	1121	1005	1421	1423	1520	1511	1510	1417
CLOOF6	132	15	758	743	722	755	754	759
CLOOF6	1132	1015	1758	1743	1722	1755	1754	1759
CLOOF6	138	17	746	757	749	766	718	761
CLOOF6	1138	1017	1746	1757	1749	1766	1718	1761
CLOOF6	156	11	334	328	320	335	312	336
CLOOF6	1156	1011	1334	1328	1320	1335	312	1336
CLOOF6	157	11	812	821	820	825	920	822
CLOOF6	1157	1011	1812	1821	1820	1825	1920	1822
CLOOF6	161	30	710	717	810	765	718	764
CLOOF6	1161	1030	1710	1717	1810	1765	1718	1764
CLOOF6	162	3	710	733	712	750	718	764
CLOOF6	1162	1003	1710	1733	1712	1750	1718	1764
CLOOF6	163	30	810	811	812	767	718	765
CLOOF6	1163	1030	1810	1811	1812	1767	1718	1765
CLOOF6	167	11	322	330	323	327	334	333
CLOOF6	1167	1011	322	1330	1323	1327	1334	1333
CLOOF6	166	20	220	226	323	327	334	227
CLOOF6	1166	1020	1220	1226	1323	1327	1334	1227
CLOOF6	168	11	322	333	334	336	312	337
CLOOF6	1168	1011	322	1333	1334	1336	312	337
CLOOF6	170	11	334	332	311	321	320	328
CLOOF6	1170	1011	1334	1332	1311	1321	1320	1328
CLOOF6	183	18	609	629	610	611	621	632
CLOOF6	1183	1018	1609	1629	1610	1611	1621	1632
CLOOF6	184	18	622	630	608	628	609	631+CF184
+CF184	90.0							
CLOOF6	1184	1018	622	1630	1608	1628	1609	1631
+CF1184	90.							+CF1184
CLOOF6	186	19	711	752	712	725	726	773
CLOOF6	1186	1019	1711	1752	1712	1725	1726	1773
CLOOF6	188	9	607	626	608	638	729	637
CLOOF6	1188	1009	1607	1626	1608	1638	1729	1637
CLOOF6	189	27	729	732	722	635	608	638
CLOOF6	1189	1027	1729	1732	1722	1635	1608	1638
CLOOF3	106	1	121	122	100	1	1	1+CL106

+CL106								+CO106
+CU106	1,390	0,0	1,390	0,0	1,390	0,0		
CLOUF3	1106	1	121	1122	1100	1	1	1+CL 106
+CL 106								+CU 106
+CU 106	1,390	0,0	1,390	0,0	1,390	0,0		
CLOUF3	207	1	223	224	200	6	6	6+CL207
+CL207								+CO207
+CU207	1,880	0,0	1,880	0,0	1,880	0,0		
CLOUF3	1207	1	223	1224	1200	6	6	6+CL 207
+CL 207								+CU 207
+CO 207	1,880	0,0	1,880	0,0	1,880	0,0		
CLOUF3	305	1	325	326	300	10	10	10+CL305
+CL305								+CU305
+CU305	1,7	1,7		1,7				
CLOUF3	1305	1	325	1326	1300	10	10	10+CL 305
+CL 305								+CU 305
+CO 305	1,7	1,7		1,7				
CLOUF3	504	1	529	528	500	18	18	18+CL504
+CL504								+CO504
+CO504	1,268	1,268		1,268				
CLOUF3	1504	1	529	1528	1500	18	18	18+CL 504
+CL 504								+CU 504
+CU 504	1,268	1,268		1,268				
CLOUF3	604	1	624	625	600	22	22	22+CL604
+CL604								+CU604
+CU604	1,55	1,55		1,55				
CLOUF3	1604	1	624	1625	1600	22	22	22+CL 604
+CL 604								+CO 604
+CU 604	1,55	1,55		1,55				
CLOUF3	703	1	737	738	700	26	26	26+CL703
+CL703								+CO703
+CU703	2,38	2,38		2,38				
CLOUF3	1703	1	737	1738	1700	26	26	26+CL 703
+CL 703								+CO 703
+CO 703	2,38	2,38		2,38				
CLOUF3	806	1	823	824	800	30	30	30+CL806
+CL806								+CO806
+CU806	1,95	1,95		1,95				
CLOUF3	1806	1	823	1824	1800	30	30	30+CL 806
+CL 806								+CO 806
+CU 806	1,95	1,95		1,95				
CLOUF3	907	1	922	923	900	34	34	34+CL907
+CL907								+CO907
+CU907	2,2	2,2		2,2				
CLOUF3	1907	1	922	1923	1900	34	34	34+CL 907
+CL 907								+CO 907
+CU 907	2,2	2,2		2,2				
CLOUF3	100	1	100	101	102	1	1	1+CL100
+CL100								+CO100
+CO100	-1,350	0,0	-1,350	0,0	-1,350	0,		

CLOOF3	1100	1	1100	1101	1102	1	1	1+CL 100
+CL 100								+CO 100
+CO 100	-1.350	0.0	-1.350	0.0	-1.350	0.		
CLOOF3	101	1	102	103	104	2	2	2+CL101
+CL101								+CO101
+CO101	-1.350	0.0	-1.350	0.0	-1.350	0.		
CLOOF3	1101	1	1102	1103	1104	2	2	2+CL 101
+CL 101								+CO 101
+CO 101	-1.350	0.0	-1.350	0.0	-1.350	0.		
CLOOF3	102	1	104	105	106	3	3	3+CL102
+CL102								+CO102
+CO102	-1.330	0.0	-1.330	0.0	-1.330	0.		
CLOOF3	1102	1	1104	1105	1106	3	3	3+CL 102
+CL 102								+CO 102
+CO 102	-1.330	0.0	-1.330	0.0	-1.330	0.		
CLOOF3	103	1	106	107	108	4	4	4+CL103
+CL103								+CO103
+CO103	-1.500	0.0	-1.500	0.0	-1.500	0.		
CLOOF3	1103	1	1106	1107	1108	4	4	4+CL 103
+CL 103								+CO 103
+CO 103	-1.500	0.0	-1.500	0.0	-1.500	0.		
CLOOF3	104	1	108	109	110	5	5	5+CL104
+CL104								+CO104
+CO104	-2.08 0.0		-2.08 0.0		-2.08 0.0			
CLOOF3	1104	1	1108	1109	1110	5	5	5+CL 104
+CL 104								+CO 104
+CO 104	-2.08 0.0		-2.08 0.0		-2.08 0.0			
CLOOF3	105	1	110	111	112	5	5	5+CL105
+CL105								+CO105
+CO105	-2.08 0.0		-2.08 0.0		-2.08 0.0			
CLOOF3	1105	1	1110	1111	112	5	5	5+CL 105
+CL 105								+CO 105
+CO 105	-2.08 0.0		-2.08 0.0		-2.08 0.0			
CLOOF3	107	1	100	119	200	61	61	61+CL107
+CL107	1. 0.		0. 1					+CO107
+CO107	0.78 0.0		0.78 0.0		0.78 0.0			
CLOOF3	2107	1	1100	1119	1200	61	61	61+CL1107
+CL1107	1. 0.0000.		1					+CO1107
+CO1107	0.78 0.0		0.78 0.0		0.78 0.0			
CLOOF3	200	1	200	201	202	6	6	6+CL200
+CL200								+CO200
+CO200	-1.880	0.0	-1.880	0.0	-1.880	0.		
CLOOF3	1200	1	1200	1201	1202	6	6	6+CL 200
+CL 200								+CO 200
+CO 200	-1.880	0.0	-1.880	0.0	-1.880	0.		
CLOOF3	201	1	202	203	204	7	7	7+CL201
+CL201								+CO201
+CO201	-2.130	0.0	-2.130	0.0	-2.130	0.		
CLOOF3	1201	1	1202	1203	1204	7	7	7+CL 201
+CL 201								+CO 201

+CU 201	-2.130	0.0	-2.130	0.0	-2.130	0.			
CUUF3	202	1	204	205	206	7	7		8+CL202
+CL202									+CU202
+CU202	-2.130	0.0	-2.130	0.0	-1.940	0.			
CUUF3	1202	1	1204	1205	1206	7	7		8+CL 202
+CL 202									+CU 202
+CU 202	-2.130	0.0	-2.130	0.0	-1.940	0.			
CUUF3	203	1	206	207	208	8	8		8+CL203
+CL203									+CU203
+CU203	-1.940	0.0	-1.940	0.0	-1.940	0.			
CUUF3	1203	1	1206	1207	1208	8	8		8+CL 203
+CL 203									+CU 203
+CU 203	-1.940	0.0	-1.940	0.0	-1.940	0.			
CUUF3	204	1	208	209	210	8	8		8+CL204
+CL204									+CU204
+CU204	-1.940	0.0	-1.940	0.0	-1.940	0.			
CUUF3	1204	1	1208	1209	1210	8	8		8+CL 204
+CL 204									+CU 204
+CU 204	-1.940	0.0	-1.940	0.0	-1.940	0.			
CUUF3	205	1	210	211	220	9	9		9+CL205
+CL205									+CU205
+CU205	-1.215	0.0	-1.215	0.0	-1.215	0.			
CUUF3	1205	1	1210	1211	1220	9	9		9+CL 205
+CL 205									+CU 205
+CU 205	-1.215	0.0	-1.215	0.0	-1.215	0.			
CUUF3	206	1	220	221	212	9	9		9+CL206
+CL206									+CU206
+CU206	-1.215	0.0	-1.215	0.0	-1.215	0.			
CUUF3	1206	1	1220	1221	212	9	9		9+CL 206
+CL 206									+CU 206
+CU 206	-1.215	0.0	-1.215	0.0	-1.215	0.			
CUUF3	208	1	200	219	300	61	61		61+CL208
+CL208	1.	0.	0.	1					+CU208
+CU208	0.78	0.0	0.78	0.0	0.78	0.0			
CUUF3	2208	1	1200	1219	1300	61	61		61+CL1208
+CL1208	1.	0.0000.	1						+CU1208
+CU1208	0.78	0.0	0.78	0.0	0.78	0.0			
CUUF3	209	1	311	217	210	47	48		48+CL209
+CL209	0.	1.	0.	1					+CU209
+CU209	-0.76	0.0	-0.923	0.0	-0.923	0.0			
CUUF3	2209	1	1311	1217	1210	47	48		48+CL1209
+CL1209	0.	-1.0000.	1						+CU1209
+CU1209	-0.76	0.0	-0.923	0.0	-0.923	0.0			
CUUF3	250	1	334	332	311	57	57		57+CL250
+CL250	0.	0.	1.	1					+CU250
+CU250	-1.41	0.0	-1.41	0.0	-1.41	0.0			
CUUF3	2250	1	1334	1332	1311	57	57		57+CL1250
+CL1250	0.	0.0001.	1						+CU1250
+CU1250	-1.41	0.0	-1.41	0.0	-1.41	0.0			
CUUF3	251	1	722	635	608	53	66		67+CL251

+CL251	0.	1.	0.	1														+CO251
+CO251	.50	0.0	.24	0.0		.113	0.0											
CLOOF3		2251		1	1722	1635	1608	53		66								67+CL1251
+CL1251	0.		-1.0000	1														+CO1251
+CO1251	.50	0.0	.24	0.0		.113	0.0											
CLOOF3		252		1	711	771	740	74		74								74+CL252
+CL252																		+CO252
+CO252	-.40	0.0	-.40	0.0		-.40	0.0											
CLOOF3		1252		1	1711	1771	740	74		74								74+CL 252
+CL 252																		+CO 252
+CO 252	-.40	0.0	-.40	0.0		-.40	0.0											
CLOOF3		300		1	300	301	302	10		11								11+CL300
+CL300																		+CO300
+CO300	-1.500	0.0	-1.500	0.0	-1.500	0.												
CLOOF3		1300		1	1300	1301	1302	10		11								11+CL 300
+CL 300																		+CO 300
+CO 300	-1.500	0.0	-1.500	0.0	-1.500	0.												
CLOOF3		301		1	302	303	304	11		11								12+CL301
+CL301																		+CO301
+CO301	-1.500	0.0	-1.500	0.0	-1.270	0.												
CLOOF3		1301		1	1302	1303	1304	11		11								12+CL 301
+CL 301																		+CO 301
+CO 301	-1.500	0.0	-1.500	0.0	-1.270	0.												
CLOOF3		302		1	304	305	306	12		12								12+CL302
+CL302																		+CO302
+CO302	-1.270	0.0	-1.270	0.0	-1.270	0.												
CLOOF3		1302		1	1304	1305	1306	12		12								12+CL 302
+CL 302																		+CO 302
+CO 302	-1.270	0.0	-1.270	0.0	-1.270	0.												
CLOOF3		303		1	306	307	309	13		13								13+CL303
+CL303																		+CO303
+CO303	-1.040	0.0	-1.040	0.0	-1.040	0.												
CLOOF3		1303		1	1306	1307	1309	13		13								13+CL 303
+CL 303																		+CO 303
+CO 303	-1.040	0.0	-1.040	0.0	-1.040	0.												
CLOOF3		304		1	309	310	311	13		13								13+CL304
+CL304																		+CO304
+CO304	-1.04	0.0	-1.04	0.0	-1.04	0.0												
CLOOF3		1304		1	1309	1310	1311	13		13								13+CL 304
+CL 304																		+CO 304
+CO 304	-1.04	0.0	-1.04	0.0	-1.04	0.0												
CLOOF3		306		1	300	319	400	61		61								61+CL306
+CL306	1.	0.	0.	1														+CO306
+CO306	0.78	0.0	0.78	00.0	0.78	0.0												
CLOOF3		2306		1	1300	1319	1400	61		61								61+CL1306
+CL1306	1.		0.0000	1														+CO1306
+CO1306	0.78	0.0	0.78	00.0	0.78	0.0												
CLOOF3		307		1	309	310	311	13		13								13+CL307
+CL307																		+CO307
+CO307	-1.040	0.0	-1.040	0.0	-1.040	0.												

CLOUF3	1307	1	1309	1310	1311	13	13	13+CL 307
+CL 307								+CO 307
+CU 307	-1.040	0.0	-1.040	0.0	-1.040	0.	.	
CLOUF3	308	1	421	340	311	45	45	46+CL308
+CL308	0.	1.	0.	1				+CO308
+CU308	-.80	0.0	-.80	0.0	-.80	0.0		
CLOUF3	2308	1	1421	1340	1311	45	45	46+CL1308
+CL1308	0.	-1.0000.	1					+CO1308
+CU1308	-.80	0.0	-.80	0.0	-.80	0.0		
CLOUF3	309	1	312	318	412	81	80	79+CL309
+CL309	1.	0.	0.	1				+CO309
+CU309	-0.57	0.0	-0.63	0.0	-0.69	0.		
CLOUF3	2309	1	312	318	412	81	80	79+CL1309
+CL1309	1.	0.0000.	1					+CO1309
+CU1309	-0.57	0.0	-0.63	0.0	-0.69	0.		
CLOUF3	400	1	400	401	402	14	14	15+CL400
+CL400								+CO400
+CU400	-1.185	0.0	-1.185	0.0	-1.440	0.		
CLOUF3	1400	1	1400	1401	1402	14	14	15+CL 400
+CL 400								+CO 400
+CU 400	-1.185	0.0	-1.185	0.0	-1.440	0.		
CLOUF3	401	1	402	403	404	15	15	16+CL401
+CL401								+CO401
+CU401	-1.440	0.0	-1.440	0.0	-2.071	0.		
CLOUF3	1401	1	1402	1403	1404	15	15	16+CL 401
+CL 401								+CO 401
+CU 401	-1.440	0.0	-1.440	0.0	-2.071	0.		
CLOUF3	402	1	404	405	406	16	16	16+CL402
+CL402								+CO402
+CU402	-2.071	0.0	-2.071	0.0	-2.071	0.		
CLOUF3	1402	1	1404	1405	1406	16	16	16+CL 402
+CL 402								+CO 402
+CU 402	-2.071	0.0	-2.071	0.0	-2.071	0.		
CLOUF3	403	1	406	407	408	16	17	17+CL403
+CL403								+CO403
+CU403	-2.071	0.0	-2.268	0.0	-2.268	0.		
CLOUF3	1403	1	1406	1407	1408	16	17	17+CL 403
+CL 403								+CO 403
+CU 403	-2.071	0.0	-2.268	0.0	-2.268	0.		
CLOUF3	404	1	408	409	410	17	17	17+CL404
+CL404								+CO404
+CU404	-2.268	0.0	-2.268	0.0	-2.268	0.		
CLOUF3	1404	1	1408	1409	1410	17	17	17+CL 404
+CL 404								+CO 404
+CU 404	-2.268	0.0	-2.268	0.0	-2.268	0.		
CLOUF3	405	1	410	411	421	17	17	17+CL405
+CL405								+CO405
+CU405	-2.268	0.0	-2.268	0.0	-2.268	0.		
CLOUF3	1405	1	1410	1411	1421	17	17	17+CL 405
+CL 405								+CO 405

+CO 405	-2,268	0.0	-2,268	0.0	-2,268	0.			
CLOUF3	406	1	400	419	500	61	61	61+CL406	
+CL406	1.	0.	0.	1				+CO406	
+CO406	0.78	0.0	0.78	0.0	0.78	0.0			
CLOUF3	2406	1	1400	1419	1500	61	61	61+CL1406	
+CL1406	1.	0.0000.	1					+CO1406	
+CO1406	0.78	0.0	0.78	0.0	0.78	0.0			
CLOUF3	407	1	510	417	421	43	43	43+CL407	
+CL407	0.	1.	0.	1				+CO407	
+CO407	0.0	-.75	0.0	-.75	0.0	-.75			
CLOUF3	2407	1	1510	1417	1421	43	43	43+CL1407	
+CL1407	0.	-1.0000.	1					+CO1407	
+CO1407	0.0	-.75	0.0	-.75	0.0	-.75			
CLOUF3	408	1	509	429	421	42	42	42+CL408	
+CL408	0.	1.	0.	1				+CO408	
+CO408	0.0	.75	0.0	.75	0.0	.75			
CLOUF3	2408	1	1509	1429	1421	42	42	42+CL1408	
+CL1408	0.	-1.0000.	1					+CO1408	
+CO1408	0.0	.75	0.0	.75	0.0	.75			
CLOUF3	409	1	410	431	422	73	73	73+CL409	
+CL409	0.	0.	1.	1				+CO409	
+CO409	0.0	.94	0.0	.94	0.0	.94			
CLOUF3	2409	1	1410	1431	422	73	73	73+CL1409	
+CL1409	0.	0.0001.	1					+CO1409	
+CO1409	0.0	.94	0.0	.94	0.0	.94			
CLOUF3	410	1	412	418	512	79	80	81+CL410	
+CL410	1.	0.	0.	1				+CO410	
+CO410	-0.69	0.0	-0.63	0.0	-0.57	0.			
CLOUF3	2410	1	412	418	512	79	80	81+CL1410	
+CL1410	1.	0.0000.	1					+CO1410	
+CO1410	-0.69	0.0	-0.63	0.0	-0.57	0.			
CLOUF3	500	1	500	501	502	19	19	19+CL500	
+CL500								+CO500	
+CO500	-1,880	0.0	-1,880	0.0	-1,880	0.			
CLOUF3	1500	1	1500	1501	1502	19	19	19+CL 500	
+CL 500								+CO 500	
+CO 500	-1,880	0.0	-1,880	0.0	-1,880	0.			
CLOUF3	501	1	502	503	504	19	20	20+CL501	
+CL501								+CO501	
+CO501	-1,880	0.0	-2,500	0.0	-2,500	0.			
CLOUF3	1501	1	1502	1503	1504	19	20	20+CL 501	
+CL 501								+CO 501	
+CO 501	-1,880	0.0	-2,500	0.0	-2,500	0.			
CLOUF3	502	1	504	505	507	20	20	20+CL502	
+CL502								+CO502	
+CO502	-2,500	0.0	-2,500	0.0	- 0.0	0.			
CLOUF3	1502	1	1504	1505	1507	20	20	20+CL 502	
+CL 502								+CO 502	
+CO 502	-2,500	0.0	-2,500	0.0	- 0.0	0.			
CLOUF3	503	1	507	530	508	21	21	21+CL503	

CLOUF3	1603	1	1607	1626	1608	25	25	25+CL 603
+CL 603								+CO 603
+CU 603	-0.941	0.0	-0.941	0.0	-0.941	0.		
CLOUF3	605	1	608	628	609	25	25	25+CL 605
+CL 605								+CO 605
+CU 605	-0.941	0.0	-0.941	0.0	-0.941	0.		
CLOUF3	1605	1	1608	1628	1609	25	25	25+CL 605
+CL 605								+CO 605
+CO 605	-0.941	0.0	-0.941	0.0	-0.941	0.		
CLOUF3	606	1	600	619	700	61	61	61+CL 606
+CL 606	1.	0.	0.	1				+CO 606
+CU 606	0.78	0.0	0.78	0.0	0.78	0.0		
CLOUF3	2606	1	1600	1619	1700	61	61	61+CL 1606
+CL 1606	1.	0.0000.	1					+CO 1606
+CU 1606	0.78	0.0	0.78	0.0	0.78	0.0		
CLOUF3	607	1	712	616	610	43	43	43+CL 607
+CL 607	0.	1.	0.	1				+CO 607
+CU 607	0.0	-.75	0.0	-.75	0.0	-.75		
CLOUF3	2607	1	1712	1616	1610	43	43	43+CL 1607
+CL 1607	0.	-1.0000.	1					+CO 1607
+CU 1607	0.0	-.75	0.0	-.75	0.0	-.75		
CLOUF3	608	1	711	636	609	42	42	42+CL 608
+CL 608	0.	1.	0.	1				+CO 608
+CU 608	0.0	.75	0.0	.75	0.0	.75		
CLOUF3	2608	1	1711	1636	1609	42	42	42+CL 1608
+CL 1608	0.	-1.0000.	1					+CO 1608
+CU 1608	0.0	.75	0.0	.75	0.0	.75		
CLOUF3	700	1	700	701	702	26	27	27+CL 700
+CL 700								+CO 700
+CU 700	-2.380	0.0	-2.490	0.0	-2.490	0.		
CLOUF3	1700	1	1700	1701	1702	26	27	27+CL 700
+CL 700								+CO 700
+CO 700	-2.380	0.0	-2.490	0.0	-2.490	0.		
CLOUF3	701	1	702	703	704	27	27	28+CL 701
+CL 701								+CO 701
+CO 701	-2.490	0.0	-2.490	0.0	-2.775	0.		
CLOUF3	1701	1	1702	1703	1704	27	27	28+CL 701
+CL 701								+CO 701
+CO 701	-2.490	0.0	-2.490	0.0	-2.775	0.		
CLOUF3	702	1	704	705	720	28	28	28+CL 702
+CL 702								+CO 702
+CO 702	-2.775	0.0	-2.775	0.0	0.0	0.		
CLOUF3	1702	1	1704	1705	1720	28	28	28+CL 702
+CL 702								+CO 702
+CO 702	-2.775	0.0	-2.775	0.0	0.0	0.		
CLOUF3	704	1	700	719	800	61	61	61+CL 704
+CL 704	1.	0.	0.	1				+CO 704
+CO 704	0.78	0.0	0.78	0.0	0.78	0.0		
CLOUF3	2704	1	1700	1719	1800	61	61	61+CL 1704
+CL 1704	1.	0.0000.	1					+CO 1704

+CU1704	0.78	0.0	0.78	0.0	0.78	0.0			
CUUF3	705		1	812	767	718	40	40	40+CL705
+CL705	1.	0.	0.	1					+CO705
+CU705	0.0	2.0	0.0	2.0	0.0	2.0			
CUUF3	2705		1	1812	1767	1718	40	40	40+CL1705
+CL1705	1.	0.0000.		1					+CO1705
+CU1705	0.0	2.0	0.0	2.0	0.0	2.0			
CUUF3	706		1	718	750	712	40	40	40+CL706
+CL706	1.	0.	0.	1					+CO706
+CU706	0.0	2.0	0.0	2.0	0.0	2.0			
CUUF3	2706		1	1718	1750	1712	40	40	40+CL1706
+CL1706	1.	0.0000.		1					+CO1706
+CU1706	0.0	2.0	0.0	2.0	0.0	2.0			
CUUF3	707		1	754	772	770	50	51	52+CL707
+CL707	0.	0.	1.	1					+CO707
+CU707	-0.73	0.0	-1.07	0.0	-1.38	0.0			
CUUF3	2707		1	1754	1772	1770	50	51	52+CL1707
+CL1707	0.	0.0001.		1					+CO1707
+CU1707	-0.73	0.0	-1.07	0.0	-1.38	0.0			
CUUF3	708	1	735	623	612	49	49	49	+CL708
+CL708	1.	0.	0.	1					+CO708
+CU708	1.28	0.0	1.28	0.0	1.28	-3.0			
CUUF3	2708		1	1735	1623	612	49	49	49+CL1708
+CL1708	1.	0.0000.		1					+CO1708
+CU1708	1.28	0.0	1.28	0.0	1.28	-3.0			
CUUF3	709		1	746	756	745	54	54	54+CL709
+CL709	1.	0.	0.	1					
CUUF3	2709		1	1746	1756	745	54	54	54+CL1709
+CL1709	1.	0.0000.		1					+CO1709
+CU1709	1.28	0.0	1.28	0.0	1.28	-3.0			
CUUF3	710		1	820	768	749	59	59	59+CL710
+CL710	0.	1.	0.	1					
CUUF3	2710		1	1820	1768	1749	59	59	59+CL1710
+CL1710	0.	-1.0000.		1					+CO1710
+CU1710	1.28	0.0	1.28	0.0	1.28	-3.0			
CUUF3	711		1	722	743	758	53	53	75+CL711
+CL711									+CO711
+CU711	-0.50	.20	-0.50	.20	-0.50	.20			
CUUF3	1711		1	1722	1743	1758	53	53	75+CL 711
+CL 711									+CO 711
+CU 711	-0.50	.20	-0.50	.20	-0.50	.20			
CUUF3	712		1	758	741	742	75	75	75 +CL712
+CL712									+CO712
+CU712	-0.50	.20	-0.50	.20	-0.50	.20			
CUUF3	1712		1	1758	1741	742	75	75	75+CL 712
+CL 712									+CO 712
+CU 712	-0.50	.20	-0.50	.20	-0.50	.20			
CUUF3	713		1	720	730	729	85	85	85+CL713
+CL713	0.	0.	1.	1					+CO713
+CU713	0.0	-1.28	0.0	-1.28	0.	-1.28			

CLUOF3	2713	1	1720	1730	1729	85	85	85+CL1713
+CL1713	0.	0.0001.	1					+CO1713
+CU1713	0.0	-1.28	0.0	-1.28	0.	-1.28		
CLUOF3	714	1	729	732	722	85	86	87+CL714
+CL714	0.	0.	1.	1				+CU714
+CU714	0.0	-1.28	0.0	-1.97	0.0	-2.45		
CLUOF3	2714	1	1729	1732	1722	85	86	87+CL1714
+CL1714	0.	0.0001.	1					+CO1714
+CU1714	0.0	-1.28	0.0	-1.97	0.0	-2.45		
CLUOF3	717	1	770	771	735	52	78	78+CL717
+CL717	0.	0.	1.	1				+CU717
+CU717	-1.38	0.0	-1.50	0.0	-1.50	0.0		
CLUOF3	2717	1	1770	1771	1735	52	78	78+CL1717
+CL1717	0.	0.0001.	1					+CO1717
+CU1717	-1.38	0.0	-1.50	0.0	-1.50	0.0		
CLUOF3	806	1	823	824	800	30	30	+CL806
+CL806								+CO806
+CU806	-1.95	0.	-1.95	0.	-1.95	0.		
CLUOF3	1806	1	823	1824	1800	30	30	30+CL 806
+CL 806								+CO 806
+CO 806	-1.95	0.	-1.95	0.	-1.95	0.		
CLUOF3	800	1	800	801	802	30	31	31+CL800
+CL800								+CO800
+CO800	-1.950	0.0	-1.535	0.0	-1.535	0.		
CLUOF3	1800	1	1800	1801	1802	30	31	31+CL 800
+CL 800								+CO 800
+CO 800	-1.950	0.0	-1.535	0.0	-1.535	0.		
CLUOF3	801	1	802	803	804	31	31	32+CL801
+CL801								+CO801
+CO801	-1.535	0.0	-1.535	0.0	-1.280	0.		
CLUOF3	1801	1	1802	1803	1804	31	31	32+CL 801
+CL 801								+CO 801
+CO 801	-1.535	0.0	-1.535	0.0	-1.280	0.		
CLUOF3	802	1	804	805	806	32	32	32+CL802
+CL802								+CO802
+CO802	-1.280	0.0	-1.280	0.0	-1.28	0.		
CLUOF3	1802	1	1804	1805	1806	32	32	32+CL 802
+CL 802								+CO 802
+CO 802	-1.280	0.0	-1.280	0.0	-1.28	0.		
CLUOF3	803	1	806	807	808	33	33	33+CL803
+CL803								+CO803
+CO803	-1.410	0.0	-1.410	0.0	-1.410	0.		
CLUOF3	1803	1	1806	1807	1808	33	33	33+CL 803
+CL 803								+CO 803
+CO 803	-1.410	0.0	-1.410	0.0	-1.410	0.		
CLUOF3	804	1	808	809	810	33	33	33+CL804
+CL804								+CO804
+CO804	-1.410	0.0	-1.410	0.0	-1.410	0.		
CLUOF3	1804	1	1808	1809	1810	33	33	33+CL 804
+CL 804								+CO 804

+CU 804	-1.410	0.0	-1.410	0.0	-1.410	0.			
CUUF3	805	1	810	811	812	33	33		33+CL805
+CL805									+CO805
+CU805	-1.410	0.0	-1.410	0.0	-1.410	0.			
CUUF3	1805	1	1810	1811	1812	33	33		33+CL 805
+CL 805									+CO 805
+CU 805	-1.410	0.0	-1.410	0.0	-1.410	0.			
CUUF3	807	1	812	821	820	58	58		58+CL807
+CL807									+CO807
+CU807	-2.31	0.0	-2.31	0.0	-2.31	0.			
CUUF3	1807	1	1812	1821	1820	58	58		58+CL 807
+CL 807									+CO 807
+CU 807	-2.31	0.0	-2.31	0.0	-2.31	0.			
CUUF3	808	1	920	822	812	39	39		39+CL808
+CL808									
CUUF3	1808	1	1920	1822	1812	39	39		39+CL 808
+CL 808									+CO 808
+CU 808	-2.31	0.0	-2.31	0.0	-2.31	0.			
CUUF3	809	1	800	819	900	61	61		61+CL809
+CL809	1.	0.	0.	1					+CO809
+CU809	0.78	0.0	0.78	0.0	0.78	0.0			
CUUF3	2809	1	1800	1819	1900	61	61		61+CL1809
+CL1809	1.	0.0000.	1						+CO1809
+CU1809	0.78	0.0	0.78	0.0	0.78	0.0			
CUUF3	810	1	920	826	746	41	41		+CL810
+CL810	1.	0.	0.	1					+CO810
+CU810	0.0	1.15	0.0	0.0	1.04	0.0			
CUUF3	2810	1	1920	1826	1746	41	41		41+CL1810
+CL1810	1.	0.0000.	1						+CO1810
+CU1810	0.0	1.15	0.0	0.0	1.04	0.0			
CUUF3	811	1	746	747	723	41	41		41+CL811
+CL811	1.	0.	0.	1					+CO811
+CU811	1.04	0.0	1.04	0.0	2.0	0.0			
CUUF3	2811	1	1746	1747	723	41	41		41+CL1811
+CL1811	1.	0.0000.	1						+CO1811
+CU1811	1.04	0.0	1.04	0.0	2.0	0.0			
CUUF3	900	1	900	901	902	35	35		35+CL900
+CL900									+CO900
+CU900	-2.61	0.0	-2.610	0.0	-2.610	0.			
CUUF3	1900	1	1900	1901	1902	35	35		35+CL 900
+CL 900									+CO 900
+CU 900	-2.61	0.0	-2.610	0.0	-2.610	0.			
CUUF3	901	1	902	903	904	36	36		36+CL901
+CL901									+CO901
+CU901	-3.120	0.0	-3.120	0.0	-3.120	0.			
CUUF3	1901	1	1902	1903	1904	36	36		36+CL 901
+CL 901									+CO 901
+CU 901	-3.120	0.0	-3.120	0.0	-3.120	0.			
CUUF3	902	1	904	905	906	36	36		36+CL902
+CL902									+CO902

+CU902	-3.120	0.0	-3.120	0.0	-3.120	0.			
CLUUF3	1902	1	1904	1905	1906	36	36	36+CL 902	
+CL 902								+CU 902	
+CU 902	-3.120	0.0	-3.120	0.0	-3.120	0.			
CLUUF3	903	1	906	907	908	37	37	37+CL903	
+CL903								+CU903	
+CU903	-2.750	0.0	-2.750	0.0	-2.750	0.			
CLUUF3	1903	1	1906	1907	1908	37	37	37+CL 903	
+CL 903								+CU 903	
+CU 903	-2.750	0.0	-2.750	0.0	-2.750	0.			
CLUUF3	904	1	908	909	910	37	37	37+CL904	
+CL904								+CU904	
+CU904	-2.750	0.0	-2.750	0.0	-2.750	0.			
CLUUF3	1904	1	1908	1909	1910	37	37	37+CL 904	
+CL 904								+CU 904	
+CU 904	-2.750	0.0	-2.750	0.0	-2.750	0.			
CLUUF3	905	1	910	911	912	37	37	38+CL905	
+CL905								+CU905	
+CU905	-2.750	0.0	-2.750	0.0	-1.500	0.			
CLUUF3	1905	1	1910	1911	1912	37	37	38+CL 905	
+CL 905								+CU 905	
+CU 905	-2.750	0.0	-2.750	0.0	-1.500	0.			
CLUUF3	906	1	912	921	920	38	38	38+CL906	
+CL906								+CU906	
+CU906	-1.500	0.0	-1.500	0.0	-1.500	0.			
CLUUF3	1906	1	1912	1921	1920	38	38	38+CL 906	
+CL 906								+CU 906	
+CU 906	-1.500	0.0	-1.500	0.0	-1.500	0.			
PLOUF	11	1	0.04						
PLOUF	12	1	.11						
PLOUF	13	1	.06						
PLOUF	15	1	.14						
PLOUF	17	1	.075						
PLOUF	19	1	.12						
PLOUF	20	1	.07						
PLOUF	1011	1	0.04						
PLOUF	1012	1	.11						
PLOUF	1013	1	.06						
PLOUF	1015	1	.14						
PLOUF	1017	1	.075						
PLOUF	1019	1	.12						
PLOUF	1020	1	.07						
PLOUFx	1	4.637+5	1.549+5	0.0	6.68+5	0.0	1.544+5	+PL1	
+PL1	5.86E-3	0.	0.	0.	-7.56+4	0.	0.	+PLL1	
+PLL1		61.83	20.65	0.0	4.75+4	0.0	72.48		
PLOUFx	2	4.637+5	1.549+5	0.0	1.031+6	0.0	1.544+5	+PL2	
+PL2	9.32E-3	0.	0.	0.	-2.10+5	0.	0.	+PLL2	
+PLL2		61.83	20.65	0.0	1.318+5	0.0	1.295+2		
PLOUFx	3	4.637+5	1.549+5	0.0	7.396+5	0.0	1.544+5	+PL3	
+PL3	6.54E-3	0.	0.	0.	-1.021+5			+PLL3	

+PLL3		61.83	20.65	0.0	64.08	0.0	84.09	
PLUOFX	4	9.411+5	1.549+5	0.0	4.637+5	0.0	1.544+5	+PL4
+PL4	8.46E-3	-2.11+5						+PLL4
+PLL4		1.607+5	20.65	0.0	61.83	0.0	3.364+2	
PLUOFX	5	4.637+5	1.549+5	0.0	7.642+5	0.0	1.544+5	+PL5
+PL5	6.78E-3	0.	0.	0.	-1.112+50.	0.	0.	+PLL5
+PLL5		61.83	20.65	0.0	6.98+4	0.0	87.91	
PLUOFX	6	1.266+6	2.323+5	0.0	6.956+5	0.0	2.316+5	+PL6
+PL6	9.34E-3	-2.576+50.	0.	0.	0.	0.	0.	+PLL6
+PLL6		1.971+5	69.69	0.0	2.087+2	0.0	4.914+2	
PLUOFX	7	4.637+5	1.549+5	0.0	7.191+5	0.0	1.544+5	+PL7
+PL7	6.34-3	0.	0.	0.	-9.451+40.	0.	0.	+PLL7
+PLL7		61.83	20.65	0.0	5.934+4	0.0	80.92	
PLUOFX	8	1.005+6	2.323+5	0.0	6.956+5	0.0	2.316+5	+PL8
+PL8	8.82-3	-6.798+40.	0.	0.	0.	0.	0.	+PLL8
+PLL8		2.976+4	69.70	0.0	2.087+2	0.0	2.355+2	
PLUOFX	9	1.015+6	2.323+5	0.0	6.956+5	0.0	2.316+5	+PL9
+PL9	8.82-3	-1.213+5	0.	0.	0.	0.	0.	+PLL9
+PLL9		7.67+4	69.70	0.0	2.087+2	0.0	1.887+2	
PLUOFX	10	4.637+5	1.549+5	0.0	6.767+5	0.0	1.544+5	+PL10
+PL10	5.04-3	0.	0.	0.	-7.876+4	0.	0.	+PLL10
+PLL10		61.83	20.64	0.0	4.946+4	0.0	74.31	
PLUOFX	14	7.547+5	1.549+5	0.0	4.637+5	0.0	1.544+5	+PL14
+PL14	6.69-3	-1.129+5	0.	0.	0.	0.	0.	+PLL14
+PLL14		7.168+4	20.65	0.0	61.83	0.0	96.17	
PLUOFX	16	1.307+6	1.549+5	0.0	4.637+5	0.0	1.544+5	+PL16
+PL16	1.19-2	-2.278+50.	0.	0.	0.	0.	0.	+PLL16
+PLL16		7.922+4	20.65	0.0	61.83	0.0	8.386+3	
PLUOFX	18	7.384+5	1.549+5	0.0	4.637+5	0.0	1.544+5	+PL18
+PL18	6.53-3	-5.768+4						+PLL18
+PLL18		2.515+4	20.65	0.0	61.83	0.0	126.97	
PLUOFX	1001	4.637+5	1.549+5	0.0	6.68+5	0.0	1.544+5	+PL1
+PL1	5.86E-3	0.	0.	0.	7.56+4	0.	0.	+PLL1
+PLL1		61.83	20.65	0.0	4.75+4	0.0	72.48	
PLUOFX	1002	4.637+5	1.549+5	0.0	1.031+6	0.0	1.544+5	+PL2
+PL2	9.32E-3	0.	0.	0.	2.10+5	0.	0.	+PLL2
+PLL2		61.83	20.65	0.0	1.318+5	0.0	1.295+2	
PLUOFX	1003	4.637+5	1.549+5	0.0	7.396+5	0.0	1.544+5	+PL3
+PL3	6.54E-3	0.	0.	0.	1.021+5			+PLL3
+PLL3		61.83	20.65	0.0	64.08	0.0	84.09	
PLUOFX	1004	9.411+5	1.549+5	0.0	4.637+5	0.0	1.544+5	+PL4
+PL4	8.46E-3	2.11+5						+PLL4
+PLL4		1.607+5	20.65	0.0	61.83	0.0	3.364+2	
PLUOFX	1005	4.637+5	1.549+5	0.0	7.642+5	0.0	1.544+5	+PL5
+PL5	6.78E-3	0.	0.	0.	1.112+50.	0.	0.	+PLL5
+PLL5		61.83	20.65	0.0	6.98+4	0.0	87.91	
PLUOFX	1006	1.266+6	2.323+5	0.0	6.956+5	0.0	2.316+5	+PL6
+PL6	9.34E-3	2.576+5						+PLL6
+PLL6		1.971+5	69.69	0.0	2.087+2	0.0	4.914+2	
PLUOFX	1007	4.637+5	1.549+5	0.0	7.191+5	0.0	1.544+5	+PL7

+PL7	6.34-3	0.	0.	0.	9.451+4			
+PLL7		61.83	20.65	0.0	5.934+4 0.0	80.92	+PLL7	
PLUOFX	1008	1.005+6	2.323+5	0.0	6.956+5 0.0	2.316+5		
+PL8	8.82-3	6.798+4					+PL8	
+PLL8		2.976+4	69.70	0.0	2.087+2 0.0	2.355+2	+PLL8	
PLUOFX	1009	1.015+6	2.323+5	0.0	6.956+5 0.0	2.316+5		
+PL9	8.82-3	1.213+5					+PL9	
+PLL9		7.67+4	69.70	0.0	2.087+2 0.0	1.887+2	+PLL9	
PLUOFX	1010	4.637+5	1.549+5	0.0	6.767+5 0.0	1.544+5		
+PL10	5.04-3				7.876+4		+PL10	
+PLL10		61.83	20.64	0.0	4.946+4 0.0	74.31	+PLL10	
PLUOFX	1014	7.547+5	1.549+5	0.0	4.637+5 0.0	1.544+5		
+PL14	6.69-3	1.129+5					+PL14	
+PLL14		7.168+4	20.65	0.0	61.83 0.0	96.17	+PLL14	
PLUOFX	1016	1.307+6	1.549+5	0.0	4.637+5 0.0	1.544+5		
+PL16	1.19-2	2.278+5					+PL16	
+PLL16		7.922+4	20.65	0.0	61.83 0.0	8.386+3	+PLL16	
PLUOFX	1018	7.384+5	1.549+5	0.0	4.637+5 0.0	1.544+5		
+PL18	6.53-3	5.768+4					+PL18	
+PLL18		2.515+4	20.65		61.83	126.97	+PLL18	
PLUOFX	25	1.266+6	2.323+5	0.0	6.956+5 0.0	2.316+5		
+PL25	9.34E-3	2.576+50.	0.		0. 0.	0.	+PL25	
+PLL25		1.971+5	69.69	0.0	2.087+2 0.0	4.914+2	+PLL25	
PLUOFX	26	1.005+6	2.323+5	0.0	6.956+5 0.0	2.316+5		
+PL26	8.82-3	6.798+40.	0.		0. 0.	0.	+PL26	
+PLL26		2.976+4	69.70	0.0	2.087+2 0.0	2.355+2	+PLL26	
PLUOFX	27	1.015+6	2.323+5	0.0	6.956+5 0.0	2.316+5		
+PL27	8.82-3	1.213+5	0.	0.	0. 0.	0.	+PL27	
+PLL27		7.67+4	69.70	0.0	2.087+2 0.0	1.887+2	+PLL27	
PLUOFX	28	1.307+6	1.549+5	0.0	4.637+5 0.0	1.544+5		
+PL28	1.19-2	2.278+50.	0.		0. 0.	0.	+PL28	
+PLL28		7.922+4	20.65	0.0	61.83 0.0	8.386+3	+PLL28	
PLUOFX	29	4.637+5	1.549+5	0.0	6.767+5 0.0	1.544+5		
+PL29	5.04-3	0.	0.	0.	7.876+4	0.	+PL29	
+PLL29		61.83	20.64	0.0	4.946+4 0.0	74.31	+PLL29	
PLUOFX	30	4.637+5	1.549+5	0.0	7.396+5 0.0	1.544+5		
+PL30	6.54E-3	0.	0.	0.	1.021+5		+PL30	
+PLL30		61.83	20.65	0.0	64.08 0.0	84.09	+PLL30	
PLUOFX	31	4.637+5	1.549+5	0.0	6.68+5 0.0	1.544+5		
+PL31	5.86E-3	0.	0.	0.	7.56+4 0.	0.	+PL31	
+PLL31		61.83	20.65	0.0	4.75+4 0.0	72.48	+PLL31	
PLUOFX	32	4.637+5	1.549+5	0.0	7.642+5 0.0	1.544+5		
+PL32	6.78E-3	0.	0.	0.	1.112+50.	0.	+PL32	
+PLL32		61.83	20.65	0.0	6.98+4 0.0	87.91	+PLL32	
PLUOFX	1025	1.266+6	2.323+5	0.0	6.956+5 0.0	2.316+5		
+PL1025	9.34E-3	-2.576+50.	0.		0. 0.	0.	+PL1025	
+PLL1025		1.971+5	69.69	0.0	2.087+2 0.0	4.914+2	+PLL102	
PLUOFX	1026	1.005+6	2.323+5	0.0	6.956+5 0.0	2.316+5		
+PL1026	8.82-3	-6.798+40.	0.		0. 0.	0.	+PL1026	
+PLL1026		2.976+4	69.70	0.0	2.087+2 0.0	2.355+2	+PLL102	

PLUOFX 1027	1.015+6	2.323+5	0.0	6.956+5	0.0	2.316+5	+PL1027
+PL1027 8.82-3	-1.213+5	0.	0.	0.	0.	0.	+PLL1027
+PLL1027	7.67+4	69.70	0.0	2.087+2	0.0	1.887+2	
PLUOFX 1028	1.307+6	1.549+5	0.0	4.637+5	0.0	1.544+5	+PL1028
+PL1028 1.19-2	-2.278+50.	0.	0.	0.	0.	0.	+PLL1028
+PLL1028	7.922+4	20.65	0.0	61.83	0.0	8.386+3	
PLUOFX 1029	4.637+5	1.549+5	0.0	6.767+5	0.0	1.544+5	+PL1029
+PL1029 5.04-3	0.	0.	0.	-7.876+4	0.	0.	+PLL1029
+PLL1029	61.83	20.64	0.0	4.946+4	0.0	74.31	
PLUOFX 1030	4.637+5	1.549+5	0.0	6.68+5	0.0	1.544+5	+PL1030
+PL1030 6.54E-3	0.	0.	0.	-1.021+5	0.	0.	+PLL1030
+PLL1030	61.83	20.65	0.0	64.08	0.0	84.09	
PLUOFX 1031	4.637+5	1.549+5	0.0	6.68+5	0.0	1.544+5	+PL1031
+PL1031 5.86E-3	0.	0.	0.	-7.56+4	0.	0.	+PLL1031
+PLL1031	61.83	20.65	0.0	4.75+4	0.0	72.48	
PLUOFX 1032	4.637+5	1.549+5	0.0	7.642+5	0.0	1.544+5	+PL1032
+PL1032 6.78E-3	0.	0.	0.	-1.112+50.	0.	0.	+PLL1032
+PLL1032	61.83	20.65	0.0	6.98+4	0.0	87.91	
PLUOF3	1	0.250	0.261	0.92-2	0.300-3		
PLUOF3	2	0.252	0.255	0.0110	0.302-3		
PLUOF3	3	0.285	0.276	0.0122	0.465-3		
PLUOF3	4	0.266	0.319	0.0112	0.319-3		
PLUOF3 5	.33	.710	.418-2	.400-3			
PLUOF3	6	5.360	9.630	0.29+1	0.898-0		
PLUOF3	7	3.890	6.960	1.1700	0.435-0		
PLUOF3	8	2.420	5.570	0.8830	0.646-1		
PLUOF3	9	1.830	2.630	0.8120	0.416-1		
PLUOF3	10	0.349	0.506	0.0116	0.654-3		
PLUOF3	11	0.319	0.370	0.0113	0.598-3		
PLUOF3	12	0.284	0.243	0.0108	0.532-3		
PLUOF3	13	0.251	0.153	0.0103	0.470-3		
PLUOF3	14	0.486	0.561	0.0845	0.104-2		
PLUOF3	15	0.330	0.353	0.0118	0.703-3		
PLUOF3	16	1.080	1.060	0.1430	0.218-1		
PLUOF3	17	1.130	1.200	0.1540	0.248-1		
PLUOF3	18	0.293	0.206	0.0165	0.550-3		
PLUOF3	19	0.337	0.559	5.35-3	0.550-3		
PLUOF3	20	0.424	1.180	7.39-3	0.693-3		
PLUOF3	21	0.284	0.310	0.76-2	0.464-3		
PLUOF3	22	0.347	0.389	0.0312	0.567-3		
PLUOF3	23	0.316	0.423	0.0108	0.516-3		
PLUOF3	24	0.438	1.270	0.0117	0.716-3		
PLUOF3	25	0.242	0.125	0.0120	0.454-3		
PLUOF3	26	0.376	1.130	0.0489	0.379-3		
PLUOF3	27	0.485	2.170	0.0646	0.489-3		
PLUOF3	28	0.830	3.250	0.0856	0.299-2		
PLUOF3	29	0.904	4.680	0.0578	0.348-2		
PLUOF3	30	0.389	0.300	0.0257	0.392-3		
PLUOF3	31	0.298	0.335	0.0173	0.301-3		
PLUOF3	32	0.184	0.145	2.21-3	0.186-3		

PLUUF3	33	0.231	0.247	0.0101	0.233-3				
PLUUF3	34	1.490	4.600	0.1870	0.508-1				
PLUUF3	35	1.700	8.200	0.2350	0.633-1				
PLUUF3	36	2.000	12.90	0.2450	0.723-1				
PLUUF3	37	1.490	7.460	0.1800	0.427-1				
PLUUF3	38	1.170	1.370	0.0691	0.366-1				
PLUUF3	39	1.950	3.700	0.8520	0.414-1				
PLUUF3	40	1.97	.447	4.27	.400-1				
PLUUF3	41	1.58	1.27	.408	.399-1				
PLUUF3	45	.504	.236	.134	.329-2				
PLUUF3	46	.486	.197	.134	.317-2				
PLUUF3	47	.347	.109	.151-4					
PLUUF3	47	.347	.109	.151-1	.227-2				
PLUUF3	48	.361	.106	.311-1	.236-2				
PLUUF3	49	.151	.569-2	.683-1	.153-3				
PLUUF3	50	0.75	.134	.017	0.052				
PLUUF3	51	2.04	.415	.024	0.082				
PLUUF3	52	3.09	1.49	.146	0.468				
PLUUF3	53	.320	.728-1	.175-1	.194-2				
PLUUF3	54	.460	.213-1	.863-1	.706-2				
PLUUF3	55	.127	.509-1	.265-2	.106-3				
PLUUF3	56	0.55	0.087	0.133	.007				
PLUUF3	57	.150	.850-1	.617-2	.125-3				
PLUUF3	58	2.14	2.50	2.08	.489-1				
PLUUF3	59	.084	.468-2	.679-2	.696-4				
PLUUF3	75	.320	.728-1	.180-1	.194-2				
PLUUF3	78	3.60	2.70	.432	1.293				
PLUUF3	61	.09	.019	2.79-5	1.09-4				
PLUUF3	42	.39	.073	.002	.008				
PLUUF3	43	.39	.002	.073	.008				
PLUUF3	39	1.95	.852	3.70	.414-1				
PLUUF3	66	.338	.267-1	.759-1	.205-2				
PLUUF3	67	.270	.579	.617-1	.164-2				
PLUUF3	73	.409	.205-1	.220	.165-2				
PLUUF3	74	.284	.42-1	.42-1	.137-2				
PLUUF3	79	.641	.182	.619-1	.692-2				
PLUUF3	80	.619	.145	.614-1	.669-2				
PLUUF3	81	.596	.11	.618-1	.643-2				
PLUUF3	82	.253	.207-1	.613-1	.843-3				
PLUUF3	83	.328	.997-1	.723-1	.109-2				
PLUUF3	84	.403	.267	.793-1	.134-2				
PLUUF3	85	.330	.284-1	.782-1	.145-2				
PLUUF3	86	.400	.314-1	.266	.157-2				
PLUUF3	87	.453	.331-1	.521	.165-2				
MAT1	1	10.3E6	3.86E6		.103				
PAKAM	WIMASS	2.59E-3							
PARAM	GRDPNT	0							
RdE2	190	151	123456	150	200	223	1200		
CUNM2	192	150	1	11.55	-0.39	-3.43	-1.16		
+CON192	3.028		8.733			134.05			+CON192

DYNRED	2	150.0				50	YES	
EIGR	2	GIV				50	1.E-9	+E2
+E2	MAX							
EIGR	1	INV	0.0	50.0	3	3	1.E-10	+E
+E	MAX							
SPUINI	5001	THRU	5050					
ASET1	0	5001	THRU	5050				
SUPORT	774	123	112	12	922	2		
ENDDATA								